Charged-particle angular correlations in XeXe collisions at $\sqrt{s_{NN}} = 5.44$ TeV

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Azimuthal correlations of charged particles in xenon-xenon collisions at a center-of-mass energy per nucleon pair of $\sqrt{s_{\mathrm{NN}}} = 5.44$ TeV are studied. The data were collected by the CMS experiment at the LHC with a total integrated luminosity of 3.42 $\mu\mathrm{b}^{-1}$. The collective motion of the system formed in the collision is parametrized by a Fourier expansion of the azimuthal particle density distribution. The azimuthal anisotropy coefficients v_2, v_3, and v_4 are obtained by the scalar-product, two-particle correlation, and multiparticle correlation methods. Within a hydrodynamic picture, these methods have different sensitivities to noncollective and fluctuation effects. The dependence of the Fourier coefficients on the size of the colliding system is explored by comparing the xenon-xenon results with equivalent lead-lead data. Model calculations that include initial-state fluctuation effects are also compared to the experimental results. The observed angular correlations provide new constraints on the hydrodynamic description of heavy ion collisions.

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I. INTRODUCTION

At sufficiently high temperatures or densities, lattice quantum chromodynamics predicts a transition from ordinary hadronic matter to a state of deconfined quarks and gluons, the so-called quark gluon plasma (QGP) (see, e.g., Ref. [1]). The QGP state can be reached through relativistic heavy ion collisions, where the collective behavior of the created medium manifests itself in azimuthal correlations among the emitted particles. These correlations have been studied in gold-gold collisions at the BNL RHIC [2–5], lead-lead (PbPb) collisions at the CERN LHC [6–8], as well as in collisions involving lighter nuclei, such as the copper-copper system studied at RHIC [9,10]. More recently, collective behavior similar to that observed in collisions of heavy nuclei has also been found in high-multiplicity events produced in the proton-lead (pPb) system, and in proton-proton (pp) collisions [11–14]. The results from these small systems raise the question as to how the size of the colliding system affects the onset of QGP formation. Measurements from xenon-xenon (XeXe) collisions, as presented here, bridge the gap between the small (pp and pPb) and large (PbPb) systems previously studied at LHC energies.

Anisotropic flow can be characterized by a Fourier expansion [15–17],

$$\frac{2\pi}{N}\frac{dN}{d\phi} = 1 + \sum_{n=1}^{\infty} 2v_n \cos[n(\phi - \Psi_n)], \qquad (1)$$

where $dN/d\phi$ is the azimuthal particle density and ϕ is the particle azimuthal angle with respect to a reference angle Ψ_n. Different reference angles can be defined. The "participant plane" angle is the direction of the semiminor axis of the region perpendicular to the beam direction spanned by the nucleons that undergo a primary interaction. The "event-plane" angle is defined by the direction perpendicular to the beam direction of the maximum outgoing particle density. In this paper the measured anisotropies are expressed in terms of the event-plane reference angle. Averaged over many events, the anisotropies measured with respect to the event plane are expected to be similar to those that would be obtained if it were possible to determine the actual participant plane.

The magnitude of the azimuthal anisotropy is characterized by the Fourier coefficients v_n. The second- and third-order Fourier coefficients are referred to as "elliptic" (v_2) and "triangular" (v_3) flow, respectively. The former reflects the lenticular shape of the collision overlap region, as well as initial-state fluctuations in the positions of nucleons at the moment of impact [18]. The latter is largely a consequence of fluctuations. While the v_2 and v_3 harmonics are believed to reflect the initial-state geometry [19], for $n \geqslant 4$ the flow harmonics are also strongly affected by the dynamics of the system expansion. Hence, studying both the lower and higher flow harmonics is important for understanding the medium created in heavy ion collisions.

This analysis presents measurements of the charged-particle collective flow in XeXe collisions at a center-of-mass energy per nucleon pair of $\sqrt{s_{\mathrm{NN}}} = 5.44$ TeV. The results are shown as functions of transverse momentum, p_{T}, for the pseudorapidity region $|\eta| < 2.4$ and for different collision overlap geometries. Spectrum-weighted values with $0.3 < p_{\mathrm{T}} < 3.0$ GeV/c, with the efficiency-corrected yield in each p_{T} interval used as the weight, are also presented. The Fourier coefficients v_2, v_3, and v_4 are obtained by two-particle correlations ($v_n\{2\}$), the scalar-product method ($v_n\{\mathrm{SP}\}$), and multiparticle cumulant analyses ($v_n\{m\}$, $m = 4$, 6, and 8).

Event-by-event fluctuations in the spatial overlap geometry lead to method-dependent differences in the extracted

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v_n values [20,21]. The fluctuations cause an increase in the deduced v_n values found using two-particle correlations and the scalar-product method, as compared to the corresponding participant plane value, while the four-particle cumulant v_n results are decreased. For fluctuations that follow a two-dimensional Gaussian behavior, the flow harmonics based on more than four particles are expected to be the same as the four-particle correlations results. Deviations from this common behavior can be used to estimate the higher-order moments of the fluctuation distribution. Comparison of flow coefficients measured by different methods probes the initial-state conditions.

The XeXe values are compared to the results from PbPb collisions at $\sqrt{s_{\mathrm{NN}}} = 5.02$ TeV. The comparison with measurements from different collision systems, but with similar collision geometry, can give insight to the system size dependence of the anisotropic flow [22]. Theoretical predictions are compared to the observed system size dependence of the flow harmonics. The results presented here provide new information on the initial-state geometry and its fluctuations, as well as the system size dependence of the medium properties.

II. CMS DETECTOR

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity coverage provided by the barrel and endcap detectors. Muons are detected in gas-ionization chambers embedded in the steel flux-return yoke outside the solenoid. The hadron forward (HF) calorimeter uses steel as an absorber and quartz fibers as the sensitive material. The two HF calorimeters are located 11.2 m from the interaction region, one on each end, and together they provide coverage in the range $3.0 < |\eta| < 5.2$. These calorimeters serve as luminosity monitors, are used to establish the event centrality, and provide the event-plane information for the scalar-product analysis. The HF calorimeters are azimuthally subdivided into 20° modular wedges and further segmented to form $0.175 \times 10^{\circ}(\Delta\eta \times \Delta\phi)$ towers. The silicon tracker measures charged particles within the pseudorapidity range $|\eta| < 2.5$. For nonisolated particles of $1 < p_{\mathrm{T}} < 10$ GeV/c and $|\eta| < 1.4$, the track resolutions are typically 1.5% in p_{T} and 25–90 (45–150) μm in the transverse (longitudinal) impact parameter [23]. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [24]. The detailed Monte Carlo (MC) simulation of the CMS detector response is based on GEANT4 [25].

III. EVENTS AND TRACK SELECTION

Results based on data recorded by CMS during the LHC runs with XeXe collisions at $\sqrt{s_{\mathrm{NN}}} = 5.44$ TeV in 2017, with an integrated luminosity of 3.42 $\mu\mathrm{b}^{-1}$, are compared to similar data obtained in 2015 from PbPb collisions at $\sqrt{s_{\mathrm{NN}}} = 5.02$ TeV with an integrated luminosity of 26 $\mu\mathrm{b}^{-1}$. In both systems, only tracks with $|\eta| < 2.4$ and $0.3 < p_{\mathrm{T}} < 10.0$ GeV/c are used.

For the XeXe events, a hardware level (level-1) trigger required at least one tower of the HF calorimeters to be above a threshold that was fixed to maximize the number of events counted, while keeping low the noise contamination from electromagnetic scattering and from pileup (i.e., multiple interactions in the same or neighboring bunch crossings). This trigger also required the presence of both colliding bunches at the interaction point. The average online pileup fraction was 0.018 per event. In addition, a high-level trigger was applied that required at least one track in the pixel detector. Events are further selected offline by requiring at least 3 GeV of energy being detected in each of three HF calorimeter towers on either side of the CMS detector and to have a reconstructed primary vertex, containing at least two tracks, located within 15 cm of the nominal collision point along the beam axis and within 0.2 cm in the transverse direction. In addition, contamination from beam-gas interactions are suppressed by applying a filter where, for each event with more than ten tracks, at least 25% of the tracks are required to satisfy a *high purity* [23] track quality criteria. The event selection efficiency is 95%. The track reconstruction algorithm is similar to that used for pp collisions [23].

For PbPb collisions, as compared to XeXe events, there is an additional level-1 trigger requirement of a coincidence between signals in the HF calorimeters on either side of the CMS detector. While offline event selection is similar for PbPb and XeXe events, for the PbPb events the filter to suppress beam-gas interaction is not applied and pileup contamination is controlled by following the procedure outlined in Ref. [26].

To ease the computational load for high-multiplicity central PbPb collisions, track reconstruction for PbPb events is done in two iterations. The first iteration reconstructs tracks from signals ("hits") in the silicon pixel and strip detectors compatible with a trajectory of $p_{\mathrm{T}} > 0.9$ GeV/c. The second iteration reconstructs tracks compatible with a trajectory of $p_{\mathrm{T}} > 0.2$ GeV/c using solely the pixel detector. In the final analysis, the first iteration tracks with $p_{\mathrm{T}} > 1.0$ GeV/c are combined with pixel-detector-only tracks with $p_{\mathrm{T}} < 2.4$ GeV/c, after removing duplicates.

In this paper only tracks from primary charged particles are considered. For the XeXe tracks and the PbPb tracks with both silicon pixel and strip hits, the impact parameter significance of the tracks with respect to the primary vertex in both the beam direction (d_Z) and the transverse plane (d_0) must be less than three standard deviations, while the relative p_{T} uncertainty ($\sigma_{p_{\mathrm{T}}}/p_{\mathrm{T}}$) must be below 10%. In addition, each track is required to have at least 11 hits in the tracker, and the chi-square per degree of freedom, associated with fitting the track trajectory, normalized to the total number of layers with hits along the trajectory, $\chi^2/\mathrm{dof}/\mathrm{layers}$, must be less than 0.15. For the PbPb pixel-only tracks, it was required that d_Z be less than eight standard deviations and that $\chi^2/\mathrm{dof}/\mathrm{layers} < 12$.

IV. ANALYSIS TECHNIQUES

The analysis techniques used in this study are fully described in previous CMS publications. A two-particle correlation analysis, as discussed in Refs. [27,28], is performed for both the XeXe and PbPb data sets. In addition, scalar-product and multiparticle cumulant analyses, as described in Ref. [29], are done for the XeXe data.

In the two-particle correlation analyses, a charged particle from one transverse momentum interval is used as a "trigger" particle, to be paired with all of the remaining charged particles from either the same or a different p_{T} interval, the "associated" particles. For a given trigger particle, the pairing is done in bins of pseudorapidity and azimuthal angle $(\Delta\eta, \Delta\phi)$. A similar pairing between the particles randomly chosen from two different events is done to establish a background distribution. A Fourier analysis of the azimuthal correlation between the trigger and associated particles leads to $V_{n\Delta}$ Fourier coefficients, where n is the Fourier order. If factorization is assumed, the two-particle coefficients can be expressed in terms of single-particle coefficients, with $V_{n\Delta}({p_{\mathrm{T}}}^{\mathrm{trig}}, {p_{\mathrm{T}}}^{\mathrm{assoc}}) = v_n\{2\}({p_{\mathrm{T}}}^{\mathrm{trig}})v_n\{2\}({p_{\mathrm{T}}}^{\mathrm{assoc}})$. The $v_n({p_{\mathrm{T}}}^{\mathrm{assoc}})$ term is given by $\sqrt{V_{n\Delta}({p_{\mathrm{T}}}^{\mathrm{assoc}}, {p_{\mathrm{T}}}^{\mathrm{assoc}})}$, thereby allowing $v_n({p_{\mathrm{T}}}^{\mathrm{trig}})$ to be determined.

In order to minimize statistical uncertainties, the associated particles are taken from a wide p_{T} range with large average anisotropic flow. In this analysis, $1.0 < {p_{\mathrm{T}}}^{\mathrm{assoc}} < 3.0$ GeV/c. To avoid short-range, nonflow correlations, a pseudorapidity gap of $|\Delta\eta| > 2$ is required for the particle pairs.

The scalar-product event-plane measurements are based on recentered flow Q vectors, defined as

$$\vec{Q}_n = \left(\sum_i^M w_i \cos(n\phi_i) - \left\langle \sum_i^M w_i \cos(n\phi_i) \right\rangle, \sum_i^M w_i \sin(n\phi_i) - \left\langle \sum_i^M w_i \sin(n\phi_i) \right\rangle \right).$$

Here, w_i is a weight for the ith particle emitted at azimuthal angle ϕ_i. The summations are over the number of particles M within a given (centrality, η range, p_{T} range) analysis bin for a given event. The averages indicated by the angular brackets are taken over all particles in all events within each analysis bin. These averages correspond to the recentering operation and are needed to minimize detector acceptance effects. If the Q vectors are presented as the corresponding complex scalars, the flow coefficients are given by

$$v_n\{\mathrm{SP}\} \equiv \frac{\langle Q_n Q_{nA}^* \rangle}{\sqrt{\frac{\langle Q_{nA} Q_{nB}^* \rangle \langle Q_{nA} Q_{nC}^* \rangle}{\langle Q_{nB} Q_{nC}^* \rangle}}}. \tag{2}$$

The particles of interest are used to obtain the Q_n vector, with unit weighting ($w_i = 1$) in the sum. The subscripts A, B, and C refer to three separate reference vectors established in different η regions. The product of Q_n with the Q_{nA} reference vector correlates the particles of interest with particles detected in the HF calorimeter (region A). For the current measurement particles of interest with $-0.8 < \eta < 0.0$ $(0.0 < \eta < 0.8)$ and within different p_{T} ranges are correlated with HF particles in the range $3 < \eta < 5$ $(-5 < \eta < -3)$. The products with Q vectors B and C are used to correct for finite resolution effects. The Q_{nC} vector corresponds to particles detected in the HF calorimeter opposite to that used to define the Q_{nA} vector. The Q_{nB} vector corresponds to particles measured in the tracker with $|\eta| < 0.5$. Since the $v_n(p_{\mathrm{T}})$ coefficients increase with p_{T} up to $\approx$3 GeV/c, the choice of either p_{T} or E_{T} weighting results in a better event-plane resolution than with unit weighting. The Q_{nA} and Q_{nC} vectors use E_{T} weighting, whereas the Q_{nB} vector uses p_{T} weighting [30].

The Q-cumulant method is used in this analysis to obtain the four- ($v_n\{4\}$), six- ($v_n\{6\}$), and eight- ($v_n\{8\}$) particle nth-order harmonic results by correlating unique combinations of four, six, and eight particles within each event. The method uses a generic framework described in Ref. [31]. This framework allows for a track-by-track weighting to correct for the detector acceptance effects. A wider pseudorapidity range with $|\eta| < 2.4$ is used for the cumulant method analysis, as compared to the scalar-product method, to reduce statistical uncertainties.

Results are presented in ranges of collision centrality. The centrality variable is defined as a fraction of the inelastic hadronic cross section, with 0% corresponding to full overlap of the two colliding nuclei. The event centrality is determined offline and is based on the total energy measured in calorimeters located in the forward pseudorapidity region $3 < |\eta| < 5$. The analysis is performed in 11 centrality classes, with intervals ranging from 0–5% to 60–70%. By comparing the XeXe and PbPb results in given centrality ranges, similar collision overlap geometries can be achieved, albeit with different numbers of participants.

In comparing the XeXe and PbPb results for more peripheral collisions, it needs to be noted that the XeXe results can be affected by an experimental bias introduced by the centrality determination. Multiplicity fluctuations in the forward region used to determine the event centrality can reduce the centrality resolution. Monte Carlo studies using the HYDJET event generator indicate this bias can be as large as 5% in the 50–60% centrality range and 10% in the 60–70% range for the $v_n\{2\}$ coefficients. For the $v_n\{4\}$ coefficients, the bias is less than 5% in the 60–70% centrality range. For more central events, the bias is found to be negligible.

V. SYSTEMATIC UNCERTAINTIES

Four different sources of systematic uncertainties are considered. To study the effect of the track selection on the final results, different track criteria are applied by varying the limits for the impact parameter significance from 2 to 5, and the relative p_{T} uncertainty from 5% to 10%. These variations are found to have a 1% influence on v_n results for peripheral collisions, increasing to 10% for the most central collisions at the lowest p_{T} values. The effect of moving the primary vertex position along the beam axis is studied by comparing the results with events from the vertex position ranges $|z_{\mathrm{vtx}}| < 3$ cm and $3 < |z_{\mathrm{vtx}}| < 15$ cm to the default range of $|z_{\mathrm{vtx}}| < 15$ cm. A 1% systematic uncertainty is attributed to this source. The systematic uncertainty resulting from the XeXe centrality calibration is estimated by varying the event selection criteria. This uncertainty is largest for the most peripheral centrality

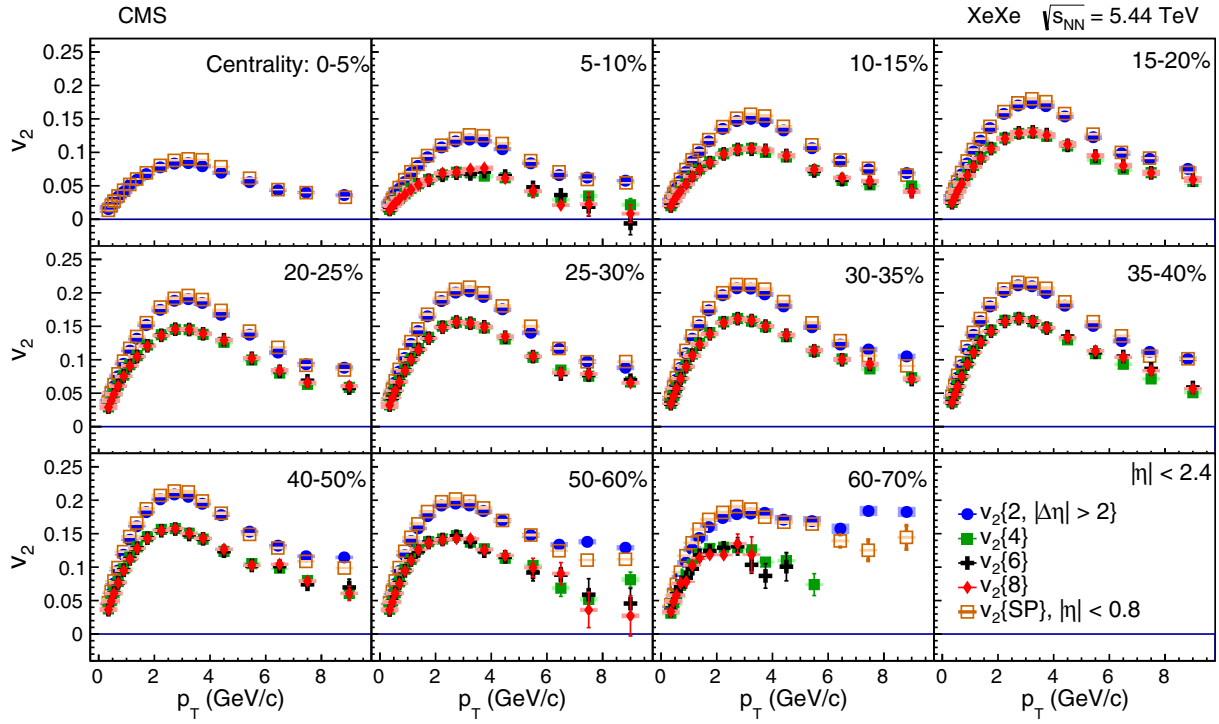

FIG. 1. Elliptic-flow coefficients v_2 based on different analysis techniques, as functions of transverse momentum and in bins of centrality, from the 5% most central (top left) to 60–70% centrality (bottom right). The results for the two-particle and multiparticle correlations correspond to the range $|\eta| < 2.4$, while the scalar-product results are for $|\eta| < 0.8$. The bars and the shaded boxes represent statistical and systematic uncertainties, respectively.

bin, where it reaches a value of 3%. To explore the sensitivity of the results to the MC simulations on which the efficiency determinations are based, analyses using the HYDJET 1.9 [32] event generator are done for generated tracks both before and after the detector effects are taken into account. The results for the two cases differ by about 2% for most centrality ranges, but the difference increases to 10% for the most central events and the lowest track transverse momenta, 0.3–0.4 GeV/c. The observed differences are included as a systematic uncertainty. The different uncertainty sources are independent and uncorrelated, therefore the total systematic uncertainty is obtained by combining the individual contributions in quadrature.

VI. RESULTS

Figure 1 shows the v_2 results, as a function of p_T and in 11 centrality bins, as measured with the different techniques. The two- and multiparticle correlation results are averaged over the pseudorapidity range of $|\eta| < 2.4$, while the scalar-product results are based on tracks with $|\eta| < 0.8$. The elliptic flow values extracted from two-particle correlations show the same pattern as with the multiparticle correlations, but with higher magnitudes. The difference in the results obtained from the two different methods can be largely ascribed to event-by-event fluctuations of the v_2 coefficient [20]. The v_2 magnitude increases with p_T, reaching a maximum value of 0.21 around 3–4 GeV/c in the 30–35% centrality range, and then slowly decreases. The maximum shifts to a lower p_T value as the events become more peripheral. Whereas $v_2\{SP\}$ is found to be generally larger than $v_2\{2, |\Delta\eta| > 2\}$, as expected for the narrower range near midpseudorapidity used for the scalar-product analysis, the situation switches at higher p_T values for centralities >30%. This might reflect a larger nonflow contribution to the two-particle correlation results. The pseudorapidity gap of two units used in the two-particle correlation analysis is less effective in removing non-flow effects, as compared to the gap of three units used for the scalar-product analysis. In the most peripheral events, the $v_2\{2\}$ distribution becomes almost flat for $p_T >$ 3.0 GeV/c. This may be a consequence of nonflow, dijet correlations dominating the results as the system size becomes small.

Figure 2 shows the v_3 values. The difference between the two- and four-particle v_3 values are larger than found for the corresponding v_2 values, exceeding a factor of 2. This suggests a larger fluctuation component to triangular flow as compared to elliptic flow. The difference in amplitude would be qualitatively expected if the v_3 correlations were dominated by initial-state fluctuations [18]. For most centralities, the four-particle distributions have no clear peak value and their p_T dependence is not as prominent as that found for the two-particle and scalar-product methods. The $v_3\{m > 4\}$ values could not be reliably determined because of their large statistical uncertainties. The $v_3\{2\}(p_T)$ distribution has a similar shape as found for the $v_2\{2\}(p_T)$ distribution, but with smaller values that approach zero, or even become negative, at higher p_T values in the most peripheral centrality ranges.

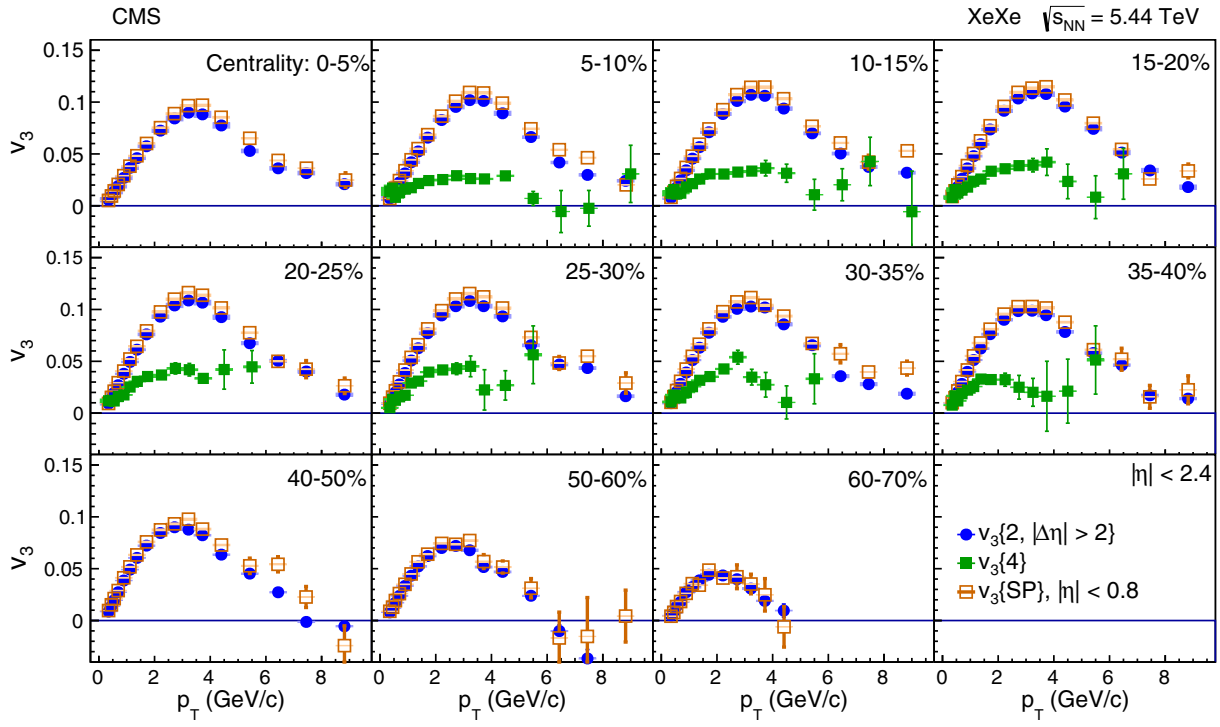

FIG. 2. Triangular-flow coefficients v_3 based on the different analysis techniques, as functions of transverse momentum and in bins of centrality, from the 5% most central (top left) to 60–70% centrality (bottom right). The results for the two-particle and multiparticle correlations correspond to the range $|\eta| < 2.4$, while the scalar-product results are for $|\eta| < 0.8$. The bars and the shaded boxes represent statistical and systematic uncertainties, respectively.

The v_4 results from the two-particle correlation and scalar-product methods are presented in Fig. 3. The Q-cumulant results are not shown because of statistical limitations. The shape of the $v_4(p_T)$ distribution is similar to those for the other measured harmonics. All three harmonics, with $n = 2$, 3, and 4, are found to have maxima at similar p_T values, but with the $n = 3$ and $n = 4$ harmonics having a reduced centrality dependence as compared to the $n = 2$ harmonic. For all three harmonics, the scalar-product values are systematically larger than the two-particle correlation results. While fluctuation effects are expected to affect both methods in a similar way, the methods measure flow in different pseudorapidity ranges, which might account for the observed difference. The similarity of the results suggests there is only a weak pseudorapidity dependence for all three harmonics.

The spectrum-weighted, single-particle anisotropy coefficients, using the two- and multiparticle correlation methods, are presented in Fig. 4. The v_2 coefficients show a strong centrality dependence with a maximum value in the 40–50% centrality bin. The v_3 and v_4 coefficients have only a weak dependence on centrality. Results based on multiparticle cumulants are below the $v_n\{2\}$ values, as expected for the influence of flow fluctuations. The predictions of the IP-GLASMA+MUSIC+UrQMD model are compared to the experimental $v_n\{2\}$ results. In this model, initial-state dynamics are described by impact parameter dependent flowing Glasma gluon fields [33]. The subsequent hydrodynamic evolution is calculated with a MUSIC simulation [34], which is a relativistic $(3 + 1)D$ model that includes shear viscosity (with a shear viscosity over entropy ratio $\eta/s = 0.16$) and a temperature-dependent bulk viscosity over entropy ratio $[\zeta/s(T)]$ [35]. The simulation finally switches from a fluid-dynamic description to a transport description using the ultrarelativistic quantum molecular dynamics (UrQMD) model at the hadronization hypersurface [36]. The theoretical calculations are in good agreement with data for the v_2 and v_4 values. For the v_3 coefficient, the calculation gives slightly larger values than observed, with the difference increasing as the size of the nuclear overlap region decreases (i.e., increasing centrality percentage).

Figure 5 shows the ratios $v_2\{6\}/v_2\{4\}$, $v_2\{4\}/v_2\{2\}$, and $v_3\{4\}/v_3\{2\}$. Theoretical predictions from a hydrodynamic model [37] calculation that uses T$_R$ENTo initial conditions [38] and from the IP-GLASMA+MUSIC+UrQMD model are compared to the experimental results. The former starts the hydrodynamic evolution at a time $\tau = 0.6$ fm/c and has a shear viscosity to entropy ratio of $\eta/s = 0.047$. Xenon is known to be a deformed nucleus with a quadrupole deformation of $\epsilon_2 = 0.15$ [39]. The T$_R$ENTo calculations are performed assuming both spherical and nominally deformed xenon nuclei. The $v_2\{4\}/v_2\{2\}$ ratio shows a strong centrality dependence, with the greatest deviation from unity, with a value of 0.625, corresponding to 5–10% central events. The $v_3\{4\}/v_3\{2\}$ and $v_2\{6\}/v_2\{4\}$ ratios show little, if any, centrality dependence. The $v_3\{4\}/v_3\{2\}$ has a value close to 0.55 for all centralities, indicating a strong influence of fluctuations on triangular flow [20]. The $v_2\{6\}/v_2\{4\}$ ratio is a few percent below unity and suggests the existence of higher-order

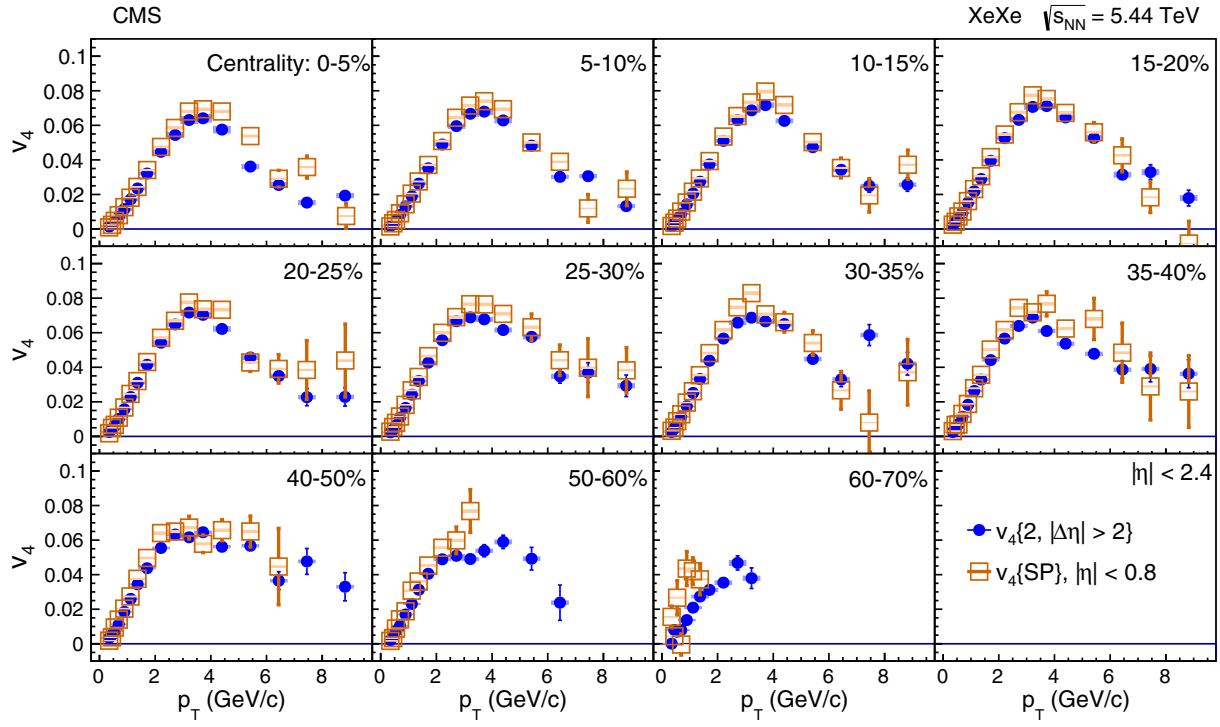

FIG. 3. The v_4 coefficients, based on the different analysis techniques, as functions of transverse momentum and in bins of centrality, from the 5% most central (top left) to 60–70% centrality (bottom right). The results for the two-particle correlations correspond to the range $|\eta| < 2.4$, while the scalar-product results are for $|\eta| < 0.8$. The bars and the shaded boxes represent statistical and systematic uncertainties, respectively.

corrections to a near-Gaussian distribution of the event-by-event flow fluctuations [40]. The IP-GLASMA+MUSIC+UrQMD and hydrodynamic models give comparable agreement with data for the flow harmonic ratios. No significant difference is found between the calculations that assume spherical and deformed Xe nuclear shapes. This suggests that the fluctuations are not sensitive to the small deformation associated with the nucleus.

The v_2 coefficients obtained by the two-particle correlations technique for XeXe collisions at $\sqrt{s_{NN}} = 5.44$ TeV are compared with corresponding PbPb data at 5.02 TeV as a function of transverse momentum in various centrality

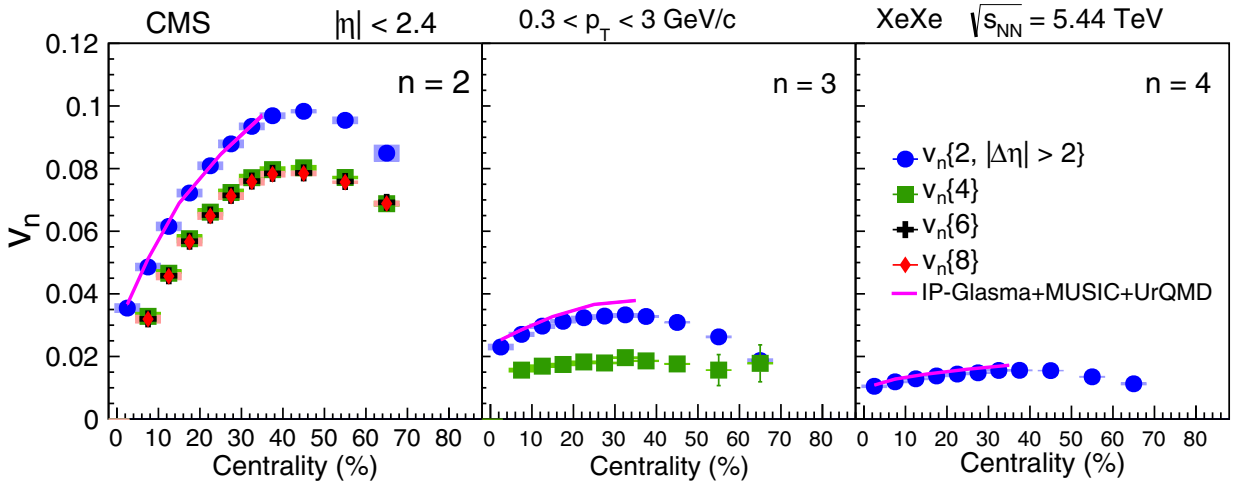

FIG. 4. Centrality dependence of the spectrum-weighted v_2, v_3, and v_4 flow harmonics with $0.3 < p_T < 3.0$ GeV/c. The v_2 results are shown for two-, four-, six-, and eight-particle correlations (left panel). The v_3 results are shown for two- and four-particle correlations (middle panel), while the v_4 values are presented for two-particle correlations technique, only. The solid curve in each panel is the IP-GLASMA+MUSIC+UrQMD prediction for $v_n\{2\}$. The shaded boxes represent systematic uncertainties.

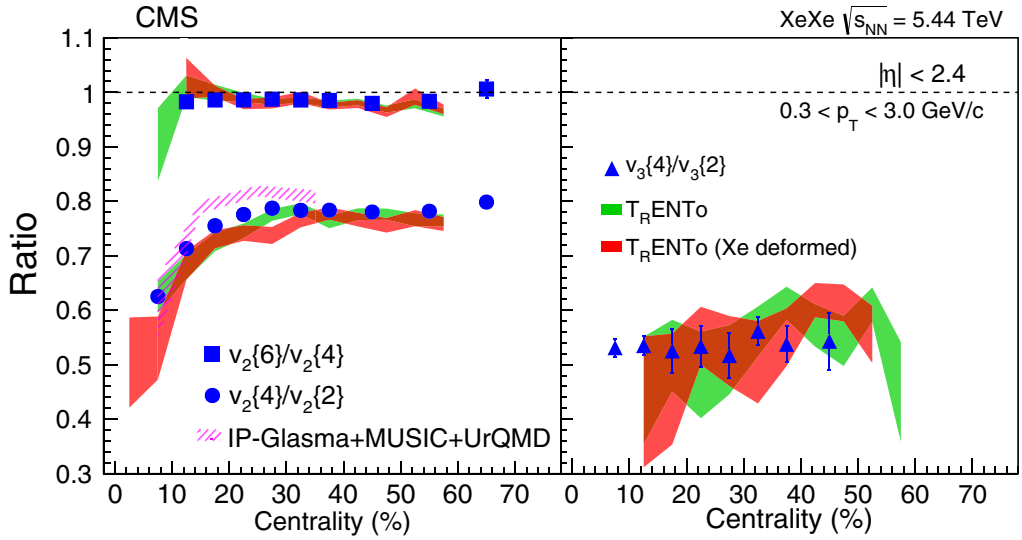

FIG. 5. Centrality dependence of $v_2\{4\}/v_2\{2\}$, $v_2\{6\}/v_2\{4\}$ (left panel), and $v_3\{4\}/v_3\{2\}$ (right panel) ratios. The shaded bands represent the theoretical predictions based on the IP-GLASMA+MUSIC+UrQMD and the relativistic hydrodynamic model from Ref. [37] considering both spherical and deformed xenon nuclei, while the widths of the areas show the statistical uncertainties of the model. The T_RENTo calculation is done for the p_T range $0.2 < p_T < 5.0$ GeV/c.

bins in Fig. 6. The v_2 values for the two systems show similar dependence on p_T. However, the maximum value of the PbPb elliptic flow coefficient is found to be greater than the corresponding XeXe value except in the 0–5% centrality bin. Since, for the most central collisions, the participant fluctuations in the initial-state geometry provide the dominant contribution to the final spacial anisotropy, lower values of v_2 in that region are expected [37] for PbPb collisions because of the larger system size. The $v_3\{2, |\Delta\eta| > 2\}$ coefficients for the two systems are compared in Fig. 7. The v_3 harmonic is entirely generated by initial participant fluctuations, so slightly larger values are expected in XeXe than in PbPb for

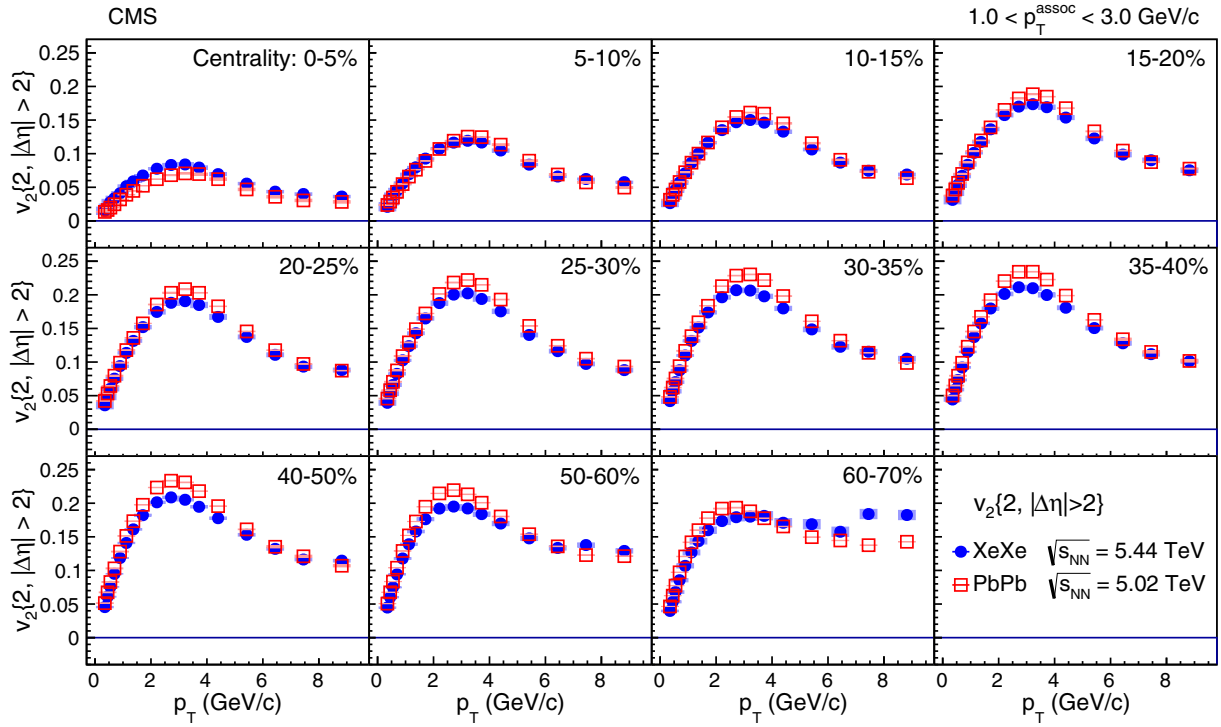

FIG. 6. Comparison of the v_2 results measured with two-particle correlations from two different systems, XeXe collisions at $\sqrt{s_{NN}} =$ 5.44 TeV and PbPb collisions at 5.02 TeV, shown as a function of p_T in eleven centrality bins. The bars (smaller than the marker size) and the shaded boxes represent statistical and systematic uncertainties, respectively.

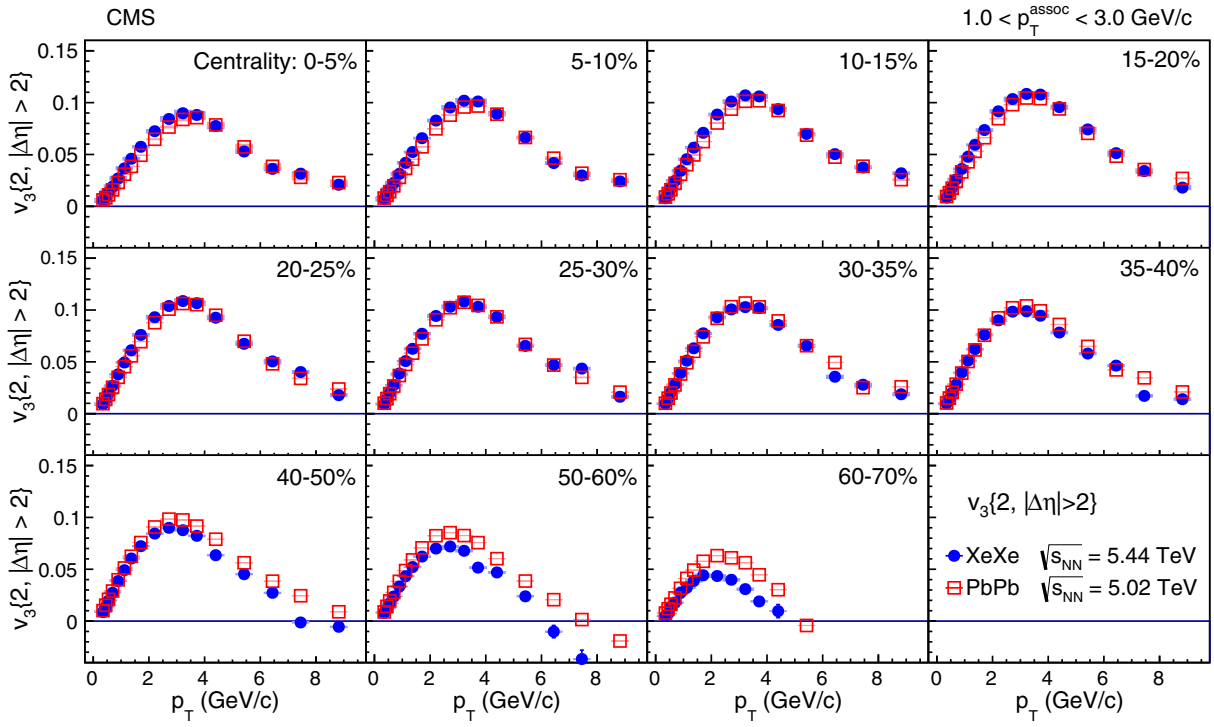

FIG. 7. Comparison of the v_3 results measured with two-particle correlations from two different systems, XeXe collisions at $\sqrt{s_{\mathrm{NN}}} = 5.44$ TeV and PbPb collisions at 5.02 TeV, shown as a function of p_{T} in 11 centrality bins. The bars (smaller than the marker size) and the shaded boxes represent statistical and systematic uncertainties, respectively.

central events (e.g., 0–30% centrality), as observed in the data. However, the v_3 harmonic has a larger sensitivity to transport coefficients (i.e., the shear viscosity) of the created medium, which tends to suppress the azimuthal anisotropy, especially for systems with a small size. This might explain the trend of v_3 where the system with the larger value is reversed in the 30–70% centrality range, with the larger PbPb system showing slightly higher v_3 values for more peripheral events. The $v_4\{2, |\Delta\eta| > 2\}$ coefficients in PbPb and XeXe collisions are shown in Fig. 8. Higher v_4 values are found for PbPb collisions, as compared to the corresponding XeXe collision results, except for the transverse momentum interval $p_{\mathrm{T}} <$ 3.0 GeV/c in the 5% most central events. The ordering of the measured harmonics between the two systems is consistent with participant fluctuations having a dominant role in central collisions, and viscosity effects becoming more important for mid-central and peripheral collisions.

Since ideal hydrodynamics is scale invariant, the XeXe and PbPb results should have similar behavior [37]. For the same percentage centrality range, the interaction regions of the two colliding systems will have similar average shapes, but will have different size. For example, in the 30–40% centrality class, the number of participating nucleons is about 1.6 times higher for the PbPb collisions. However, initial-state fluctuations and viscosity corrections can cause scale invariance breaking. Fluctuations of the initial state are proportional to $A^{-1/2}$, where A is the atomic mass, and, therefore, one can expect a larger fluctuation component for XeXe collisions than for PbPb collisions [41]. However, the influence of the localized fluctuations will decrease with increasing viscosity. The viscosity is thought to be proportional to $A^{-1/3}$ [42] and is therefore also expected to be larger for XeXe collisions. Although the hydrodynamic model simulations do not suggest a large effect on the $v_n\{4\}/v_n\{2\}$ and $v_2\{6\}/v_2\{4\}$ ratios based on the Xe deformation, this deformation can influence the ratio of the XeXe and PbPb results. The quadrupole deformation of the colliding nuclei is expected to have the greatest influence for the XeXe v_2 values corresponding to the most central collisions [37].

Figure 9 shows the p_{T} dependent ratios of XeXe and PbPb harmonic coefficients for different centrality ranges. The ratios reach a maximum value between 1 and 2 GeV/c, within the current uncertainties, and then decrease up to $p_{\mathrm{T}} \sim$ 6 GeV/c, at which point they start to increase again. The increasing trend above 6 GeV/c, which is most pronounced for the v_2 coefficient, might be a consequence of back-to-back dijet correlations that cannot be fully eliminated with the $|\Delta\eta| > 2$ requirement. This nonflow behavior is increasingly significant as the system size becomes smaller, with correspondingly smaller particle multiplicities.

Figure 10 compares the spectrum-weighted v_2, v_3, and v_4 values with $0.3 < p_{\mathrm{T}} < 3.0$ GeV/c for the XeXe and PbPb systems. The largest difference between the two systems is found for the v_2 coefficients corresponding to the most central events, where the XeXe results are larger by a factor of about 1.3. For centralities above 10%, the PbPb results become higher and the ratio has only a weak centrality dependence. For the v_3 and v_4 coefficients, the ratio $v_n[\mathrm{XeXe}]/v_n[\mathrm{PbPb}]$

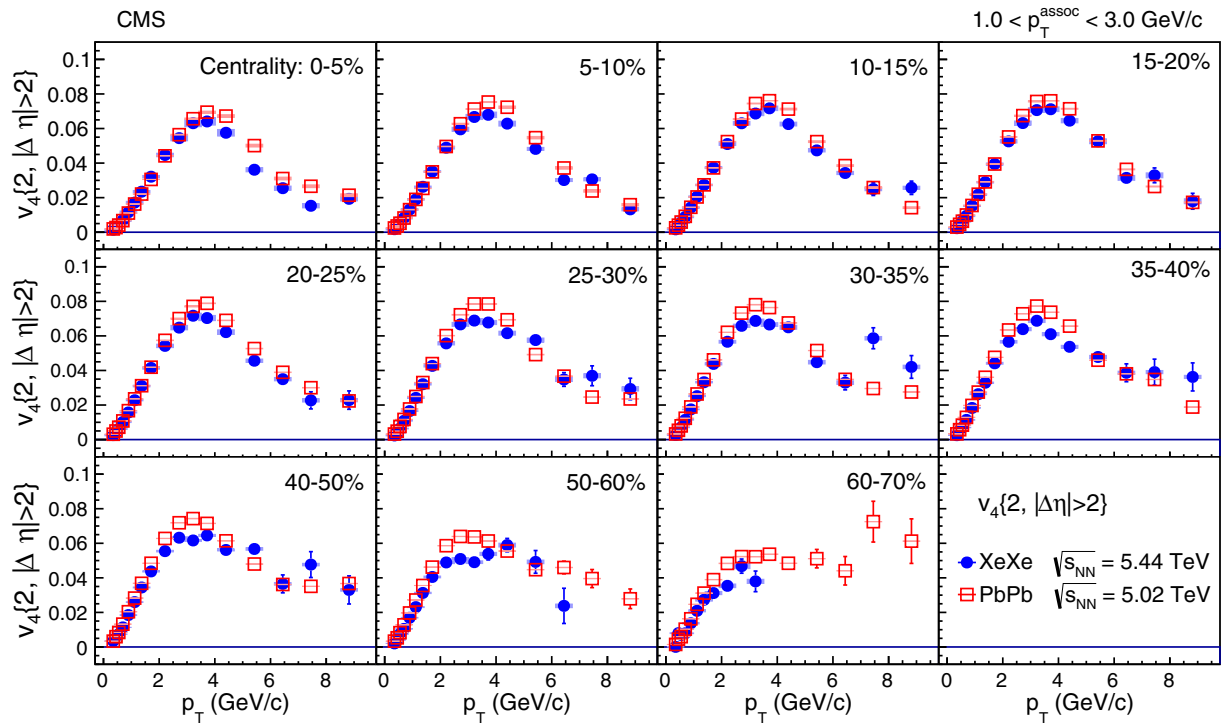

FIG. 8. Comparison of the v_4 results measured with two-particle correlations from two different systems, XeXe collisions at $\sqrt{s_{\mathrm{NN}}} = 5.44$ TeV and PbPb collisions at 5.02 TeV, shown as a function of p_{T} in 11 centrality bins. The bars and the shaded boxes represent statistical and systematic uncertainties, respectively.

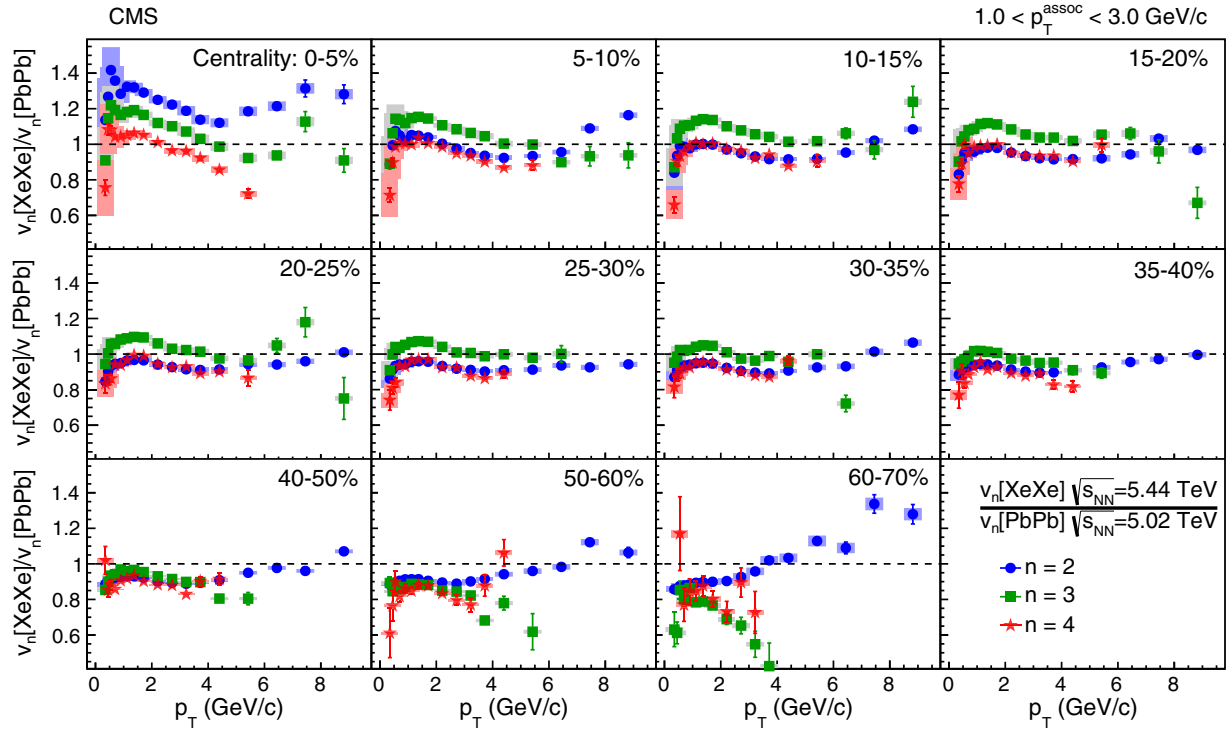

FIG. 9. Ratios of the v_2, v_3, and v_4 harmonic coefficients from two-particle correlations in XeXe and PbPb collisions as functions of p_{T} in 11 centrality bins. The bars and the shaded boxes represent statistical and systematic uncertainties, respectively.

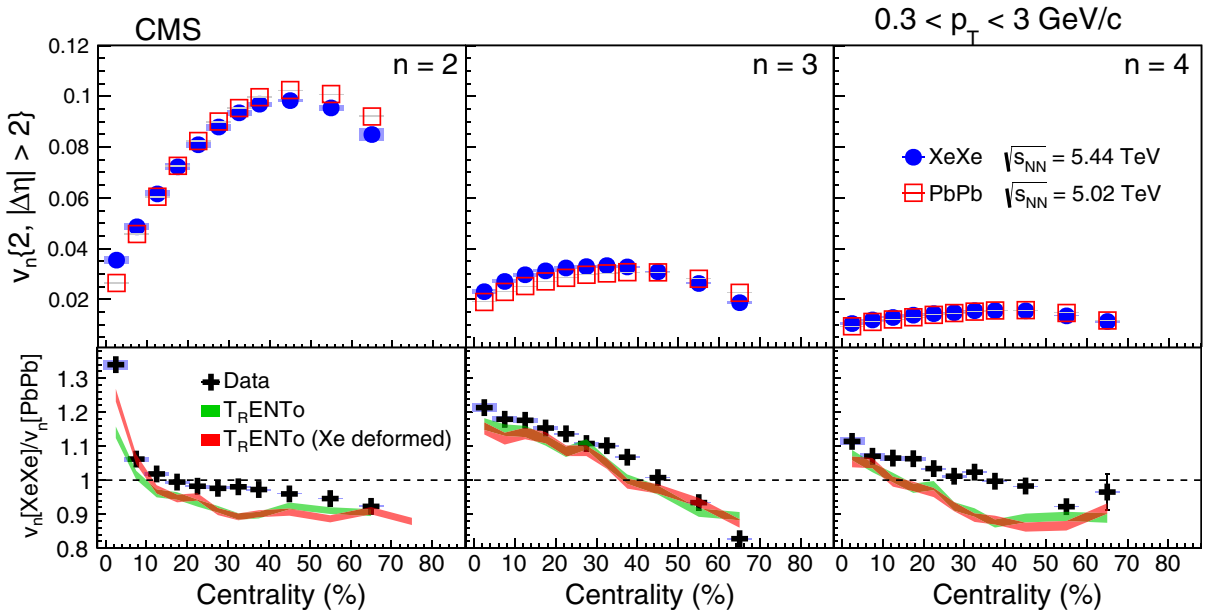

FIG. 10. Centrality dependence of the spectrum-weighted v_2, v_3, and v_4 harmonic coefficients from two-particle correlations method for $0.3 < p_T < 3.0$ GeV/c for XeXe collisions at $\sqrt{s_{NN}} = 5.44$ TeV and PbPb collisions at 5.02 TeV. The lower panels show the ratio of the results for the two systems. The bars and the shaded boxes represent statistical and systematic uncertainties, respectively. Theoretical predictions from Ref. [37] are compared to the data (shaded bands). The model calculation is done for the p_T range $0.2 < p_T < 5.0$ GeV/c.

decreases with centrality with an almost constant slope. The relativistic hydrodynamic model calculations of Ref. [37] are also shown in Fig. 10. Compared to calculations assuming a spherical Xe shape, including the xenon nuclear deformation in hydrodynamic models has little effect on the predicted flow characteristics over the centrality range 10–70%, as expected. For the most central 0–10% range, the v_2[XeXe]/v_2[PbPb] model ratio shows a greater sensitivity to the xenon nuclear deformation, with the calculation including deformation in better agreement with experiment. For all measured harmonics, the model values lie below the experimental results, with the greatest difference found for the v_4 coefficients.

VII. SUMMARY

In this paper, the v_2, v_3, and v_4 azimuthal flow harmonics are shown for xenon-xenon (XeXe) collisions at a center-of-mass energy per nucleon pair of $\sqrt{s_{NN}} = 5.44$ TeV based on data obtained with the CMS detector. Three analysis techniques with different sensitivities to flow fluctuations, including two-particle correlations, the scalar-product method, and the multiparticle cumulant method, are used to explore the event-by-event fluctuations. The harmonic coefficients are compared to those found with lead-lead (PbPb) collisions at $\sqrt{s_{NN}} = 5.02$ TeV to explore the effect of the system size. The magnitude of the v_2 coefficients for XeXe collisions are larger than those found in PbPb collisions for the most central collisions. This is attributed to a larger fluctuation component in the lighter colliding system. In more peripheral events, the PbPb v_n coefficients are consistently larger than those found for XeXe collisions. This behavior is qualitatively consistent with expectations from hydrodynamic models. A clear ordering $v_2\{2\} > v_2\{4\} \approx v_2\{6\} \approx v_2\{8\}$ is observed for XeXe collisions, with $v_2\{6\}$ and $v_2\{4\}$ values differing by 2–3%. The $v_3\{4\}/v_3\{2\}$ ratio is found to be significantly smaller than the $v_2\{4\}/v_2\{2\}$ ratio, suggesting a dominant fluctuation component for the v_3 harmonic. Hydrodynamic models that consider the xenon nuclear deformation are able to better describe the v_2[XeXe]/v_2[PbPb] ratio in central collisions than those assuming a spherical Xe shape, although the deformation appears to have little effect on the fluctuation-sensitive ratio of the cumulant orders. These measurements provide new tests of hydrodynamic models and help to constrain hydrodynamic descriptions of the nuclear collisions.

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