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Controlling Molecular Conformation for Highly Efficient and Stable Deep-Blue Copolymer Light-Emitting Diodes

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ABSTRACT

We report a novel approach to the achievement of deep-blue, high-efficiency, and long-lived solution processed polymer light-emitting diodes (PLEDs) via a simple molecular-level conformation change whereby we introduce rigid β-phase segments into a 95% fluorene - 5% arylamine copolymer emission layer (EML). The arylamine moieties at low density act as efficient exciton formation sites in PLEDs whilst the conformational change alters the nature of the dominant luminescence from a broad, charge-transfer like emission to a significantly blue-shifted and highly vibronically structured, excitonic emission. As a consequence, we observe a significant improvement in Commission International de L’Eclairage (CIE) (x, y) co-ordinates from (0.149, 0.175) to (0.145, 0.123) whilst maintaining high efficiency and improving stability. We achieve peak luminous efficiency, $\eta = 3.60$ cd/A and luminous power efficiency, $\eta_w = 2.44$ lm/W; values that represent state of the art performance for single copolymer deep-blue PLEDs. These values are five-fold better than for otherwise-equivalent, β-phase poly(9,9-dioctylfluorene) (PFO) EML PLEDs (0.70 cd/A and 0.38 lm/W). This report represents the first demonstration of the use of molecular conformation as a vector to control the optoelectronic properties of a fluorene copolymer; previous examples have been confined to homopolymers.
INTRODUCTION

Since the discovery of electroluminescence from conjugated polymers in 1989, there has been significant interest in solution processable PLEDs as potential candidates for low cost, energy efficient display and lighting applications. Tremendous efforts have been made to develop deep-blue (usually defined by the electroluminescence (EL) emission having CIE (x, y) coordinates both \( \leq 0.15 \))² light emitting polymers (LEPs) for use in high-luminance and high-efficiency PLED displays,³–⁵ with this requirement essential to achieving the colour gamut needed for high quality display applications. An additional commercialisation challenge is the limited stability of blue fluorescent LEPs,⁶ for which the operational device lifetime is relatively short compared to red and green phosphorescence-based polymer emitters; the latter have encapsulated lifetimes of over 100,000 hours.⁷ In this paper we report a novel approach to the achievement of deep-blue, high-efficiency, and long-lived PLEDs via the introduction of a conformation change in the conjugated backbone of a fluorene-arylamine copolymer.

Conjugated polymers based on fluorene backbones have been extensively studied as blue OLED emission materials on account of both (i) their wide optical gaps, e.g. ~3.0 eV for PFO,⁸ that are favourable for deep-blue emission and (ii) their high photoluminescence (PL) quantum yields, e.g. up to 50-60% for glassy phase poly(9,9-di-octylfluorene) (PFO) (see Figure 1(a)).⁹,¹⁰ They can, however, suffer from poor operational and colour stability arising from fluorenone defects, giving rise to ‘green band’ excimer emission at 535 nm.¹¹–¹³ PFO is an especially well-studied member of the fluorene-homopolymer family and can be prepared with a number of distinct microstructures¹⁴,¹⁵, namely the glassy-,¹⁴–¹⁶ crystalline-,¹⁴,¹⁵,¹⁷,¹⁸ liquid-crystalline (LC)¹⁵,¹⁹–²¹ and chain-segment-extended \( \beta \)-phases.¹⁴,¹⁵,²²–²⁶ The \( \beta \)-phase has an increased
backbone planarity within a fraction of chain segments, with the corresponding torsion angle between adjacent fluorene units $\approx 180^\circ$, resulting in the octyl substituent groups for neighbouring monomers lying on opposite sides of the chain. In contrast, the glassy phase is a disordered phase with a broad distribution of torsion angles between monomers.\textsuperscript{14–16} The different phases can be identified using both optical absorption and PL spectroscopy measurements.\textsuperscript{14,21,25–28}

![Chemical structures](image)

**Figure 1.** Chemical structures for (a) poly(9,9-dioctylfluorene) (100F8 or PFO) and (b) 9,9-dioctylfluorene:butyl substituted phenylenediamine (F8:BSP) copolymers ($x = 0.97, 0.95, 0.90, 0.80$ for 97F8:3BSP, 95F8:5BSP, 90F8:10BSP, 80F8:20BSP and $x = 0.50$ for 50F8:50BSP or PFB).

The $\beta$-phase has attracted much attention due to the action of its extended chain segments as a 'self-dopant'\textsuperscript{20} within an otherwise glassy matrix. These segments constitute the most ordered,
lowest energy states and trigger an efficient energy transfer from the surrounding high energy state glassy segments.\textsuperscript{14,20,22,27,28} The β-phase chain segments have been shown to act as charge carrier-trapping and exciton formation sites, with the extension in conjugation length reducing the optical gap of PFO by some 0.3 eV.\textsuperscript{20,24,27,29–32} As a consequence, β-phase PFO PLEDs have been reported with a nearly two-fold increase in luminous efficiency from 1.0 cd/A to 1.9 cd/A compared to their glassy counterparts.\textsuperscript{30,33} As yet such an approach has not been applied to copolymers. Finally, there are numerous processing methods by which to induce β-phase chain segments, with reported methods including (i) thermal cycling,\textsuperscript{20,21,34} (ii) Langmuir-Bodgett film formation and transfer,\textsuperscript{35} (iii) deposition of films from mixed solvent,\textsuperscript{23} solvent/additive\textsuperscript{33} and high boiling point solvent\textsuperscript{28} solutions, (iv) post-deposition film exposure to solvent vapours,\textsuperscript{21} and (v) dipping in/flooding with solvent.\textsuperscript{23,30}

Approaches to high efficiency within simple PLED device architectures require that individual layers be optimized to perform more than one function. In contrast to blending materials with different functionality\textsuperscript{36}, covalently linked copolymers combine functionalities in a way that is resistant to phase separation. As a consequence, state of the art polymer LED emission materials are now invariably complex copolymers featuring emission and electron- and hole-transport moieties with optimized fractional compositions and chain architectures.\textsuperscript{37–39} In this study we focus on 9,9-diocetylfluorene (F8):butyl-substituted phenylenediamine (BSP) copolymers (see Figure 1(b)) and in particular the 95% F8:5% BSP, or for short 95F8:5BSP, copolymer. During synthesis, the insertion of BSP units into each copolymer chain is subject to the constraint that since BSP-BSP couplings are not possible each BSP unit must have F8 neighbours. The chain formation process is otherwise statistical in nature, dependent on BSP monomer concentration. The low fraction of BSP monomer units in the reaction mixture then ensures that the copolymer
will contain long sequences of F8 units, interrupted only by sparsely distributed BSP units. Further details of the F8:BSP copolymer synthesis are given in the experimental section below. We directly compare copolymer properties with the 100F8 homopolymer PFO and a blend of 90% PFO and 10% poly(9,9-dioctylfluorene-alt-bis-N, N’-(4-butylphenyl)-bis-N,N’-phenyl-1,4-phenylenediamine) (PFB) (with a corresponding volume fraction of 95% F8 and 5% BSP units) which we label 90PFO/10PFB. Further results are presented in the supporting information (SI) for the alternating 50F8:50BSP (PFB) and statistical 97F8:3BSP, 90F8:10BSP and 80F8:20BSP copolymers.

Copolymerisation of BSP and other arylamine moieties into an otherwise F8 polymer backbone reduces the ionization potential from - 5.8 eV versus vacuum for PFO to, for example, - 5.09 eV for PFB and - 4.98 eV for the corresponding alternating copolymer with methoxy substituted phenylenediamine moieties (poly(9,9-dioctylfluorene-co-bis-N,N’-(4-methoxyphenyl)-bis-N,N’-phenyl-1,4-phenylenediamine) (PFMO)). Varying both the arylamine moiety and its fractional incorporation within an F8-based copolymer thus provides an approach to tuning the hole injection and carrier balance properties of blue PLEDs. The associated optical gap remains large but because the emission acquires a more charge-transfer-like character with a broadened spectrum it is no longer suitable to address the display requirement for deep-blue luminescence despite the improvement in efficiency.

Our study looks to combine the electrical benefits of arylamine incorporation with a conformational approach to spectral control. We demonstrate deep-blue, high-efficiency and stable PLEDs by inducing the β-phase conformation within long, uninterrupted F8 chain segments of the 95F8:5BSP copolymer. This allows a significant improvement in CIE (x, y) from (0.149, 0.175) to (0.145, 0.123) and yields peak luminous efficiency, \( \eta = 3.60 \text{ cd/A} \) (at
146.5 cd/m$^2$) and luminous power efficiency, $\eta_w = 2.44$ lm/W (at 10.8 cd/m$^2$). The latter efficiency values represent state of the art performance for simple copolymer deep-blue PLEDs,$^{3,4,46}$ being more than five-fold better than for otherwise-equivalent, glassy PFO EML PLEDs (0.70 cd/A and 0.38 lm/W) and 14 and 60 times higher, respectively, than for our first-reported PFO EML PLEDs (0.25 cd/A and 0.04 lm/W).$^{10}$ Even at 1000 cd/m$^2$ the efficiencies remain high, with $\eta = 3.50$ cd/A and $\eta_w = 1.50$ lm/W.

RESULTS AND DISCUSSION

Optical absorption and photoluminescence spectroscopy. Optical absorption spectra for PFO, 95F8:5BSP copolymer and 90F8/10PFB blend films on spectrosil substrates are shown in Figures 2(a), (b), and (c), respectively, before and after solvent treatment to generate F8 $\beta$-phase chain segments. The glassy-phase spectra of the 95F8:5BSP copolymer (Figure 2 (b)) and 90F8/10PFB blend (Figure 2(c)) films closely resemble that of PFO with the main $\pi-\pi^*$ absorption peak located at $\sim$ 384 nm. After solvent vapor annealing (SVA) with toluene, the characteristic $\beta$-phase absorption peak at 433 nm was observed for all three samples, including the copolymer. This confirms that the BSP moieties within the copolymer backbone do not prevent $\beta$-phase chain segment formation. The small BSP fraction (5 wt%) ensures that there are sufficiently long F8 segments within which the $\beta$-phase can form; oligofluorenes with as few as five 9,9-dioctylfluorene repeat units are reported to show $\beta$-phase spectral features.$^{47}$
Figure 2. Peak normalized optical absorption (solid) and PL emission (dashed) spectra for (a) PFO, (b) 95F8:5BSP copolymer and (c) 90PFO/10PFB blend films spin coated on spectrosil substrates. Glassy phase data are shown with black lines whilst \( \beta \)-phase data are shown with red lines. PL emission spectra were excited at \( \lambda_{\text{ex}} = 385 \text{ nm} \).

The proportion of F8 \( \beta \)-phase segments can be estimated by subtraction of a suitably normalized glassy phase absorption spectrum and comparison of the integrated residual (\( \beta \)-phase) and subtracted areas (Figure S1). Table 1 summarizes the estimated \( \beta \)-phase fractions for each film type.
Table 1. Estimated β-phase fractions and β-phase PL contributions in film samples of PFO, 95F8:5BSP and 90PFO/10PFB

<table>
<thead>
<tr>
<th>Film type</th>
<th>β-phase fraction</th>
<th>β−phase PL contribution</th>
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<tbody>
<tr>
<td>PFO</td>
<td>10% ± 2%</td>
<td>100%</td>
</tr>
<tr>
<td>95F8:5BSP copolymer</td>
<td>5% ± 1%</td>
<td>62%</td>
</tr>
<tr>
<td>90PFO/10PFB blend</td>
<td>12% ± 2%</td>
<td>82%</td>
</tr>
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A higher fraction of β-phase chain segments is formed in films of the PFO homopolymer (10%) and 90PFO/10PFB blend (12%) than of the 95F8:5BSP copolymer (5%). The β-phase fraction in the homopolymer is broadly consistent with previous results in the literature. The 50F8:50BSP alternating copolymer (PFB) does not have any extended sequences of F8 units in which the β-phase can form and hence in the blend only the PFO chains support β-phase segments. We also know that bulky BSP moieties disrupt close chain packing, leading to a more disordered glassy microstructure for PFB with no observed crystallization on thermal annealing and no β-phase formation. It is perhaps not, therefore, surprising that the 95F8:5BSP copolymer has a smaller fraction of β-phase chain segments formed during SVA than PFO. Our study additionally shows that β-phase segment formation still occurs (to a lesser degree) for 90F8:10BSP and (marginally) 80F8:20BSP copolymer films (see Figure S2). In the case of the blend films, the 10% fraction of PFB chains (with 50% BSP content) will not support β-phase segment formation so one might expect a proportionate reduction in overall β-phase fraction relative to PFO. This is not, however, seen; most likely as a result of an increase in free volume that compensates by facilitating conformation change in the PFO chains.
The PL emission spectra for each of the three film types are shown in Figures 2(a), (b) and (c) before and after solvent treatment to generate F8 β-phase chain segments. All spectra were excited at $\lambda_{\text{ex}} = 385$ nm. The glassy homopolymer PFO film spectrum in Figure 2(a) is consistent with literature reports, with $S_1$-$S_0$ 0-0 and 0-1 vibronic peaks at 421 nm and 447 nm, respectively.\textsuperscript{14,24,25,27,28} The glassy 95F8:5BSP copolymer film shows a red-shifted, broad, asymmetric (with long wavelength tail), and largely featureless PL spectrum, which is very similar (but slightly red shifted (455 nm peak) and broadened) to that of the PFB component (450 nm peak) in the blend film and, indeed, the 97F8:3BSP, 90F8:10BSP and 80F8:20BSP copolymers (see Figure S2). The 95F8:5BSP copolymer film spectrum doesn’t reveal an obviously PFO-like (F8-localised) component indicating efficient energy transfer from locally excited F8 excitons to BSP centered excitons.\textsuperscript{49}

The excited states responsible for the glassy copolymer PL have been shown to have significant charge transfer (CT) character\textsuperscript{49} as observed for PFB and poly(9,9-dioctylfluorene-co-N-(4-butylphenyl)diphenylamine) (TFB).\textsuperscript{50} The dilute solution PL spectra of PFO homopolymer, PFB and 95F8:5BSP for solvents of different polarity (toluene, THF and dichlorobenzene) are shown in Figure S3. The PFO PL shows little change with solvent polarity, whilst for PFB there is a large red shift and broadening with increasing polarity, strongly indicative of CT character. For 95F8:5BSP the emission comprises both vibronically-structured F8 and broadband BSP related contributions with the former experiencing no solvatochromic shift whilst the latter red-shifts and broadens, confirming a coexistence of both bound neutral exciton and CT emission states. The CT character originates from a differential spatial partitioning of the HOMO and LUMO wavefunctions across the BSP and F8 units, leading to a displacement in associated hole
and electron charge densities.\textsuperscript{49} This is supported by cyclic voltammetry (CV) measurements (see supporting information Figure S4).

The PL spectra are significantly altered by the generation of $\beta$-phase chain segments. All three film-types then display well-resolved vibronic structure with a close match of the blend and copolymer PL spectral features to those of PFO; the characteristic $\beta$-phase vibronic peaks appear at 437, 465 and 498 nm. However, differences do exist, with the vibronic peaks best resolved for PFO, less so for the blend and least for the copolymer. In addition, the apparent strength of the $S_1$-$S_0$ 0-1 and 0-2 vibronic peaks relative to the 0-0 is greater for the 90PFO/10PFB blend than PFO and greatest for the 95F8:5BSP copolymer, consistent with the blend and copolymer spectra comprising a superposition of $\beta$-phase PFO-like structured excitonic emission and residual PFB-like broadband CT emission. Separation of the 95F8:5BSP copolymer and 90PFO/10PFB blend spectra in this way (Figure S5) reveals that 62\% of the copolymer emission is $\beta$-phase structured emission whilst 38\% is residual PFB-like emission. The 90PFO/10PFB blend shows 82\% $\beta$-phase emission and 18\% residual PFB emission.

Efficient energy transfer from high energy glassy to low energy $\beta$-phase segments in PFO has been extensively studied, with only a few \% of $\beta$-phase segments needed for dominant $\beta$-phase emission.\textsuperscript{14,25,28} Intriguingly, we observe similar behavior here with the presence of $\beta$-phase segments in the copolymer and blend films leading to a strong promotion of structured vibronic emission. This is despite the fact that as a consequence there is a net increase in mean PL emission energy relative to the glassy film spectra (with dominant CT-like emission); this can be explained by the very small Stokes shift ($\lambda \sim 5$nm) for $\beta$-phase segments. $\beta$-phase PL spectral components are also evident in the SVA 97F8:3BSP, 90F8:10BSP and 80F8:20BSP copolymer films but not for PFB (50F8:50BSP alternating copolymer) (Figure S2). As the fraction of BSP
units increases, the fraction of β-phase segments formed after SVA decreases (Table S1),
resulting in larger residual fractions of PFB-like emission. Radiative decay times for the
competing emissive species are likely to be important in this context.

Despite both having the same volume fraction of F8 and BSP units, the glassy phase
90PFO/10PFB blend film PL spectrum contains emission components from both PFO
(evidenced by the shoulder at 425 nm) and PFB (main peak at 450 nm) polymer chains, whereas
the 95F8:5BSP copolymer seemingly does not. This indicates that energy transfer from majority
F8- to minority BSP-centered sites is less efficient in the blend (PFO to PFB inter-chain transfer)
than found in the copolymer (for combined inter- and intra-chain transfer); blend microstructure
will clearly also play a role in this.

The copolymer, in addition, shows less intense emission around 530 nm (see Figure S6(d)),
where ‘green-band’ (fluorenone defect enabled) excimer emission occurs for PFO and related
materials.12,13 A reduction in green-band emission is an advantage of incorporating bulky BSP
moieties within the copolymer backbone, reducing fluorenone-to-fluorenone cofacial π-stacking,
and hence excimer formation and emission.

**Time dependent PL spectroscopy.** To further characterize the emissive species contributing
to the PL spectra, time correlated single photon counting (TCSPC) measurements were used to
record PL decay transients under (unless otherwise indicated) 404 nm excitation and 460 nm
detection. Figure 3 shows results for, from top to bottom, 95F8:5BSP copolymer, 90PFO/10PFB
blend and PFO homopolymer thin films with both glassy (left) and β-phase (right panel)
microstructures. Table 2 collects together the fitted decay times, fractional amplitudes, average
decay times and photoluminescence quantum efficiencies (PLQEs). Additional decay transients
collected at a range of different wavelengths between 420 and 540 nm are shown in Figure S7
and equivalent data for PFB is shown in Figure S8. PL spectra and decay transients were also measured for dilute toluene solution samples of PFO, PFB and 95F8:5PFB and are shown in Figure S9.

**Figure 3.** PL decay transients for (a) glassy- and (b) β-phase films of, from top to bottom, 95F8:5BSP copolymer (red), 90PFO/10PFB blend (blue) and PFO homopolymer (black). The measured instrument response function (IRF) is also shown (dashed grey line). Samples were excited at 404 nm and the resulting PL emission intensity was monitored at 460 nm.

**Table 2.** Decay times and percentage amplitudes obtained by fitting the transient PL data presented in Figure 3 to a multi-exponential decay model. The $\chi^2$ values are indicative of the fit-quality. The average decay time is found as the weighted sum of the fitted lifetimes $\tau_{av} = \sum A_i \tau_i$. Photoluminescence quantum efficiency (PLQE) values are also included.

<table>
<thead>
<tr>
<th>Sample</th>
<th>$\tau_1$(ns):(A$_1$)</th>
<th>$\tau_2$(ns):(A$_2$)</th>
<th>$\tau_3$(ns):(A$_3$)</th>
<th>$\tau_{av}$ (ns)</th>
<th>$\chi^2$</th>
<th>PLQE (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>PFO Glassy</td>
<td>0.314:(97%)</td>
<td>1.34:(3%)</td>
<td>-</td>
<td>0.35</td>
<td>1.21</td>
<td>35 ± 5</td>
</tr>
<tr>
<td>PFO β-phase</td>
<td>0.270:(98%)</td>
<td>1.15:(2%)</td>
<td>-</td>
<td>0.29</td>
<td>1.13</td>
<td>35 ± 5</td>
</tr>
</tbody>
</table>
The glassy PFO films yield a predominantly fast neutral singlet exciton decay with ~ 314 ps time constant, in agreement with decay times reported previously for this microstructure (~300 - 400 ps). A minor (~ 3%) fraction of longer-lived (~1.34 ns) PL is also observed, attributed to inter-chain/-segment states, including non-geminate pairs and fluorenone-defect-based excimers (yielding ‘green-band’ emission) (Figure S7(a)). Consistent with this assignment, the PFO decay transients for dilute solutions are mono-exponential at all detection wavelengths, with ~ 356 ps decay time (Figure S9(b)). For PFO β-phase film samples, the 460 nm decay times reduce marginally to ~270 ps and ~1.15 ns, which (given no PLQE decrease (Table 2)) points to an increase in transition dipole moment, consistent with the known increase in conjugation length.

The PL transients of the 95F8:5BSP copolymer in dilute solution (Figure S9(c)) show a bi-exponential decay at shorter wavelengths (420 and 440 nm), attributed to combined PFO-like excitonic emission (with \( \tau_1 \approx 140 \) ps) and PFB-like CT emission (with \( \tau_2 \approx 1.4 \) ns). At longer wavelengths, beyond 460 nm, a mono-exponential decay with \( \tau \approx 1.4 \) ns is observed, identical to the dilute solution decay for PFB (Figure S9(d)). In the case of glassy 95F8:5BSP copolymer films the overall decay at 460 nm can be described by three components with \( \tau_1 \approx 580 \) ps (43%), \( \tau_2 \approx 2.18 \) ns (45%) and \( \tau_3 \approx 11.9 \) ns (12%), each attributable to CT emission. The average lifetime
is then $\tau_{av} \sim 2.66$ ns. These time constants all differ from the solution CT-state decay, consistent with the influence of heterogeneity in solid-state packing and inter-chain/-segment interactions, as also reported in previous studies of fluorene-amine copolymers.\textsuperscript{49,56,57} We further note that the decay times vary substantially with emission wavelength (Figure S7(c)), again suggesting a distribution of CT lifetimes within the copolymer. Interestingly, the longest-time constant ($\tau_3 \sim 11.9$ ns) decay is not observed in PFB (50F8:50BSP alternating copolymer) (Figure S8) and its observation here then points to a potentially more substantial spatial separation of electron and hole wavefunctions in glassy 95F8:5BSP. One possibility would be inter-chain excitations formed between high electron affinity F8 units in one chain and low ionization potential BSP units in a neighboring chain. However, as no distinct exciplex peak is seen in either PL or EL this remains speculative.

The striking spectral changes that occur when $\beta$-phase chain segments are induced in 95F8:5BSP copolymer films (Figure 2(b)) are accompanied by a strong change in 460 nm PL decay dynamics (Figure 3 and Table 2). Each of the fitted decay times reduces, to $\tau_1 \sim 319$ ps (56%), $\tau_2 \sim 1.32$ ns (35%) and $\tau_3 \sim 7.43$ ns (9%), with a large shift in fractional weighting towards the $\tau_1$ component. The average lifetime correspondingly reduces to $\tau_{av} \sim 1.31$ ns. The majority sub-320 ps decay component is consistent with the spectral dominance of vibronic F8-based emission (Figure 2(b)) whilst the 1.32 and 7.43 ns time constants signal the presence of residual CT emission. As for the glassy case, the decays are slower for longer collection wavelengths (Figure S7(d)). However, unlike the situation for PFO and blend films, a modest decrease in PLQE from 45% to 40% was observed when the $\beta$-phase was induced in copolymer films (Table 2), albeit that the PLQE itself remains relatively high (c.f. 35% for PFO and 25%
for the blend). Among possible explanations, fluorenone-centered quenching is plausible but remains unproven.

The 460 nm PL decay transient for glassy 90PFO/10PFB blend films can also be fit to three exponentials, although their relative fractions and decay times, not surprisingly, differ from those of the 95F8:5BSP copolymer (Table 2). Excitonic emission from the PFO chains is dominant, with CT-like emission accounting for much smaller fractions than in the copolymer; corresponding decay times are $\tau_1 \sim 295$ ps (64%), $\tau_2 \sim 1.48$ ns (27%) and $\tau_3 \sim 8.2$ ns (9%). As for the copolymer, when emission is collected at longer wavelengths, we see that longer lifetime emissive species are increasingly important (Figure S7(e)). In the blend case, the longest time constant excited states ($\tau_3 \sim 8.2$ ns) have previously been shown to be exciplexes generated between PFO and PFB, with thermally assisted energy transfer to a PFB CT-like exciton. Upon $\beta$-phase induction, the decay becomes almost mono-exponential (Figure 3) with excitonic emission from PFO $\beta$-phase chain segments totally dominant; decay times are $\tau_1 \sim 247$ ps (97%) and $\tau_2 \sim 1.93$ ns (3%). The lack of a longer lived time constant indicates the majority excitons generated on PFO chains no longer form exciplexes with PFB and instead efficiently transfer to $\beta$-phase segments before undergoing radiative decay. As detection wavelength is increased (Figure S7(f)), the short-lived decay time remains relatively constant ($\tau \sim 271$-325 ps) until 500 nm, where CT emission increases. The PLQE values for glassy and $\beta$-phase 90PFO/10PFB blends were 20% and 25% respectively, smaller than for both PFO and 95F8:5BSP copolymer films, indicating more substantial non-radiative decay.

**Display-related Device characteristics.** To test the effect of $\beta$-phase segment formation on device performance, PLEDs were fabricated (see experimental methods) with a conventional bottom-emitting device architecture, comprising glass substrate / indium tin oxide (ITO) anode /
poly(3,4-ethylenedioxythiophene) (PEDOT:PSS) hole injection layer / TFB electron blocking interlayer / emission layer (EML) / LiF / Ca / Al cathode. Schematic device and energy level diagrams are shown in Figure 4. The energy levels of PFO, 95F8:5BSP and PFB were deduced from CV measurements (Figure S4 and Table S2).

Figure 4. (a) Device structure of blue PLEDs and (b) corresponding schematic energy level diagram. The polymer energy levels were deduced from CV measurements. The shaded grey area indicates the smaller energy gap for β-phase segments in PFO and 95F8:5BSP, whilst the pale green area indicates the energy gap for BSP units within the 95F8:5BSP copolymer.
Figure 5. PLED characteristics for glassy (filled symbols, (a), (c) and (e)) and β-phase (open symbols, (b), (d) and (f)) EML film microstructure devices. J-V-L data are plotted in (a) and (b), Luminous (cd/A) and luminous power (lm/W) efficiency data as a function of luminance in (c) and (d) and associated EQE data in (e) and (f). 95F8:5BSP copolymer EML data are plotted as
red circles, 90PFO/10PFB blend data as blue triangles and PFO homopolymer data as black squares.

Figures 5 (a) and (b) compare the current density and luminance vs voltage (J-V-L) characteristics for, respectively, glassy and β-phase PFO homopolymer, 95F8:5BSP copolymer and 90PFO/10PFB blend EML devices. Figures 5 (c) and (d) show the corresponding luminance-dependent glassy and β-phase PLED efficiencies \( \eta \) (cd/A) and \( \eta_w \) (lm/W) and Figures 5 (e) and (f) the associated external quantum efficiencies \( \eta_{eqe} \) (EQE). Other parameters (turn-on voltage, peak \( \eta \), \( \eta_w \) and \( \eta_{eqe} \)) for these devices are collated in Table 3 and their electroluminescence (EL) spectra are shown in Figure 6(a).

**Table 3.** Summary of best PLED performance for the current study showing turn-on voltages (defined as the voltage at which luminance reaches 1 cd/m\(^2\)), peak luminous (cd/A), luminous power (lm/W) and external quantum efficiencies (%) and their values at 100 cd/m\(^2\) and 1000 cd/m\(^2\), and CIE (x, y) color coordinates for PFO, 95F8:5BSP copolymer and 90PFO/10PFB blend EML films with both glassy and β-phase microstructures.
<table>
<thead>
<tr>
<th>EML Type</th>
<th>Turn-on Voltage (V)</th>
<th>Luminous efficiency $\eta$ (cd/A)</th>
<th>Luminous Power efficiency $\eta_w$ (lm/W)</th>
<th>External quantum efficiency $\eta_{eqe}$ (%)</th>
<th>CIE (x,y)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>Peak @ 100 cd/m$^2$</td>
<td>Peak @ 1000 cd/m$^2$</td>
<td>Peak @ 100 cd/m$^2$</td>
<td>@ 1000 cd/m$^2$</td>
</tr>
<tr>
<td>PFO Glassy</td>
<td>4.2</td>
<td>0.43 @ 6.8 V</td>
<td>0.36 @ 5.6 V</td>
<td>0.22 @ 6.6 V</td>
<td>0.31 @ 6.6 V</td>
</tr>
<tr>
<td></td>
<td></td>
<td>0.43 @ 1000 cd/m$^2$</td>
<td>0.19 @ 1000 cd/m$^2$</td>
<td></td>
<td>0.26 @ 1000 cd/m$^2$</td>
</tr>
<tr>
<td>PFO $\beta$-phase</td>
<td>4.0</td>
<td>0.70 @ 6.2 V</td>
<td>0.48 @ 5.4 V</td>
<td>0.32 @ 6.0 V</td>
<td>0.43 @ 6.0 V</td>
</tr>
<tr>
<td></td>
<td></td>
<td>0.69 @ 1000 cd/m$^2$</td>
<td>0.37 @ 1000 cd/m$^2$</td>
<td></td>
<td>0.30 @ 1000 cd/m$^2$</td>
</tr>
<tr>
<td>90PFO/10PFB Glassy</td>
<td>3.8</td>
<td>1.52 @ 5.6 V</td>
<td>1.47 @ 4.6 V</td>
<td>0.97 @ 4.6 V</td>
<td>0.89 @ 4.6 V</td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.52 @ 1000 cd/m$^2$</td>
<td>1.48 @ 1000 cd/m$^2$</td>
<td></td>
<td>0.86 @ 1000 cd/m$^2$</td>
</tr>
<tr>
<td>90PFO/10PFB $\beta$-phase</td>
<td>3.8</td>
<td>1.31 @ 5.6 V</td>
<td>1.21 @ 4.4 V</td>
<td>0.79 @ 4.4 V</td>
<td>0.80 @ 4.4 V</td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.28 @ 1000 cd/m$^2$</td>
<td>1.28 @ 1000 cd/m$^2$</td>
<td></td>
<td>0.73 @ 1000 cd/m$^2$</td>
</tr>
<tr>
<td>95F8:5BSP Glassy</td>
<td>3.2</td>
<td>4.05 @ 6.2 V</td>
<td>3.80 @ 4.0 V</td>
<td>2.62 @ 4.0 V</td>
<td>2.40 @ 4.0 V</td>
</tr>
<tr>
<td></td>
<td></td>
<td>4.00 @ 1000 cd/m$^2$</td>
<td>4.00 @ 1000 cd/m$^2$</td>
<td></td>
<td>2.30 @ 1000 cd/m$^2$</td>
</tr>
<tr>
<td>95F8:5BSP $\beta$-phase</td>
<td>3.2</td>
<td>3.60 @ 5.6 V</td>
<td>3.60 @ 4.2 V</td>
<td>2.44 @ 4.2 V</td>
<td>2.40 @ 4.2 V</td>
</tr>
<tr>
<td></td>
<td></td>
<td>3.50 @ 1000 cd/m$^2$</td>
<td>3.50 @ 1000 cd/m$^2$</td>
<td></td>
<td>2.44 @ 1000 cd/m$^2$</td>
</tr>
<tr>
<td></td>
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<td></td>
<td></td>
<td></td>
<td>(0.149, 0.175)</td>
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<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>(0.145, 0.123)</td>
</tr>
</tbody>
</table>
In terms of PLED efficiency, PFO glassy EML devices show a peak $\eta = 0.43$ cd/A at 6.8 V, a peak $\eta_w = 0.22$ lm/W at 5.6 V and a peak $\eta_{eqe} = 0.31\%$ at 6.6 V, with 1000 cd/m$^2$ luminance at 7 V. The PFO $\beta$-phase EML devices show a significant improvement, with peak efficiencies of $\eta = 0.70$ cd/A at 6.2 V, $\eta_w = 0.38$ lm/W at 5.4 V, $\eta_{eqe} = 0.43\%$ at 6.0 V and 1000 cd/m$^2$ luminance now at 5.9 V. This improvement is consistent with previous reports for $\beta$-phase PFO devices.\textsuperscript{30,31,33,59,60} Figure 6(b) shows, however, that the efficiency gains are at the expense of a detrimental change in EL emission color; a consequence not previously emphasized. The shift from glassy (peak $\lambda = 425$ nm) to $\beta$-phase (peak $\lambda = 440$ nm) alters the CIE (x, y) coordinates from (0.155, 0.098) to (0.157, 0.117), with a corresponding shift in dominant wavelength from 470 to 475 nm and a decrease in color saturation from 86 to 82%. This behavior limits the achievable display color gamut. It also helps to explain why the $\eta_{eqe}$ enhancement is more modest than the gains in $\eta$ and $\eta_w$; a shift to the green leads to a better overlap with the photopic eye sensitivity function that peaks at 555nm.

The 90PFO/10PFB blend EML PLEDs show significantly better efficiency characteristics than for the corresponding PFO devices (Figure 5), with the PFB fraction strongly assisting hole injection (Figure 4(b)).\textsuperscript{40,45} The glassy blend EML gives peak efficiencies $\eta = 1.52$ cd/A at 5.6 V, $\eta_w = 0.97$ lm/W at 4.6 V and $\eta_{eqe} = 0.89\%$ at 5.6 V, with a luminance of 1000 cd/m$^2$ at 6.3 V, whilst $\beta$-phase blend EML devices show peak efficiencies $\eta = 1.31$ cd/A at 5.6 V, $\eta_w = 0.79$ lm/W at 4.4 V and $\eta_{eqe} = 0.80\%$ at 5.6 V and reach 1000 cd/m$^2$ at 6.2 V. $\beta$-phase devices are, therefore, somewhat less efficient than glassy devices but in terms of CIE (x, y) color coordinates $\beta$-phase segment formation leads to a shift from (0.162, 0.174) to (0.151, 0.118), resulting in a highly desirable, deeper-blue emission (Figure 6(b)). The corresponding dominant wavelength shifts from 478 to 474 nm and color saturation increases from 73 to 83%. This color
shift also helps to explain at least in part the proportionately larger decrease in $\eta$ and $\eta_w$ values than in $\eta_{eqe}$ as the emission then has reduced overlap with the photopic eye sensitivity function.

Glassy and $\beta$-phase 95F8:5BSP copolymer EML devices show yet further enhanced PLED efficiency (Table 3). Glassy devices give peak efficiencies $\eta = 4.05$ cd/A at 6.2 V, $\eta_w = 2.62$ lm/W at 4.0 V and $\eta_{eqe} = 2.40\%$ at 6.0 V, with a luminance of 1000 cd/m$^2$ at 6.6 V, whilst for $\beta$-phase devices $\eta = 3.60$ cd/A at 5.4 V, $\eta_w = 2.44$ lm/W at 4.2 V and $\eta_{eqe} = 2.40\%$ at 5.6 V, with 1000 cd/m$^2$ at 7.2 V. Interestingly, here $\eta_{eqe}$ is unaltered by $\beta$-phase induction whilst (see Figure 6(b)) the CIE (x, y) coordinates still shift positively from (0.149, 0.175) to (0.145, 0.123), resulting in a dominant wavelength decrease from 479 nm to 474 nm and a color saturation increase from 77% to 85%. Figure S10 shows the deconvolution of the EL spectra of $\beta$-phase copolymer and blend devices, showing $\beta$-phase emission accounts for ~67% of the total EL emission in the copolymer device, and ~80% emission in the blend device.
Figure 6. (a) EL spectra at 6 V for PFO homopolymer, 95F8:5BSP copolymer and 90PFO/10PFB blend devices with glassy (black) and β-phase (red) EML microstructures. (b) Color coordinate (CIE(x,y)) diagram showing the EL emission coordinates for PFO (square), 95F8:5BSP (circle) and 90PFO/10PFB (triangle) for glassy (filled black symbols) and β-phase (filled white symbols) EMLs. Inset shows EL emission coordinates on expanded scale; arrows
indicate color shift from glassy to β-phase for PFO (green arrow), 95F8:5BSP copolymer (blue arrow) and 90PFO/10PFB blend (purple arrow).

The copolymer EML efficiency improvement is largely attributed to the BSP units both assisting hole injection and, due to their sparse distribution, acting as deep hole-traps. Glassy PFO EML devices display the highest turn-on voltages (defined as the applied bias at which L = 1 cd/m²), namely 4.2 V, reducing to 4.0 V on induction of β-phase chain segments, consistent with previous reports.³⁰,³¹,³⁹ The 90PFO/10PFB blend devices conversely show the same turn-on voltage (3.8 V) irrespective of glassy or β-phase microstructure, as also do the 95F8:5BSP copolymer devices (3.2 V); injection is clearly controlled by BSP rather than β-phase units. The better turn-on behavior for the copolymer devices is likely due to the more uniform distribution of BSP units; they are present within every 95F8:5BSP polymer chain. Atomic force microscopy (AFM) topography images show evidence of phase segregation between PFO and PFB chains in 90PFO/10PFB blend samples (Figure S11) along with an increase in roughness (Table S3) whilst the 95F8:5BSP microstructure appears featureless. Evidence of phase segregation between PFO and PFB chains has also been previously observed in the literature.⁵⁸

Other studies have additionally shown that such copolymers have three to four orders of magnitude lower time-of-flight photocurrent hole mobility (µ⁺ToF) than equivalent films of PFO and alternating fluorene-arylamine copolymers.⁴⁰,⁴³,⁴⁵ The measured µ⁺ToF ≈ 10⁻⁷ cm²/Vs for glassy films indicates substantially deeper hole trapping than for β-phase segment formation in PFO. Electrochemical data (Figure S4 and Table S2) for 95F8:5BSP films suggest that the trap depth is ~ 0.3 eV above the ~ 5.80 eV HOMO level of PFO. This trapping effect is also evident in the reduction in current density between the PFO homopolymer and 95F8:5BSP copolymer EML devices. We find J = 231.4 mA/cm² at 7 V for glassy PFO, and only 41.4 mA/cm² at the
same voltage for glassy 95F8:5BSP despite the turn-on voltage for the copolymer (3.2 V) being 1 V lower than for PFO (4.2 V).

A well-known drawback to using PFO as the EML in PLEDs is its low spectral stability under operation due to the appearance of a low energy ‘green-band’ emission when driven at higher voltages, the origin of which is inter-chain emission from fluorenone-based defects. Figure 7 shows EL spectra as a function of applied voltage for glassy and β-phase PFO and 95F8:5BSP devices. As the voltage is driven beyond 7 V for the glassy PFO device and beyond 9 V for the β-phase PFO device, the emergence of a broad, low energy component (‘green-band’) is clearly observed. However for both glassy and β-phase 95F8:5BSP devices, no green-band is observed when driving devices between 5 and 13 V, indicating that 95F8:5BSP devices show much improved color stability compared with PFO PLED devices. This is likely to be due to the bulky BSP units causing increased disruption in the packing structure between the polymer chains, thereby preventing green band emission; similar reductions in green band have been observed when small amounts of carbazole units or large amine group end caps have been incorporated into polyfluorene chains.
Figure 7. EL spectra as a function of voltage for (a) glassy phase PFO, (b) β-phase PFO, (c) glassy phase 95F8:5BSP and (d) β-phase 95F8:5BSP PLED devices

The BSP-centered hole trapping increases the likelihood of exciton formation and improves charge carrier balance.\textsuperscript{6,64} Copolymer devices based on 97F8:3BSP and 90F8:10BSP EMLs were also fabricated (Figures S12 and S13), revealing that the 95F8:5BSP copolymer gives the optimal EML efficiency performance in this composition sequence (Table S4); one needs enough but not too much BSP incorporation and the way in which the BSP units are incorporated also matters.

In summary, the device efficiency and display color parameter data (Table 4) show that in the absence of BSP units (PFO homopolymer) the induction of β-phase chain segments is advantageous to device efficiency but only at the expense of a shift away from desirable deep
blue emission. The introduction of BSP units into the EML (blend and copolymer) yields a significant overall enhancement in device efficiency and color stability with the copolymer performing substantially better than the blend. Again, though, the efficiency enhancement is achieved at the expense of color response with the BSP-related CT-like emission yielding less deep-blue color coordinates (green shift in dominant wavelength and reduction in saturation). Induction of β-phase chain segments is then strongly beneficial to the color response and for copolymer EML devices this occurs without any appreciable decrease in quantum efficiency; the β-phase copolymer EML PLEDs thus provide the best overall combination of efficiency and color performance.

Table 4. Display color parameters for PFO, 90PFO/10PFB blend and 95F8:5BSP copolymer EML PLEDs. Note the close agreement for β-phase EML color saturation and dominant wavelength for all three EML types and likewise the close agreement for glassy EML copolymer and blend devices.

<table>
<thead>
<tr>
<th>EML Type</th>
<th>EQE (%) @ 1000 cd/m²</th>
<th>CIE (x,y)</th>
<th>Dominant Wavelength (nm)</th>
<th>Color Saturation (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>PFO Glassy</td>
<td>0.31</td>
<td>0.155, 0.098</td>
<td>470</td>
<td>86</td>
</tr>
<tr>
<td>PFO β-phase</td>
<td>0.43</td>
<td>0.157, 0.117</td>
<td>475</td>
<td>82</td>
</tr>
<tr>
<td>90PFO/10PFB Glassy</td>
<td>0.87</td>
<td>0.162, 0.174</td>
<td>478</td>
<td>73</td>
</tr>
<tr>
<td>90PFO/10PFB β-phase</td>
<td>0.78</td>
<td>0.151, 0.118</td>
<td>474</td>
<td>83</td>
</tr>
<tr>
<td>95F8:5BSP Glassy</td>
<td>2.39</td>
<td>0.149, 0.175</td>
<td>479</td>
<td>77</td>
</tr>
<tr>
<td>95F8:5BSP β-phase</td>
<td>2.37</td>
<td>0.145, 0.123</td>
<td>474</td>
<td>85</td>
</tr>
</tbody>
</table>
As a final device test, motivated by literature reports that β-phase formation in PFO increases device lifetime, encapsulated pre- and post-SVA 95F8:5BSP copolymer PLEDs were subjected to accelerated lifetime testing under nitrogen using a constant current source set to deliver 4 mA (i.e. $J \approx 90$ mA/cm$^2$ for the 4.5 mm$^2$ pixels under test). The luminance was measured at 60-second intervals starting from 2821 cd/m$^2$ for glassy and 2300 cd/m$^2$ for β-phase devices. Initially (c.f. Figure S14), the glassy and β-phase 95F8:5BSP PLED luminance values decayed at a similar rate, with half decay times $T_{50\%}$ (glassy) = 176 mins and $T_{50\%}$ (β-phase) = 180 mins. The subsequent decay was much more rapid in the glassy copolymer EML devices, especially beyond 400 mins. The luminance took 490 mins to drop to 30% of its starting value for the glassy 95F8:5BSP PLED, but took 630 mins to reach the same fractional output for the β-phase device, and whilst the latter was still emitting some 250 cd/m$^2$ at 1200 mins the glassy device luminance had fallen below 4 cd/m$^2$ by that time. The operational stability of β-phase 95F8:5BSP copolymer devices is, therefore, significantly greater than that of otherwise equivalent glassy devices. Several factors are expected to contribute to this including a desirable distribution of the recombination sites vertically through the EML, enhanced charge carrier balance and a faster radiative decay time for the exciton.

We now consider further the effects of β-phase formation on the energy transfer of emissive species in order to explain the desirable PLED emission for 95F8:5BSP copolymer EML devices. We explore, in particular, the origin of efficient energy transfer to β-phase F8-centred excitons, with concomitant net increase in mean photon energy and deeper-blue emission.
The HOMO and LUMO of glassy phase PFO homopolymer films were deduced from CV measurements to be 5.80 eV and 2.10 eV respectively (Figure S4 and Table S2), similar to previously reported values.\textsuperscript{42} Upon β-phase chain segment formation a smaller optical gap component is introduced into the ensemble of absorbing chromophores, with the resolved $S_0$-$S_1$ 0-0 peak at 433 nm.\textsuperscript{24,27} Additionally, as a result of rapid energy migration to β-phase segments, there is a ~ 16 nm red-shift in $S_1$-$S_0$ 0-0 PL emission. CV-measurement-based HOMO and LUMO values were also determined for PFB and 95F8:5BSP, yielding 5.05 eV and 2.00 eV and 5.49 eV and 2.10 eV, respectively (Table S2).

The schematic energy level diagram for glassy 95F8:5BSP copolymer chains (Figure 8(a)) shows how the energy levels vary spatially along the copolymer chain, with the electron-rich BSP unit being raised in energy relative to the F8 units. Whilst electrical excitation should (at least initially) predominantly produce CT excitons due to the strong hole trapping nature of the BSP units,\textsuperscript{6,38,43} under optical excitation exciton states will also form on longer F8 segments (labelled I in Figure 8). In this case, efficient inter- (not shown) and intra-chain energy transfer of F8 excitons to CT states (labelled II in Figure 8) is observed.\textsuperscript{49}
Figure 8. Schematic energy level diagrams (left column) for intra-chain energy transfer processes in 95F8:5BSP (a) glassy- and (b) β-phase chains following optical excitation. Jablonski diagrams for each process are shown to the right of each schematic. See text for explanation of numbering I, II, and III.

Upon β-phase formation, the corresponding F8 HOMO level will move up and the LUMO down, forming a Type I quantum-well-like structure along the polymer chain (Figure 8(b)). Under optical excitation, excitons formed locally on glassy F8 segments undergo efficient energy transfer to either a β-phase F8 segment (labelled III in Figure 8) or BSP-based CT states before decaying (Figure 8(b)). As such, the 95F8:5BSP β-phase PL spectrum is a superposition of both PFO-like β-phase emission (~62% of total emission) and residual CT emission (~38%). The reason β-phase emission dominates the spectrum despite the film having roughly the same proportion of β-phase segments and BSP units (~5%) is likely to be a result of the faster decay for β-phase F8 excitons. Longer-lived CT states will be more prone to thermal energy transfer.
before decay. Nevertheless, as expected, the fraction of CT type emission increases (and β-phase emission correspondingly decreases) as more BSP units are incorporated into the copolymer chain (see Figure S2 for 90F8:10BSP and 80F8:20BSP UV-vis absorption and PL emission data).

To further probe the energy transfer mechanism in the 95F8:5BSP copolymer, additional low temperature PL measurements (λ<sub>ex</sub> = 400 nm) were taken for β-phase PFO and both glassy and β-phase 95F8:5BSP copolymer samples (Figure S15) from 290K down to 10 K. Low temperature β-phase 95F8:5BSP PL measurements (Figure S15(c) and (d)) reveal that the spectrum remains a superposition of both vibronic β-phase emission (Figure S15(a)) and CT-like emission from BSP centered states (Figure S15(b)) over the whole 10 to 290 K range. Spectral deconvolution reveals that the component spectra are identical to the PFO β-phase and glassy 95F8:5BSP spectra at each temperature, as for example shown at 10 K Figure S15(e). By integrating the component spectra at each temperature we find that (Figure S15(f)) the fraction of β-phase emission decreases as temperature decreases, from ~64% at 290 K to ~32% at 10 K.

These results support the energy transfer model outlined above since energy transfer in β-phase PFO has been proposed to be a two-step process of thermally assisted exciton diffusion followed by Forster resonance energy transfer from high energy glassy phase F8 to low energy β-phase segments. The decrease in β-phase emission at low temperature could then be explained as a result of energy transfer to β-phase units being more dependent on thermally assisted diffusion than is the case for BSP sites. Alternatively, enhanced polaron formation has been previously observed for β-phase PFO and polarons are known to act as emission quenching sites, especially at low temperatures where their lifetimes are long, leading to an increasing β-
phase exciton quenching at low temperatures.\textsuperscript{24} Similar effects are seen in other polymers when segmentation by conjugation breaks leads to polaron trapping.\textsuperscript{66}

Under electrical excitation of copolymer EML devices, it is likely that a much greater fraction of initially formed excitations are CT states, due to the hole trapping nature of the BSP unit (trap depth \textasciitilde0.3 eV). Despite this, F8 $\beta$-phase emission provides the dominant (\textasciitilde67\%) contribution to the EL spectrum, with residual CT-like emission delivering a more modest (\textasciitilde33\%) contribution. This suggests that either endothermic energy transfer from BSP-centered CT states to $\beta$-phase F8 excitons or ‘trap-filling’ of BSP sites occurs. In the latter case, as the BSP ‘traps’ fill, $\beta$-phase F8 charge localization is expected to play an increasing role. Evidence for this is clearly seen in Figure S16, where the relative fraction of $\beta$-phase exciton EL emission increases with voltage. The shorter emission decay time for F8 $\beta$-phase segment- than BSP-localized states also feeds into the higher relative fraction of $\beta$-phase EL.

A key difference for 90PFO/10PFB blend samples is that inter-molecular energy transfer between neighboring PFB and PFO chains becomes important. Under optical excitation of the glassy-phase blend, excited states are formed directly on both PFO and PFB chains. However, due to the raised HOMO and LUMO energies of PFB relative to PFO, a type II heterojunction occurs at the PFO/PFB interfaces, resulting in exciplex formation.\textsuperscript{58} At room temperature, the exciplex undergoes endothermic energy transfer to the emissive PFB CT state yielding its characteristic spectrum (Figure S17).\textsuperscript{58} However, not all the PFO excitons generated will form an exciplex and we consequently also observe PFO exciton emission as a shoulder at 425 nm (Figure 2(c)). Under electrical excitation, as holes will tend to localize on PFB chains and electrons on PFO (see Figure 4 energy levels) it is likely that exciplexes are directly formed by electron-hole Coulomb capture\textsuperscript{67} before undergoing endothermic energy transfer to PFB CT.
states (Figure S18). Consistent with this, Figure 6(a) indeed shows a much weaker 425 nm PFO emission shoulder for glassy phase blend EL than PL.

For β-phase blends, the majority of optically excited glassy segment F8 excitons will tend to transfer their energy to β-phase sites without forming exciplexes (Figure S17). Additionally, under electrical excitation, exciplex states that form between PFB and β-phase PFO chains (from direct electron-hole capture) will likely undergo endothermic energy transfer to PFO β-phase excitons (Figure S18). As a consequence, both the PL (Figure 2(c)) and EL (Figure 6(a)) β-phase spectra are dominated by structured vibronic emission. Some residual CT emission remains, respectively 18% and 20% for PL and EL (Figures S5 and S10), but these values are significantly lower than for the 95F8:5BSP copolymer case despite the same volume fraction of BSP units being present. This is most likely due to the greater fraction of β-phase segments generated and the consequences of phase separation.

CONCLUSION

In summary, we have demonstrated a novel route to high-efficiency, deep-blue emitting PLEDs by introducing a simple molecular-level conformation change in the F8 sequences of 95F8:5BSP copolymers from the disordered glassy phase to the rigid β-phase microstructure via annealing in a toluene solvent atmosphere. UV-visible absorption and PL spectroscopy measurements of solvent annealed F8:BSP copolymers at 3%, 5%, 10% and 20% BSP fraction showed β-phase formation in the F8 segments, with the appearance of characteristic red-shifted absorption peaks and the promotion of well-structured vibronic emission.

Incorporating 5% BSP units into the conjugated backbone of an otherwise F8 polymer (thus yielding 95F8:5BSP) produces a five-fold performance enhancement in PLED luminous
efficiency and luminous power efficiency relative to PFO (i.e., 100F8). The BSP units (with significantly lower ionization potential) enhance hole injection and act as hole-trapping sites that assist efficient exciton formation. The BSP units have an additional benefit in increasing the color stability of the PLEDs by suppressing green-band emission when driven at higher voltages. The BSP units also, however, undesirably shift the EL to a CT-state-based sky-blue emission spectrum (CIE \((x, y) = (0.149, 0.175)\)). Subsequent introduction of \(\beta\)-phase chain segments within the copolymer restores a more desirable deep-blue, vibronically-structured EL emission with CIE \((x, y) = (0.145, 0.123)\) whilst retaining the bulk of the efficiency enhancement; it also increases operational stability. The overall best PLEDs, using \(\beta\)-phase 95F8:5BSP EMLs, then have \(\eta = 3.60 \text{ cd/A at 5.4 V and } \eta_w = 2.44 \text{ lm/W at 4.2 V, with 1000 cd/m}^2 \text{ luminance at 7.2 V.}\)

The spatial distribution of BSP units in the PLED active layer is also found to be significant to device function, with 95F8:5BSP copolymer (homogeneous) and 90PFO/10PFB blend (heterogeneous) EMLs showing distinct properties despite containing the same volume fraction of BSP units. Both the glassy- and \(\beta\)-phase 90PFO/10PFB blend devices performed less well than equivalent copolymer devices with luminous efficiency, luminous power efficiency and EQE lower by a factor 2.5-3.0 but still significantly better than for PFO-only EML devices.

Our study represents the first demonstration of the use of a simple molecular-level chain conformation change as a vector to control the optoelectronic properties of a fluorene-based copolymer. It will be interesting to see how broadly our conformation control approach can be applied to materials systems for other device applications, including solar energy conversion, electronics and sensing.

**EXPERIMENTAL**
**Materials.** PFO, 95F8:5BSP and 50F8:50BSP were supplied by Cambridge Display Technology Ltd and used as received. The weight-average molecular weights ($M_w$) and polydispersity indices (PDIs) were $1.04 \times 10^5$ g/mol and 2.3 for PFO, $2.77 \times 10^5$ g/mol and 2.4 for 95F8:5BSP and $1.32 \times 10^6$ g/mol and 2.0 for PFB (50F8:50BSP). Additional samples of 97F8:3BSP, 95F8:5BSP, 90F8:10BSP, 80F8:20BSP and 50F8:50BSP copolymers were provided by the Sumitomo Chemical Company Tsukuba Research Laboratory and again used as received. All (100-x)F8:xBSP polymers were synthesized via Suzuki coupling, with the fraction x controlled by the monomer composition of the reaction mixture. This comprised 50% boronic ester disubstituted F8, (50-x)% bromine disubstituted F8 and x% bromine disubstituted BSP. The coupling process links carbon atoms with a boronic ester substituent to carbon atoms with a bromine substituent and thus creates F8-F8 and F8-BSP linkages but not BSP-BSP linkages.

**UV-Vis absorption, photoluminescence spectroscopy and transient photoluminescence decay measurements.** Optical spectroscopy was undertaken on ~60 nm thickness PFO homopolymer, F8:BSP copolymer and polymer/polymer blend films spin-coated from toluene solution (10 mg/ml) onto quartz substrates. The substrates had been pre-cleaned by sequential 15 minute sonications in acetone, isopropanol, and detergent (Hellmanex III, 2% by volume in DI water) prior to a 3 minute plasma ash in air at 80 watts.

Steady state transmittance was measured using a Shimadzu UV-2550 UV-visible spectrophotometer. Absorbance was calculated directly from transmittance (no scattering or reflection correction) based on the natural logarithm and the substrate contribution simply subtracted. Photoluminescence spectra were recorded in reflection geometry using a Jobin Yvon Horiba Fluoromax-3 spectrofluorometer (excitation wavelength $\lambda_{ex} = 385$ nm). PLQE was
measured using a Jobin Yvon Horiba Fluoromax-3 spectrofluorometer equipped with a diffusely reflecting integrating sphere.

Time-resolved photoluminescence decay measurements used an IBH fluorescence lifetime spectrometer operating in time-correlated single photon counting (TSPC) mode. The excitation source was a 404 nm pulsed LED operating at 1 MHz rep-rate with a pulse temporal width of 200 ps. IBH Datamax software was used to deconvolve the instrument response function from the data and to fit multi-exponential decay functions.

**Low temperature photoluminescence measurements.** Samples were held inside a helium filled, closed-cycle cryostat, with spectra recorded at 20 K intervals. The cryostat temperature was held constant for 5 minutes prior to each measurement, enabling the sample to reach thermal equilibrium. Excitation (0.13 mW) was with 400 nm light from a monochromated supercontinuum laser source (Fianium). PL was collected at right angles to excitation using a 100 µm diameter optical fibre and fed into a spectrometer (Andor SR-163) equipped with a CCD detector (Andor i-Dus). To enhance the signal to noise ratio more than one hundred 0.1 s duration measurements were averaged. A dark background was subtracted from all spectra before correcting with a calibration file measured for a standard light source that accounts for the detector spectral response. The optical geometry of the experiment was unaltered between sequential temperature measurements.

**PLED fabrication and characterization:** The PLED device architecture consists of a multilayer stack comprising ITO/PEDOT:PSS/TFB/EML/LiF/Ca/Al. ITO anode structures on glass substrates (size 12 mm x 8 mm) were cleaned for 15 minutes each in a sequence of ultrasonic baths using acetone, isopropanol and detergent (Hellmanex III, 2% by volume in DI water). This was followed by oxygen plasma ashing in an Emitech K1050X. Next, a 35 nm
thickness film of PEDOT:PSS (Clevios P VP) was deposited as hole injecting layer by spin-coating at 3000 rpm and annealing in air for 15 min at 135 °C. This was followed by spin-coating (at 1000 rpm) a 15 nm thickness electron-blocking TFB interlayer from 2 mg/ml toluene solution, and then baking in nitrogen at 180 °C for one hour. The 95F8:5BSP EML (60 nm thickness) was deposited on top of the TFB interlayer, again by spin coating, at 2500 rpm from a 10 mg/ml toluene solution. For the 90PFO/10PFB blend EML, PFO and PFB toluene solutions were separately prepared and mixed to give the desired weight ratio before spin-coating (2500 rpm) to a thickness of 60 nm. Finally for PFO EML samples, toluene solutions were spin-coated (2500 rpm) to a thickness of 60 nm. To induce β-phase chain segments in the EMLs, each sample was solvent vapor annealed in a toluene atmosphere at 50 °C for 2 hours. An MBraun thermal evaporator was used to deposit the top cathode comprising a triple layer of LiF (2 nm), calcium (30 nm) and aluminum (100 nm).

PLEDs were characterized at room temperature in a sealed sample chamber under nitrogen, using a computer-controlled Keithley Source Measure unit to apply a bias voltage to the chosen pixel (each substrate accommodated 6 PLED pixels) and to measure the resultant current. A Minolta LS100 spot luminance meter measured the corresponding pixel luminance and electroluminescence spectra were recorded using an Ocean Optics USB 2000 CCD spectrometer equipped with a fiber light collection bundle. Accelerated lifetime testing was performed using the same experimental apparatus.

**ASSOCIATED CONTENT**

**Supporting Information.** The supporting information (PDF) contains: UV-vis absorption and photoluminescence measurements for 97F8:3BSP, 90F8:10BSP, 80F8:20BSP and PFB;
additional TCSPC measurements in solution and thin film; deconvolution of UV-vis and PL measurements; characterization of 97F8:3BSP and 90F8:10BSP PLEDs; and additional energy transfer diagrams for 90PFO/10PFB blends.

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Notes

The authors declare no competing financial interest.

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