

1     **A hybrid macro-modelling strategy with multi-objective calibration for**  
2     **accurate simulation of multi-ring masonry arches and bridges**

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4             B. Pantò<sup>1,2\*</sup>, C. Chisari<sup>3</sup>, L. Macorini<sup>1</sup>, B.A. Izzuddin<sup>1</sup>

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6             <sup>1</sup> Department of Civil and Environmental Engineering, Imperial College London  
7             South Kensington Campus, London SW7 2AZ, United Kingdom

8             <sup>2</sup> Department of Engineering, Durham University, United Kingdom

9             <sup>3</sup> Department of Architecture and Industrial Design, University of Campania “Luigi  
10             Vanvitelli”, Abazia di S. Lorenzo, via S. Lorenzo, Aversa CE (Italy)

11  
12             **ABSTRACT**

13     This paper presents an efficient hybrid continuum-discrete macro-modelling strategy with an  
14     enhanced multiscale calibration procedure for realistic simulations of brick/block-masonry  
15     bridges. The response of these structures is affected by the intrinsic nonlinearity of the masonry  
16     material, which in turn depends upon the mechanical properties of units and mortar joints and  
17     the bond characteristics. Finite element approaches based upon homogenised representations  
18     are widely employed to assess the nonlinear behaviour up to collapse, as they are generally  
19     associated with a limited computational demand. However, such models require an accurate  
20     calibration of model material parameters to properly allow for masonry bond. According to the  
21     proposed approach, the macroscale material parameters are determined by an advanced multi-  
22     objective strategy with genetic algorithms from the results of mesoscale "virtual" tests through  
23     the minimisation of appropriate functionals of the scale transition error. The developed  
24     continuum-discrete finite element macroscale description and the calibration procedure are  
25     applied to simulate the nonlinear behaviour up to collapse of multi-ring arch-bridge specimens  
26     focusing on the 2D planar response. The results obtained are compared to those achieved using  
27     detailed mesoscale models confirming the effectiveness and accuracy of the proposed approach  
28     for realistic nonlinear simulations of masonry arch bridges.

29     **Keywords:** nonlinear analysis; continuous finite element models; mesoscale models; multi-

30 ring masonry arches; optimisation procedures; genetic algorithms.

## 31 **1 INTRODUCTION**

32 Old masonry arch bridges belong to the cultural and architectural heritage and still play a  
33 critical role within railway and roadway networks in Europe and worldwide. These structures  
34 were built following empirical rules and were not designed to resist current traffic loading and  
35 the loads induced by extreme events, such as earthquakes. An accurate assessment of the  
36 ultimate performance of these complex structural systems represents a crucial step to prevent  
37 future failures and preserve such historical structures for the next generations.

38 Masonry arch barrels are the key structural components of masonry arch bridges. Their  
39 nonlinear behaviour is strongly influenced by the mechanical properties of the two constituents,  
40 masonry units and mortar joints, and their arrangement to form the brick/blockwork of the arch  
41 (i.e. masonry bond). Two main categories of masonry arch bridges can be identified: stone  
42 masonry and brick masonry bridges [1]. In the first group, the arches are built from large  
43 voussoirs organised in a single arch ring. Conversely, in the case of brick masonry bridges, a  
44 multi-ring arrangement is usually utilised, where the number of rings depends on the span  
45 length of the arch. The rings are typically bonded together using the stretcher method, where  
46 the connection between adjoining rings is guaranteed by continuous mortar joints. To date,  
47 numerous laboratory and in-situ tests have been performed to investigate the failure  
48 mechanisms of masonry arches and bridges, considering also the influence of backfill, under  
49 monotonic and cyclic loading conditions [2]. Specific studies on multi-ring arches showed how  
50 ring separation and shear sliding generally affect the ultimate strength and failure mode [3]-  
51 [6], where weak circumferential mortar joints have been found to lead to an ultimate strength  
52 reduction of about 30% for short spans and up to 70% in the case of longer span arches.

53 In previous research, different numerical strategies have been proposed to simulate the  
54 nonlinear behaviour of masonry arches and bridges [2]. Generally, approaches based on limit

55 analysis principles can be effectively used to estimate the ultimate load capacity [7]-[9].  
56 However, such strategies do not provide information about the nonlinear response before  
57 collapse, and they are often based upon crude assumptions, e.g. the representation of masonry  
58 as a no-tension material, which may lead to underestimating the ultimate resistance of masonry  
59 arches. Previous studies also comprised simplified 2D finite element (FE) limit-analysis  
60 descriptions to simulate the arch-backfill interaction [10],[11] and 3D nonlinear FE strategies  
61 with elasto-plastic solid elements [1]-[14], where masonry is assumed as a homogeneous  
62 isotropic material disregarding its anisotropic nature. Isotropic modelling approaches are  
63 widely employed in engineering practice due to their computational efficiency, especially for  
64 the analysis of large bridges. However, they may lead to an unrealistic representations of  
65 typical failure modes not directly associated with longitudinal bending. Also, their application  
66 to masonry may require complex calibration procedures to account even in a simplified way  
67 for its anisotropic nature, as shown in [15] with reference to masonry walls. Furthermore, the  
68 use of more complex damage/plasticity orthotropic models [16]-[20], based on damage and/or  
69 plasticity formulations where tensile, shear and compressive failure mechanisms are described,  
70 are still largely applied for research and not yet considered for the practical assessment of  
71 realistic structures.

72 More recent numerical models for masonry arched structures and bridges include the micro-  
73 model strategy proposed by Milani et al. [21] using triangular rigid elements and nonlinear  
74 links, the discrete macro-element method (DMEM) [22]-[24] and the distinct element method  
75 (DEM) [25],[26]. A detailed 3D mesoscale modelling strategy for masonry arch bridges has  
76 been developed at Imperial College London [27],[28], which is used as the reference solution  
77 for the calibration of the proposed macroscale approach hereinafter. According to this strategy,  
78 the masonry parts of the bridge are simulated by using linear solid elements and 2D nonlinear  
79 interface elements to explicitly allow for the masonry bond [29]. The backfill is modelled by

80 elasto-plastic solid elements, and the connection between the masonry components and the  
81 backfill is represented through nonlinear interfaces allowing for the actual frictional  
82 interaction. This approach generally leads to accurate response predictions, including under  
83 extreme loading, but it is associated with significant computational cost which can hinder its  
84 use for the practical assessment of real large structures.

85 With the aim of achieving a suitable compromise between accuracy and efficiency, this paper  
86 proposes a hybrid continuum-discrete macroscale description for multi-ring masonry arches  
87 and masonry arches bridges. Elasto-plastic-damage continuum solid elements interacting with  
88 2D nonlinear interfaces are employed to model a masonry arch, although, unlike the mesoscale  
89 strategy, mesh discretisation is not directly related to the dimensions of units and mortar joints.  
90 The damage-plasticity model proposed in [30] and a multi-surface cohesive-frictional model  
91 [30] are employed for solid and interface elements, respectively. Furthermore, an innovative  
92 multi-objective optimisation procedure, based on virtual tests developed adopting detailed  
93 mesoscale descriptions, is put forward and applied to evaluate the mechanical parameters of  
94 the hybrid model.

95 The proposed modelling strategy with the advanced calibration procedure is validated against  
96 mesoscale simulations considering multi-ring arches and masonry arch-backfill specimens  
97 with different geometrical and mechanical properties. The numerical results confirm the  
98 accuracy and high efficiency of the developed hybrid approach, which can be used for practical  
99 and accurate assessment of realistic, including long span, masonry arch bridges.

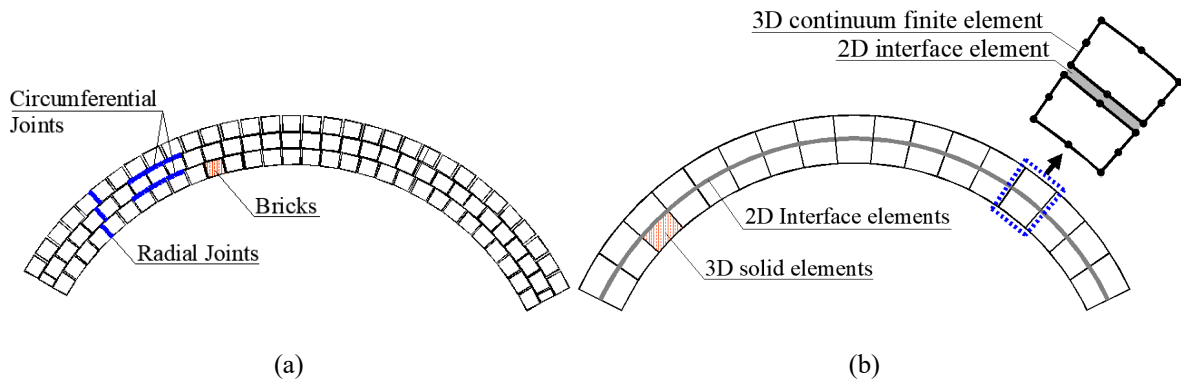
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## 101 **2 THE HYBRID MACRO-MODELLING APPROACH**

102 In the proposed FE modelling strategy, the arch is discretised using a regular mesh of nonlinear  
103 continuum 20-noded solid elements. In addition, 2D nonlinear zero-thickness interface  
104 elements are arranged along the circumferential mid-thickness surface of the arch to simulate



105 damage associated with potential ring sliding and separation. In the simplest case where only  
 106 one circumferential layer of interfaces is considered (Figure 1), each interface lumps the linear  
 107 deformability and non-linear behaviour of n-1 ring joints, with n being the number of rings of  
 108 the physical arch. Importantly, the characteristics of the FE mesh with solid elements are not  
 109 directly linked to the masonry bond. Thus, an arbitrary number of solid elements can be  
 110 employed along the length of the arch, according to the desired level of response detail, but at  
 111 least two solid elements should be arranged along the thickness of the arch to accommodate  
 112 the mid-thickness nonlinear interfaces. The accuracy due to different discretisation along the  
 113 circumferential direction is explored in the numerical applications described in the following  
 114 sections.



115  
 116  
 117 Figure 1. 2D view of (a) a generic multi-ring arch and (b) its 3D macro-modelling description.

## 118 2.1 The 3D damage-plasticity model

119 In the macroscale representation implemented in ADAPTIC [31], the isotropic plastic-damage  
 120 material model presented in [15] is used for the 20-noded solid elements. A standard  
 121 decomposition of total strains ( $\boldsymbol{\epsilon}$ ) in elastic ( $\boldsymbol{\epsilon}_e$ ) and plastic ( $\boldsymbol{\epsilon}_p$ ) components is considered, and  
 122 the stress tensor ( $\boldsymbol{\sigma}$ ) is obtained from the effective stress tensor ( $\bar{\boldsymbol{\sigma}}$ ) and a scalar damage variable  
 123  $d(\bar{\boldsymbol{\sigma}}, \kappa_t, \kappa_c)$ . The latter variable depends on the stress state and two historical variables ( $\kappa_t, \kappa_c$ )  
 124 representing the evolution of plastic strains in tension and in compression. The material  
 125 relationship is expressed analytically by:

126

$$\boldsymbol{\sigma} = (1 - d) \bar{\boldsymbol{\sigma}} = (1 - d) \mathbf{E}_0 (\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}_p) \quad (1)$$

127

128 where  $\mathbf{E}_0$  is the initial fourth-order isotropic elastic tensor.

129 The local plastic problem is solved at each integration point of the domain to evaluate the

130 effective stress, adopting a non-associated elasto-plastic constitutive law with Drucker-Prager-

131 like plastic flow potential, according to the approach proposed in [32].

132 The plastic behaviour is governed by the evolution of the yield surface:

133

$$F(\bar{\boldsymbol{\sigma}}, \boldsymbol{\kappa}) = \frac{1}{1 - \alpha} \cdot (\alpha I_1 + \sqrt{3J_2} + \beta(\boldsymbol{\kappa}) \langle \bar{\sigma}_{max} \rangle - \gamma \langle -\bar{\sigma}_{max} \rangle) + \bar{f}_c(\kappa_c) \quad (2)$$

134 where:

135     –  $\beta(\boldsymbol{\kappa}) = -\frac{\bar{f}_c(\kappa_c)}{\bar{f}_t(\kappa_t)} (1 - \alpha) - (1 + \alpha) ;$

136     –  $\alpha = \frac{\tilde{f}_{b0}-1}{2\tilde{f}_{b0}-1} ;$

137     –  $\gamma = \frac{3(1-K_c)}{2K_c-1} ;$

138     –  $\bar{\sigma}_{max} = \max(\bar{\sigma}_1, \bar{\sigma}_2, \bar{\sigma}_3)$  with  $\bar{\sigma}_i$  principal effective stress;

139     –  $\langle x \rangle = \frac{x+|x|}{2}.$

140     –  $\bar{f}_c(\kappa_c), \bar{f}_t(\kappa_t)$  effective strength in compression and tension, respectively;

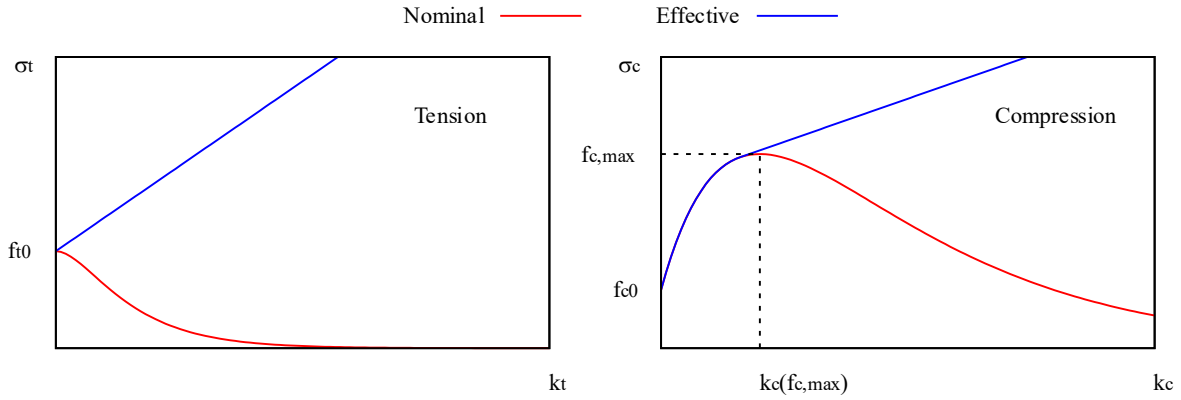
141     –  $K_c$  ratio of the second stress invariant on the tensile meridian to that on the compressive  
142 meridian at initial yield;

143     –  $\tilde{f}_{b0}$  ratio between biaxial and uniaxial compressive strength.

144 To improve the computational robustness, both tensile and compressive strengths,  $\bar{f}_\chi(\kappa_\chi)$  with

145  $\chi = t, c$ , allow for hardening behaviour, while the softening response is obtained for the

146 nominal strength  $f_{\chi}(\kappa_{\chi})$  by introducing an appropriate damage law  $d_{\chi}(\kappa_{\chi}) = 1 - \frac{f_{\chi}(\kappa_{\chi})}{\bar{f}_{\chi}(\kappa_{\chi})}$ , as  
 147 shown in Figure 2.



148

149 Figure 2. Uniaxial constitutive relationships in tension and compression.

150 The global damage variable is obtained as:

151

$$d(\bar{\boldsymbol{\sigma}}, \boldsymbol{\kappa}) = 1 - [1 - s_t(\bar{\boldsymbol{\sigma}}) d_c(\kappa_c)][1 - s_c(\bar{\boldsymbol{\sigma}}) d_t(\kappa_t)] \quad (3)$$

152 where:

153 –  $s_t(\bar{\boldsymbol{\sigma}}) = 1 - w_t r(\bar{\boldsymbol{\sigma}});$

154 –  $s_c(\bar{\boldsymbol{\sigma}}) = 1 - w_c (1 - r(\bar{\boldsymbol{\sigma}}));$

155 –  $r(\bar{\boldsymbol{\sigma}}) = \begin{cases} 0 & \text{if } \bar{\sigma}_1 = \bar{\sigma}_2 = \bar{\sigma}_3 = 0 \\ \frac{\sum_{i=1}^3 \langle \bar{\sigma}_i \rangle}{\sum_{i=1}^3 |\bar{\sigma}_i|} & \text{otherwise} \end{cases}$ , scalar parameter ranging from 0 (all principal

156 stresses are negative) to 1 (all principal stresses are positive) expressing the state stress;

157 –  $w_t, w_c$  are parameters governing the stiffness recovery from compression to tension and

158 vice versa.

159 Since as well-known softening behaviour may lead to mesh sensitivity, a fracture-energy

160 approach has been adopted to maintain objectivity in the results. In particular, the stress-strain

161 constitutive relationship is defined at element level starting from a stress-crack opening curve

162 based on fracture energy, assumed as material parameter and a characteristic length evaluated  
 163 as a function of the element volume.

164 The model has been extensively used to simulate the mechanical behaviour of concrete [32][33]  
 165 and masonry [34],[35],[36]. Some inherent model characteristics, however, hinder its use to  
 166 represent specific shear failure modes typical of multi-ring masonry arches. More specifically,  
 167 the adopted damage-plasticity continuum description does not enable the definition of the shear  
 168 strength independently from the tension and compression strengths. It can be seen by applying  
 169 Eq. (2) assuming a pure shear 2D stress state ( $\bar{\sigma}_x = \bar{\sigma}_y = \bar{\sigma}_z = \bar{\tau}_{xy} = \bar{\tau}_{yz} = 0, \bar{\tau}_{xz} = \bar{\tau}$ ) which  
 170 leads to the yield function:

171

$$F(\bar{\tau}, \boldsymbol{\kappa}) = \frac{1}{1 - \alpha} \cdot (\sqrt{3}\bar{\tau} + \beta(\boldsymbol{\kappa})\bar{\tau}) + \bar{f}_c(\kappa_c) \quad (4)$$

172

173 Imposing  $F(\bar{\tau}, \boldsymbol{\kappa}) = 0$  in Eq. (4), the effective shear strength  $\bar{f}_v(\boldsymbol{\kappa})$  can be evaluated as:

174

$$\bar{f}_v(\boldsymbol{\kappa}) = \frac{1 - \alpha}{\sqrt{3} + \beta(\boldsymbol{\kappa})} |\bar{f}_c(\kappa_c)| \quad (5)$$

175

176 and, since  $r(\bar{\boldsymbol{\sigma}}) = 0.5$  for pure shear, the damage parameter becomes:

177

$$d(\boldsymbol{\kappa}) = 1 - [1 - (1 - 0.5 w_t) d_c(\kappa_c)][1 - (1 - 0.5 w_c) d_t(\kappa_t)] \quad (6)$$

178

179 Assuming that damage in compression has not developed,  $d_c(\kappa_c) = 0$ , the equivalent damage  
 180 parameter becomes:

181

$$d(\boldsymbol{\kappa}) = (1 - 0.5 w_c) d_t(\kappa_t) \quad (7)$$

182 and

$$\begin{aligned} f_v(\boldsymbol{\kappa}) &= (1 - d(\boldsymbol{\kappa})) \bar{f}_v(\boldsymbol{\kappa}) \\ &= [1 - (1 - 0.5 w_c) d_t(\kappa_t)] \frac{1 - \alpha}{\sqrt{3} + \beta(\boldsymbol{\kappa})} |\bar{f}_c(\kappa_c)| \end{aligned} \quad (8)$$

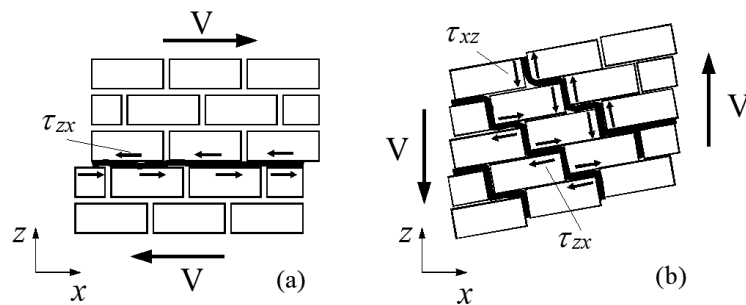
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184 From Eq. (8), some considerations can be made on the shear behaviour of the model. Firstly,  
 185 the initial shear strength (Eq. (5) with  $\boldsymbol{\kappa} = \mathbf{0}$ ) is governed by the initial compression and tension  
 186 strengths and the parameter  $\tilde{f}_{b0}$ , which in practice is always in the range 1.12÷1.16. This  
 187 confirms that it is not possible to define a specific shear strength independent from tension and  
 188 compression strengths, as for instance a shear strength relating to the sliding of mortar joints,  
 189 which is a typical shear failure mode for multi-ring masonry arches. A workaround to have  
 190 some freedom in the definition of initial shear strength could be to calibrate  $\bar{f}_c(0) = f_{c0}$   
 191 appropriately and independently from the observed compressive behaviour, while both  $f_{t0}$  and  
 192  $f_{c,max}$  would still be determined based on their specific failure modes.

193 The second consideration is that the evolution of nominal shear strength depends on the  
 194 parameter  $w_c$  (see Eq. (8)) which is defined based on the expected cyclic response (stiffness  
 195 recovery). A typical value, leading to complete stiffness recovery from tension to compression,  
 196 is  $w_c = 1.0$  [37],[38]. Inserting this in Eq. (7), the expression for damage in pure shear in the  
 197 absence of compression damage is obtained  $d(\kappa_t) = 0.5 d_t(\kappa_t)$ . The conclusion is that the  
 198 evolution of nominal shear strength is completely governed by damage in tension, without the  
 199 possibility for specifying an alternative more realistic constitutive relationship.

200 Finally, it is worth mentioning that the macroscale damage-plasticity continuum representation  
 201 is not capable of distinguishing failure due to shear parallel to the mortar bed joint  $\tau_{zx}$  from  
 202 that due to shear orthogonal to the mortar bed joint  $\tau_{xz}$ , as in the Cauchy solid these two stresses

203 are equal, and the yield surface cannot consider separate contributions. In reality, while the  
 204 former failure mode is governed by sliding of the units on the weak planes represented by the  
 205 mortar joints, the latter is governed by the internal rotation of bricks depending on their  
 206 geometric shape ratio and brick interlocking, as schematically shown in [39]. To allow for these  
 207 different phenomena enriched continuum representations, e.g. Cosserat continuum [39], would  
 208 need to be employed.



209

210 Figure 3. Shear stress and failure mode in a brick/block masonry sample under (a) pure shear parallel  
 211 and (b) orthogonal to the bed joints.

212 To overcome these intrinsic limitations of typical continuum damage-plasticity constitutive  
 213 models, an alternative hybrid macroscale representation is proposed, in which, as outlined  
 214 before, shear sliding along the continuous circumferential mortar joints of multi-ring arches is  
 215 described by introducing nonlinear interfaces whose material characteristics are defined based  
 216 on the calibration strategy described in Section 3.

217 **2.2 Constitutive model for nonlinear interface elements**

218 2D 16-noded interface elements [29] are employed for the mid-thickness circumferential  
 219 interfaces using the plasticity-damage constitutive model proposed in [30]. According to this  
 220 description, interface tractions and relative displacements describing the static and kinematics  
 221 of the element, are composed of a normal component in the direction orthogonal to the interface  
 222 and two shear components on the plane of the interface. The effective stresses are evaluated at

223 each Gauss point by solving a linear hardening elasto-plastic problem considering multi-  
 224 surface plasticity. Then, the nominal stresses are obtained by multiplying the effective stresses  
 225 by the damage matrix  $\mathbf{D}$ , containing the damage index in tension, shear and compression  
 226 ranging from 0 (no-damage) to 1 (complete damage).

227 Similarly to the solid elements, a standard decomposition between elastic and plastic  
 228 deformations is considered and the concept of effective stress  $\bar{\mathbf{s}} = \mathbf{K}_0(\mathbf{e} - \mathbf{e}_p)$  is introduced,  
 229 where  $\mathbf{K}_0 = \text{diag}\{k_n \quad k_t \quad k_t\}$  is the diagonal initial stiffness matrix with  $k_n$  and  $k_t$  the  
 230 normal and shear stiffness,  $\bar{\mathbf{s}} = [\bar{\sigma} \quad \bar{\tau}_1 \quad \bar{\tau}_2]$ ,  $\mathbf{e} = [\varepsilon \quad \gamma_1 \quad \gamma_2]$  and  $\mathbf{e}_p = [\varepsilon_p \quad \gamma_{p1} \quad \gamma_{p2}]$  the  
 231 effective stress, the total strains and the plastic strains, respectively. The nominal stresses are  
 232 evaluated from the effective stress according to:

$$233 \quad \mathbf{s} = (\mathbf{I}_3 - \mathbf{D})\bar{\mathbf{s}} = (\mathbf{I}_3 - \mathbf{D})\mathbf{K}_0(\mathbf{e} - \mathbf{e}_p) \quad (9)$$

234 where  $\mathbf{D}$  represents an anisotropic damage tensor, containing distinct variables for the normal  
 235 ( $D_n$ ) and the tangential ( $D_t$ ) directions. A tri-linear plastic yield domain is considered to  
 236 simulate the tensile (Mode I), shear (Mode II) and crushing (Model III) failure models. Three  
 237 distinct plastic works, corresponding to each fracture mode rule the evolution of the damage  
 238 variables.

239 The plastic yield domain (Figure 4) is composed of three surfaces,  $F_t$ ,  $F_c$ , and  $F_s$  respectively,  
 240 associated with the tensile (mode I), compression and shear (mode II) failure modes, as defined  
 241 by:

$$243 \quad F_s(\bar{\mathbf{s}}, q) = \sqrt{\bar{\tau}_1^2 + \bar{\tau}_2^2} + \bar{\sigma} \tan(\phi) - c' \quad (10a)$$

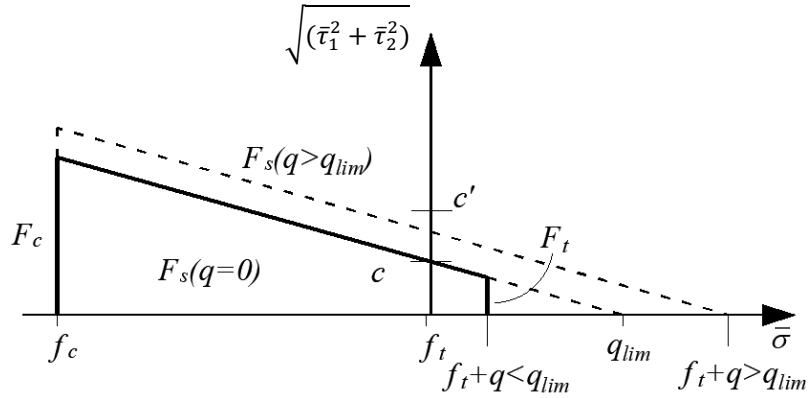
$$244 \quad F_t(\bar{\mathbf{s}}, q) = \bar{\sigma} - (f_t - q) \quad (10b)$$

$$245 \quad F_c(\bar{\mathbf{s}}) = -\bar{\sigma} + f_c \quad (10c)$$

246

247 where  $f_t$  and  $f_c$  are the tensile and compression material strengths,  $\phi$  the friction angle and  $q$  a  
 248 linear hardening variable, ranging from 0 (initial value) to the limit value  $q_{lim} = \frac{c}{\tan(\phi)} - f_t$ .  
 249 Moreover,  $c' = c$  if  $q \leq q_{lim}$  and  $c' = c + (q - q_{lim})\tan(\phi)$  if  $q > q_{lim}$ . With the increase  
 250 of  $q$  the surface  $F_t$  reduces until becoming a point when  $q$  reaches the value  $q_{lim}$ . On the other  
 251 hand,  $F_s$  increases with the increase of  $q$  (Figure 4). Two associated plastic flows are related to  
 252  $F_t$  and  $F_c$ , while a plastic potential  $G_s$  in shear, obtained from  $F_s$  substituting  $\phi$  to  $\phi_g$ , is  
 253 considered to take into account the effects of masonry dilatancy.

254



255

256

Figure 4. Yield surface of the material model for nonlinear interfaces.

257 Following the solution of the plastic problem, the damage evolution is evaluated as a function  
 258 of the three ratios  $r_c = W_{pc}/G_c$ ,  $r_t = W_{pt}/G_t$  and  $r_s = W_{ps}/G_s$  where  $W_{pc}$ ,  $W_{pt}$  and  $W_{ps}$  are  
 259 the plastic works in compression, tension and shear, respectively, and  $G_c$ ,  $G_t$ ,  $G_s$  the  
 260 corresponding fracture energies. Finally, the nominal stresses are given by Eq. (8). More details  
 261 on the model formulation can be found in [30].

262

### 263 3 CALIBRATION PROCEDURE

264 The mechanical calibration of the proposed model requires the determination of several  
 265 material parameters defining the linear and nonlinear behaviour of the 3D solid elements and



266 the 2D interfaces, as described Sections 2.1 and 2.2, and reported in [15],[29] and [30]. For this  
 267 reason, an objective and robust calibration procedure represents a fundamental step to  
 268 guarantee the model accuracy and applicability.

269 This work presents an original multi-objective calibration procedure based on the multiscale  
 270 approach proposed in [15] which considers the representation of a structure under suitable  
 271 boundary conditions according to two scales: mesoscale, indicated hereinafter by the  
 272 superscript  $m$ , and macroscale, with the superscript  $M$ . The considered setup is called *virtual*  
 273 *test*, and it is assumed there exists a mapping  $\mathcal{M}: \Omega^m \rightarrow \Omega^M$  between the mesoscale and the  
 274 macroscale domains. As elaborated subsequently, different to the procedure proposed in [15]  
 275 which consisted of a single-objective optimisation algorithm, the newly proposed procedure  
 276 leads to a multi-objective optimisation problem allowing for a set of optimum solutions (Pareto  
 277 Front) which improves the robustness and accuracy of the model calibration procedure.

278 According to the original formulation put forward in [15], stress power equivalence between  
 279 the two scales is approximately enforced on the entire domain of the virtual test. The stress  
 280 power equivalence reads:

$$\int_{\Omega^M} \boldsymbol{\sigma}^M : \dot{\boldsymbol{\epsilon}}^M d\Omega^M = \int_{\Omega^m} \boldsymbol{\sigma}^m : \dot{\boldsymbol{\epsilon}}^m d\Omega^m + \dot{\epsilon} \quad (11)$$

281 where  $\dot{\epsilon}$  represents the error rate due to the approximations induced by the specific macromodel  
 282 utilised. Considering pseudo-static stress states, the equality between internal and external  
 283 work implies:

$$\begin{aligned} \int_{\Gamma^M} \mathbf{t}^M \cdot \dot{\mathbf{u}}^M d\Gamma^M + \int_{\Omega^M} \mathbf{b}^M \cdot \dot{\mathbf{u}}^M d\Omega^M \\ = \int_{\Gamma^m} \mathbf{t}^m \cdot \dot{\mathbf{u}}^m d\Gamma^m + \int_{\Omega^m} \mathbf{b}^m \cdot \dot{\mathbf{u}}^m d\Omega^m + \dot{\epsilon} \end{aligned} \quad (12)$$

284

285 where  $\mathbf{t}$  are the surface forces on the boundary  $\Gamma$ , while  $\mathbf{b}$  are volume forces. Neglecting the  
 286 contribution of these latter and considering the chain rule of differentiation, Eq. (12) finally  
 287 reads:

$$\dot{\epsilon} = \int_{\Gamma^M} \left( \mathbf{t}^M \cdot \dot{\mathbf{u}}^M - \mathbf{t}^m \cdot \dot{\mathbf{u}}^m \frac{\partial \Gamma_i^m}{\partial \Gamma_i^M} \right) d\Gamma_i^M \quad (13)$$

288 Eq. (13) represents the error rate at time  $t$  due to the scale transition. In [15], a global non-  
 289 negative monotonically increasing error function was defined:

$$\begin{aligned} \epsilon(t) &= \int_0^t [\dot{\epsilon}(\tau)]^2 d\tau \\ &= \int_0^t \left[ \int_{\Gamma^M} \left( \mathbf{t}^M(\tau) \cdot \dot{\mathbf{u}}^M(\tau) - \mathbf{t}^m(\tau) \cdot \dot{\mathbf{u}}^m(\tau) \frac{\partial \Gamma_i^m}{\partial \Gamma_i^M} \right) d\Gamma_i^M \right]^2 d\tau \end{aligned} \quad (14)$$

290

291 The extension of the original procedure, proposed in this paper, consists of partitioning the  
 292 error defined as in Eq. (11) or in Eq. (13) as:

$$\begin{aligned}
\dot{\epsilon} &= \dot{\epsilon}_1 + \dot{\epsilon}_2 + \dots \\
&= \int_{\Omega_1^M} (\boldsymbol{\sigma}^M : \dot{\boldsymbol{\epsilon}}^M - \boldsymbol{\sigma}^m : \dot{\boldsymbol{\epsilon}}^m) d\Omega_1^M \\
&\quad + \int_{\Omega_2^M} (\boldsymbol{\sigma}^M : \dot{\boldsymbol{\epsilon}}^M - \boldsymbol{\sigma}^m : \dot{\boldsymbol{\epsilon}}^m) d\Omega_2^M + \dots
\end{aligned} \tag{15}$$

$$\begin{aligned}
\dot{\epsilon} &= \dot{\epsilon}_1 + \dot{\epsilon}_2 + \dots \\
&= \int_{\Gamma_1^M} (\mathbf{t}^M \cdot \dot{\mathbf{u}}^M - \mathbf{t}^m \cdot \dot{\mathbf{u}}^m) d\Gamma_1^M \\
&\quad + \int_{\Gamma_2^M} (\mathbf{t}^M \cdot \dot{\mathbf{u}}^M - \mathbf{t}^m \cdot \dot{\mathbf{u}}^m) d\Gamma_2^M + \dots
\end{aligned} \tag{16}$$

293 The contributions  $\dot{\epsilon}_1, \dot{\epsilon}_2, \dots$  respectively refer to a volume partitioning in Eq. (15), with  
294  $\Omega_1^{M|m} + \Omega_2^{M|m} + \dots = \Omega^{M|m}$ , or load-based partitioning in Eq. (16). For the sake of simplicity,  
295 in Eq. (15), (16) it is assumed that there is not any modification of volumes and surfaces in the  
296 scale transition, i.e.,  $\partial\Gamma_i^m / \partial\Gamma_i^M = \partial\Omega_i^m / \partial\Omega_i^M = 1$ .

297 In this case several error functions can be defined as:

298

$$\omega_i = \int_0^T [\dot{\epsilon}_i(\tau)]^2 d\tau \quad i = 1, 2, \dots \tag{17}$$

299

300 The solution of the calibration procedure is given by the solution of the multi-objective  
301 minimisation problem:

302

$$\tilde{\mathbf{p}} = \arg \min_p [\omega_1, \omega_2, \dots] \tag{18}$$

303 The error partitioning defined in Eq. (15) or (16) has two consequences. The first consequence  
304 is that it allows defining the granularity of the homogenisation, avoiding the possible error  
305 compensations given by different parts of the structure. For instance, if the nonlinearities in the  
306 mesoscale model are concentrated in one small region of the domain, it is possible to use  
307 volume partitioning in Eq. (15) to focus the calibration of the parameters governing the  
308 nonlinear behaviour of the macroscale representation in that region, while controlling the  
309 elastic parameters by matching the response in the remaining domain. The second consequence  
310 is that the calibration problem is turned into a multi-objective optimisation problem, in contrast  
311 to the original formulation [15] which was a single-objective optimisation procedure. As shown  
312 in [40] and [41], using multiple objectives in a calibration problem may strongly increase the  
313 robustness of the procedure. In the numerical applications reported in Sections 4 and 5, two  
314 partitions of the global error are considered to simplify the interpretation of the optimisation  
315 results. However, the use of a larger number of partitions may be considered.

316 The multi-objective optimisation problem is solved by means of a Non-Dominated Sorting  
317 Genetic Algorithm [42], implemented in TOSCA-TS software [43]. The optimum is given by  
318 the Pareto Front (PF), which represents the set of non-dominated solutions. A careful  
319 investigation on the features of the Pareto Front may highlight possible inconsistencies of the  
320 model to calibrate [40] and represents a key part of the calibration procedure towards the  
321 definition of the most representative solutions and a significant improvement of the original  
322 procedure presented in [15].

323 Finally, it is worth noting that the calibration strategy considers the evolution of the stress  
324 power over time, and thus it cannot properly allow for the additional work contribution due to  
325 initial loading. Thus, it is preferable to avoid initial loading in the virtual test. However, this  
326 does not limit the applicability of the procedure, as multiple loads with independent loading  
327 paths can be introduced without any modifications in the methodology.

328 The proposed strategy enables an objective evaluation of the macroscale model parameters  
329 given the masonry mesoscale properties which can be obtained directly from simple in-situ or  
330 laboratory tests performed on masonry units and mortar (or tests performed on small  
331 assemblages of units such as triplets), following consolidated methodologies already reported  
332 in the literature [44].

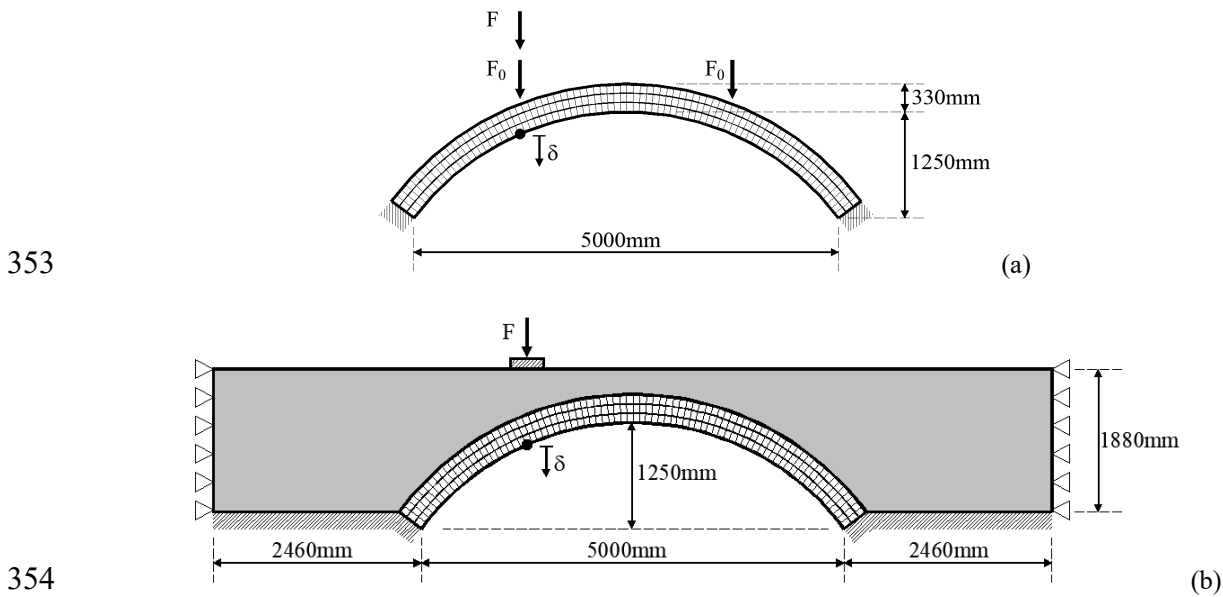
## 333 **4 NUMERICAL SIMULATIONS OF MEDIUM SPAN MASONRY ARCHES AND** 334 **BRIDGES**

335 Two medium-span masonry arch specimens, one interacting with backfill as found in typical  
336 masonry arch bridges, are analysed by means of detailed mesoscale models [29]. The results  
337 of the mesoscale analyses are then used as reference solutions to highlight some limits of a  
338 typical continuum macroscale description, and to investigate the improved accuracy  
339 guaranteed by the proposed continuum-discrete hybrid representation for multi-ring masonry  
340 arches.

### 341 **4.1 Masonry arch and bridge specimens**

342 The first specimen (Figure 5a) consists of a 5 m span three-ring brick-masonry arch. The arch  
343 is characterised by 1250 mm rise, 330 mm thickness and 675 mm width. Adjacent rings with  
344  $215 \times 102.5 \times 65$  mm<sup>3</sup> bricks are connected according to the stretcher method by continuous  
345 circumferential 10 mm thick mortar joints. The second specimen (confined arch) comprises a  
346 brick-masonry arch with the same geometrical characteristics of the bare arch interacting with  
347 backfill material, which, extends 2460 mm horizontally from the two supports of the arch and  
348 300 mm vertically above the crown, according to the experimental layout considered in [4],  
349 (Figure 5b). Full supports are assumed at the base of the arch and the backfill, while simple  
350 supports against the horizontal longitudinal displacements are applied on the two vertical sides

351 of the backfill. Moreover, the horizontal transverse displacements on the two lateral faces of  
352 the arch and backfill are restrained to represent a plane strain condition (Figure 5b).



354

355 Figure 5. Geometrical characteristics and loading conditions for the (a) bare and (b) confined arch  
356 specimens.

## 357 4.2 Mesoscale simulations

358 In the numerical mesoscale description, 20-noded elastic solid elements are used to simulate  
359 the brick units and 16-noded interfaces [29] are employed to represent both the radial and the  
360 circumferential mortar joints. As the focus is on the 2D response, a mesh with only one element  
361 along the representative 1m width of the arch specimens is considered. The mesoscale  
362 description of the arch requires 240 3D solid elements, 403 2D interface elements and 6453  
363 nodes to which correspond 19359 DOFs. The backfill is modelled adopting a FE mesh with  
364 15-noded tetrahedral elements. Finally, nonlinear interface elements are utilised to model the  
365 physical interface connecting the arch to the backfill.

366 Two masonry types have been considered in the analyses: a *strong* masonry to represent  
367 modern good quality brickwork, and a *weak* masonry to represent historical masonry [45]. The  
368 mesoscale mechanical parameters are reported in Tables 1 and 2. These parameters have been

369 selected based on previous numerical studies where the adopted mesoscale description for  
370 masonry arches and bridges was validated against experimental tests. In particular, the  
371 parameters of the strong masonry have been adopted in [27] and [28] to reproduce the response  
372 of a two-ring 3 m span arch, while the parameters for weak masonry were adopted in [47] in  
373 further validations against physical experiments.

374 Following [28], an elasto-plastic material model with a modified Drucker-Prager yield criterion  
375 is employed for the backfill, assuming a Young's modulus  $E_b = 500MPa$ , a cohesion  $c_b =$   
376  $0.001MPa$ , a friction and a dilatancy coefficient  $\tan\phi_b = 0.95$  and  $\tan\psi_b = 0.45$ . The  
377 nonlinear interfaces simulating the interaction between the arch and the backfill at the extrados  
378 of the arch have tensile strength  $f_{fi} = 0.002MPa$ , cohesion  $c_{fi} = 0.0029MPa$ , friction  
379 coefficient  $\tan\phi_{fi} = 0.6$  and zero dilatancy.

380 Table 1. Mechanical parameters of the bricks adopted in the analyses.

Masonry	$E_b$	$\nu$	$w$
	[MPa]	[-]	[kN/m <sup>3</sup> ]
Weak	6000	0.15	16
Strong	16000	0.15	22

381

382 Table 2. Interface mechanical parameters adopted in the analyses.

Masonry	$k_n - k_t$	$f_t - f_c - c$	$G_t - G_s - G_c$	$tg\phi - tg\phi_g$
	[N/mm <sup>3</sup> ]	[MPa]	[N/mm]	
Weak	60.0 - 30.0	0.05 - 9.1 - 0.085	0.02 - 0.125 - 5.0	0.5 - 0.0
Strong	90.0 - 40.0	0.26 - 24.5 - 0.40	0.12 - 0.125 - 5.0	0.5 - 0.0

383

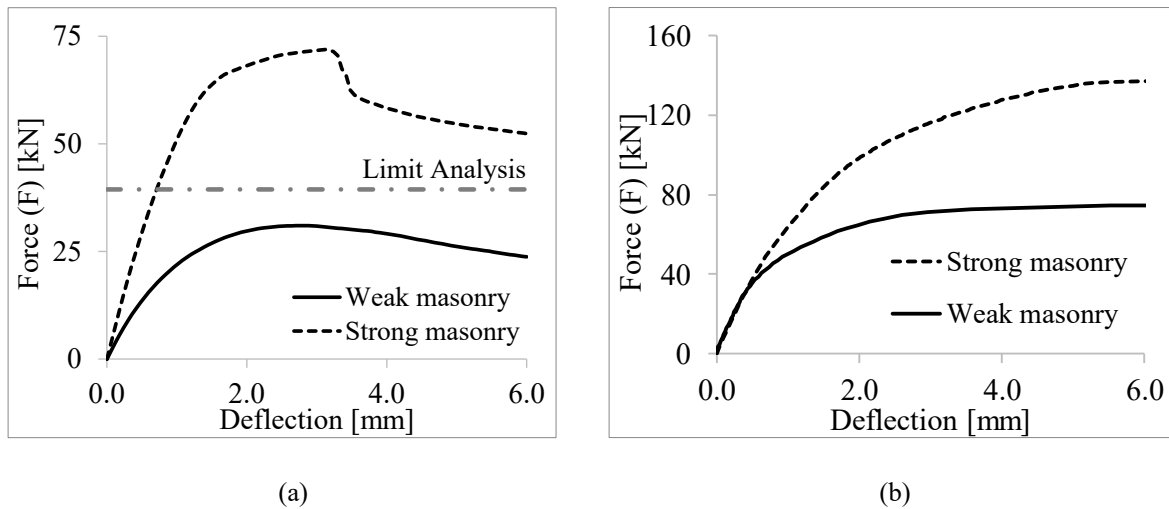
384 In the numerical simulations of the bare arch, two initial vertical forces  $F_0 = 22.5kN$  are applied  
385 at the quarter and three-quarter span and then maintained constant during the subsequent

386 loading stage, when a vertical force  $F$  is applied at quarter span and monotonically increased  
387 up to collapse. Both forces  $F_0$  and  $F$  are uniformly distributed on a patch area of  $210 \times 675 \text{mm}^2$ .  
388 When the arch interacts with the backfill, the initial load corresponds to the weight of the arch  
389 and the backfill both with a specific weight of  $22 \text{kN/m}^3$ , while the force  $F$  is applied on the top  
390 surface of the backfill on a patch area of  $400 \times 675 \text{mm}^2$  centred at the quarter span of the arch  
391 (Figure 5b). To improve the numerical stability, nonlinear dynamic analysis is performed by  
392 imposing an initial velocity of  $0.1 \text{mm/s}$  at the loaded nodes, which is maintained constant  
393 during the simulation up to collapse. Zero viscous damping is considered in the analyses.

394 Figure 6 shows the load-displacement curves of the bare arch (Figure 6a) and the arch  
395 interacting with backfill (Figure 6b), where the force  $F$  is plotted against the vertical deflection  
396 at the quarter span of the arch. In Figure 6a, the ultimate load evaluated through the limit-  
397 analysis, based on the classic Heyman's hypotheses and evaluated through an ad-hoc tool [48],  
398 is reported for comparison.

399 A significant influence of the masonry typology on the global response is observed both in the  
400 case of the bare arch, where the ultimate load ranges from  $31 \text{kN}$  to  $66 \text{kN}$ , and for the arch with  
401 backfill, where the peak force varies from  $74 \text{kN}$  to  $137 \text{kN}$ . As expected, the initial stiffness of  
402 the bare arch is significantly affected by the masonry characteristics. Conversely, the confined  
403 arch shows almost the same initial stiffness for weak and strong masonry. Considering the  
404 specific weight of the strong masonry (Table 1), standard limit analysis provides a prediction  
405 of the peak-load ( $39.43 \text{kN}$ ) significantly lower compared to the strong-masonry model due to  
406 the hypothesis of no-tension material. At the same time, it provides an overestimated peak-load  
407 compared to the weak-masonry model because it neglects the sliding between the rings.





408 Figure 6. Load-deflection curves for the (a) bare arch and (a) the arch interacting with backfill.

409 The failure mechanisms and the equivalent von-Mises stress contours of the two specimens,  
 410 with both masonry typologies, are reported in Figure 7. Finally, Figure 8 shows the tensile  
 411 damage contours at the last step of the analysis obtained by the different models. The failure  
 412 mechanism of the models with weak masonry is characterised by shear sliding along the  
 413 circumferential interfaces, mainly concentrated in the zone between the left support of the arch  
 414 and the loading area at quarter span, and close to the three-quarter span of the arch. This  
 415 mechanism prevents the activation of flexural plastic hinges. In the case of strong masonry,  
 416 flexural failure is observed with the opening of four radial cracks in both the bare arch (Figure  
 417 7) and the arch with backfill (Figure 7d). In the models with weak masonry, significant damage  
 418 in the radial joints is observed close to the load. In the models with backfill, large portions at  
 419 the extrados of the arch are affected by shear-sliding damage, both at the radial and  
 420 circumferential interfaces. This damage develops also at the frame-backfill interfaces.

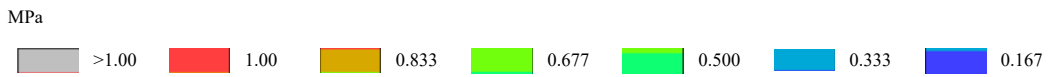
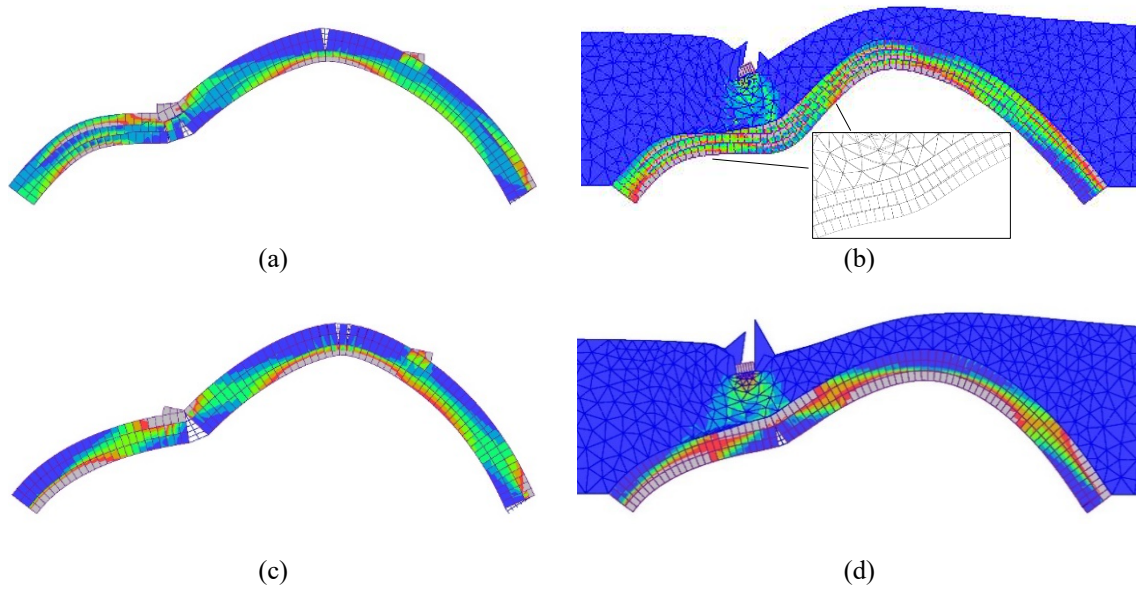
421 In the following, the mesoscale solutions are assumed as the baseline results for the assessment  
 422 of more efficient macroscale models and the proposed hybrid continuous-discrete descriptions  
 423 for multi-ring arches.

424

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428

Figure 7. Ultimate deformed shape and von-Mises stress contours for the bare arch with (a) weak

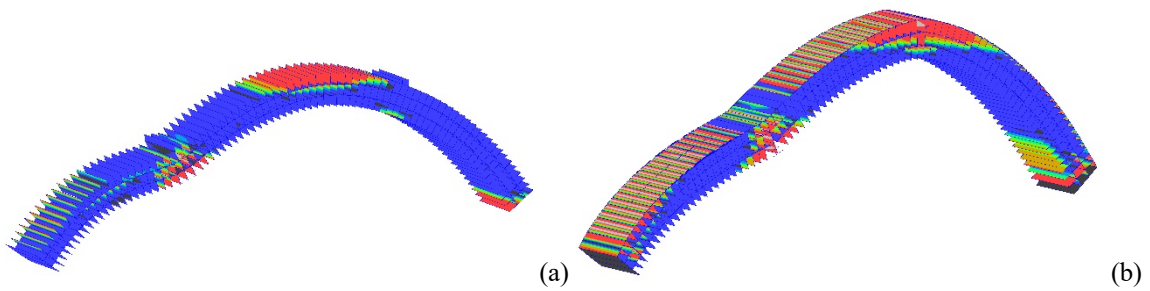
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masonry and (c) strong masonry, and for the arch interacting with backfill with (b) weak masonry and

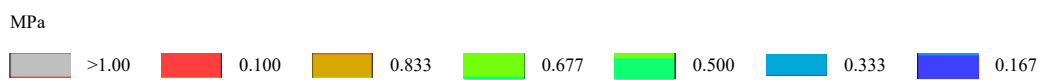
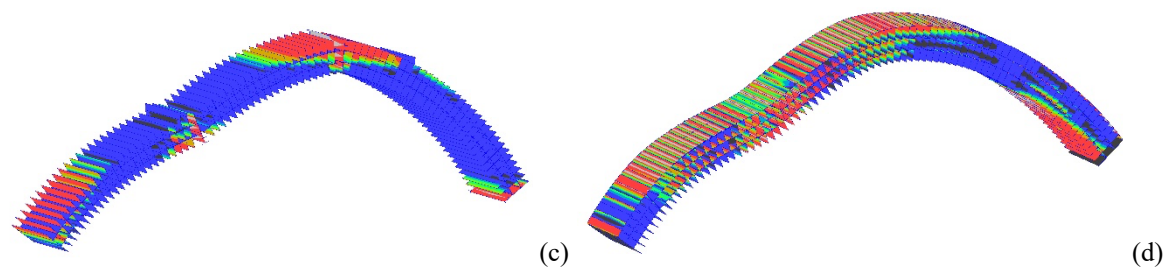
430

(d) strong masonry.

431



432



433

Figure 8. Interface tensile damage contours for the bare arch with (a) weak masonry and (c) strong

434

masonry, and for the arch interacting with backfill with (b) weak masonry and (d) strong masonry.

### 435 4.3 Macroscale simulations

436 The two specimens presented in Section 4.1 have been analysed by a continuum macroscale  
437 description for masonry, using the isotropic damage-plasticity constitutive law described  
438 before. Since the model is developed employing quadratic elements, a relatively coarse mesh  
439 can be used to improve computational efficiency. More specifically, a mesh with two elements  
440 along the thickness of the arch with a length in the circumferential direction approximately  
441 equal to half the thickness of the arch has been considered in the numerical simulations.  
442 Moreover, as for the mesoscale model, only one solid element is arranged along the 1m width  
443 of the arch. As a result of this, the masonry arch is represented by 80 3D solid elements, 40 2D  
444 interface elements and 981 nodes corresponding to 2943 DOFs. It can be observed that the  
445 macro-modelling description allows a reduction of 85% of DOFs compared to the mesoscale  
446 description demonstrating the potential for considerable reduction in computational demands  
447 with the proposed model.

448 The aim of this investigation is to explore the accuracy and potential limitations of a standard  
449 continuum isotropic macroscale approach to predict the response of multi-ring masonry arches,  
450 where the model material parameters are calibrated according to two alternative *simplified* and  
451 *advanced* procedures.

#### 452 4.3.1. Simplified calibration procedure

453 In initial macroscale simulations, the material model parameters for masonry have been  
454 evaluated through a simplified calibration procedure considering the mesoscale material  
455 properties reported in Tables 1 and 2. The macroscopic Young's modulus for the masonry  
456 material  $E$  has been determined by combining in series the stiffness of brick units with that of  
457 the mortar interfaces along the direction of the arch. The tensile strength and fracture energy  
458  $f_t, G_t$  and the compressive strength  $f_c$  are assumed coincident to the corresponding values of

459 the mesoscale interfaces. The remaining parameters for the damage plasticity model are  
460 assumed equal to standard values used in previous studies for modelling masonry materials.

461 More specifically:

- 462 • The ratio between initial and maximum compressive strength  $\tilde{f}_y = \frac{f_{c0}}{f_{c,max}}$  is assumed  
463 equal to 0.3 according to [32][33]**Error! Reference source not found.;**
- 464 • The dilatancy angle  $\psi$  is taken equal to  $35^\circ$  which is consistent with the value adopted  
465 for modelling quasi-brittle material as concrete [32][33] and corresponds  
466 approximately to the median of the values (ranging from  $10^\circ$  to  $50^\circ$ ) typically used for  
467 masonry [34][35][36];
- 468 • The eccentricity of the plastic flow potential is taken as  $\epsilon = 0.1$  to improve  
469 computational robustness as suggested in [37];
- 470 •  $\mu$  governing the relative influence of damage and plasticity in tension ( $\mu = 0$  for fully  
471 damage material) is assumed equal to 0.2;
- 472 • The plastic strain at maximum compression stress  $k_{c,f_c}$  is taken as 0.002 following [47];
- 473 • The ratio between the plastic strain at damage onset in compression and the plastic  
474 strain at maximum compression  $\rho_c$  is considered equal to 1.0, as damage is assumed to  
475 develop in the softening branch of the stress-strain response.

476 As noted in Section 2.1, preliminary numerical simulations showed a significant influence of  
477 the parameter governing the stiffness recovery in compression  $w_c$  on the global response of the  
478 arch. Thus, two limit values (0,1) are considered, while parameter  $w_t$  determining the stiffness  
479 recovery in tension is assumed as zero. The complete set of mechanical properties for the  
480 continuum macro-modelling description are reported in Table 3.

481

482

483

Table 3. Macroscopic mechanical parameters resulting from the simplified calibration procedure.

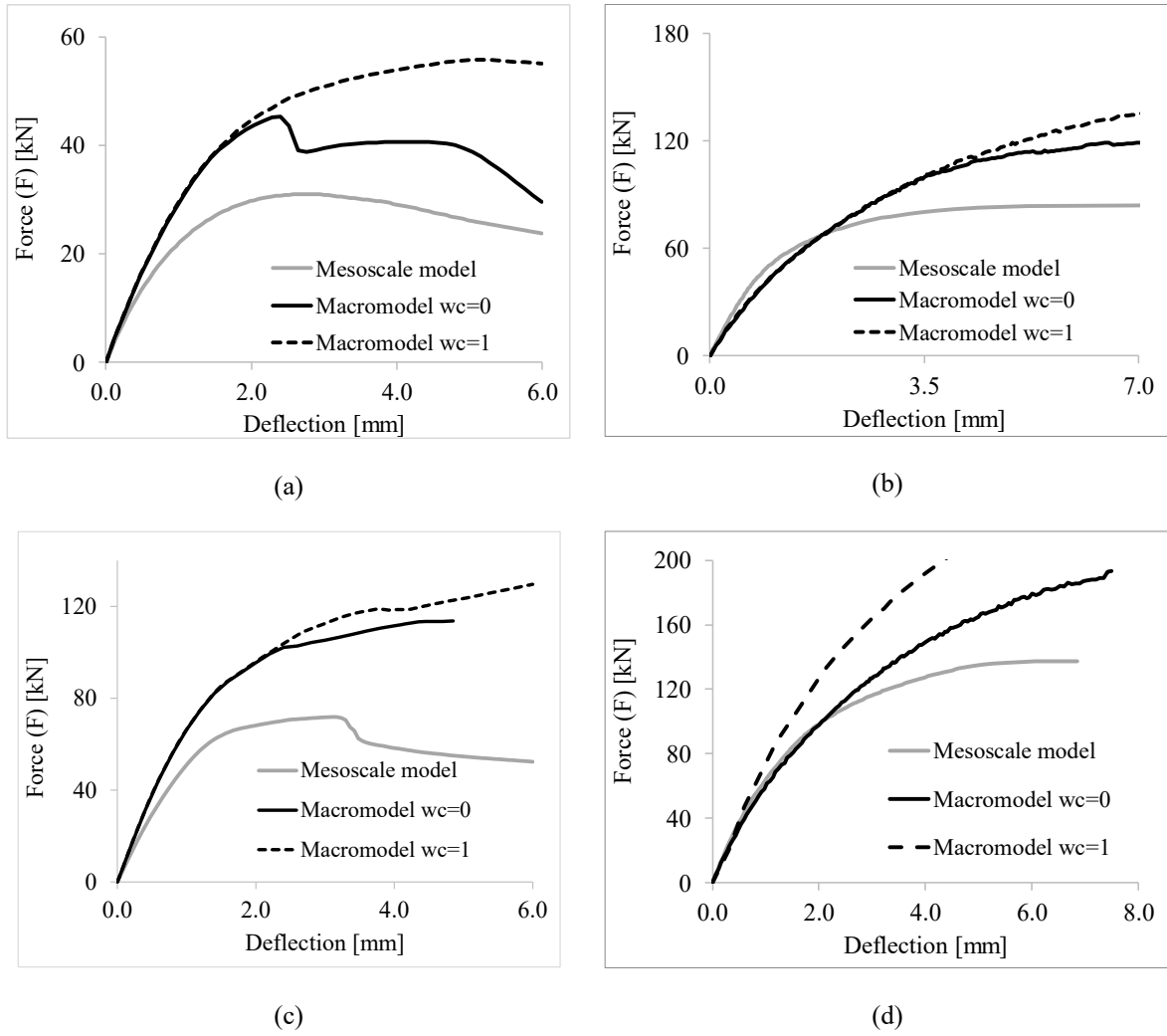
$E$	$\nu$	$\tilde{f}_{b0}$	$\tilde{f}_y$	$\psi$	$\epsilon$	$K_c$	$f_{mt}$	$f_{mc}$	$G_{mt}$	$\mu$	$k_{c,fc}$	$\rho_c$	$w_c$	$w_t$
[MPa]	[-]	[-]	[-]	[°]	[-]	[-]	[MPa]	[MPa]	[N/mm]	[-]	[-]	[-]	[-]	[-]
weak masonry														
2571	0.15	1.16	0.3	35	0.1	0.66	0.05	9.1	0.02	0.2	2E-3	1.0	0.0	0.0
													1.0	
strong masonry														
4747	0.15	1.16	0.3	35	0.1	0.66	0.26	24.0	0.12	0.2	2E-3	1.0	0.0	0.0
													1.0	

484

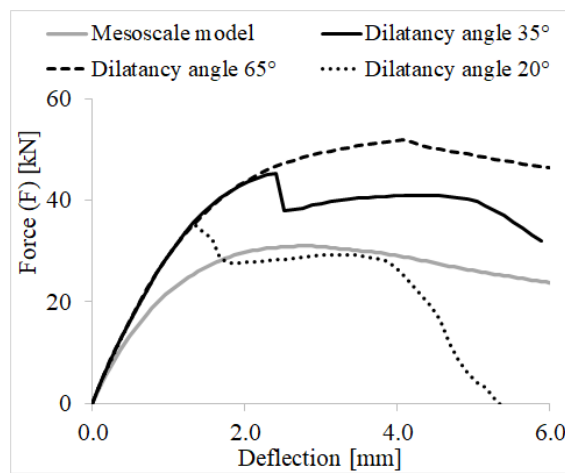
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486 Figure 9 shows the load-displacement responses predicted by the macroscale descriptions  
487 which are compared against the reference mesoscale curves. In the main, the macromodels  
488 predict the initial stiffness of the masonry arches accurately, yet significantly overestimating  
489 the peak strengths without providing a realistic representation of the post-peak behaviour as  
490 given by the reference mesoscale models. Furthermore, very different macromodel curves are  
491 obtained depending on the adopted value for  $w_c$ . In particular, the largest differences between  
492 the mesoscale models and the corresponding macromodels are achieved when  $w_c=1$ . It should  
493 be noted that this value is recommended by most software implementations [37][38] to model  
494 the cyclic response of quasi-brittle materials.

495 In Figure 10, the influence of the dilatancy angle on the global response of the weak masonry  
496 bare arch is shown. Since this parameter governs the normal plastic deformation due to shear,  
497 it is expected that by increasing  $\psi$  the global behaviour becomes more ductile due to the  
498 confinement effects exerted by the surrounding elements. Given its high influence on the global  
499 behaviour, it is apparent that more accurate calibration is needed for such critical parameter.



500 Figure 9. Load-displacement curves for the (a) bare arch and (b) the confined arch with weak masonry  
 501 and (c) the bare arch and (d) confined arch with strong masonry.



502  
 503 Figure 10. Influence of dilatancy angle on the global response.

504 The deformed shapes at failure (at the last step of the analysis) with the tensile damage contour  
505 distributions are depicted in Figure 11. The models with strong masonry exhibit flexural failure  
506 (Figure 11c,d) which is in agreement with the failure mode predicted by the mesoscale models  
507 (Figure 7c,d) and mostly characterised by damage in tension concentrated at the intrados and  
508 extrados of the arch corresponding to the opening of plastic hinges. The models with weak  
509 masonry show a mixed failure mechanism with a clear local shear in the arch with a marked  
510 punching shear effect developing underneath the area where the external load is applied and  
511 flexural damage at the extrados of the arch at the side opposite to the loaded area (Figure 11a,b).  
512 This is not predicted by the mesoscale model, which shows ring separation at failure (Figure  
513 7a,b). This main difference confirms the inability of the continuum isotropic damage-plasticity  
514 model to represent shear sliding between adjacent rings, which is a characteristic failure  
515 mechanism of multi-ring arches well captured by detailed mesoscale models. Moreover for the  
516 arches with strong masonry, the use of the macroscale continuum isotropic model leads to a  
517 significant overestimation of the ultimate strength and ductility, where the numerical  
518 predictions are affected significantly by some model parameters (e.g.  $w_c$  and  $\psi$ ) which cannot  
519 be determined via simplified calibrations.

#### 520 *4.3.2. Advanced calibration procedure*

521 To improve the accuracy of the macroscale predictions, the advanced calibration procedure  
522 described in Section 3 has been applied to determine the macromodel material parameters,  
523 focusing first on the specimens with weak masonry where the initial macroscale predictions,  
524 based on a simplified calibration of the model material parameters, were not in good agreement  
525 with the mesoscale results.

526 The bare arch in Figure 5a, subjected to two constant initial forces at the quarter span ( $L/4$ ) and  
527 three-quarter span ( $3/4L$ ), both equal to 16kN, and to a patch load applied at  $L/4$  and increased  
528 up to collapse, is used as the virtual test for the calibration of the model parameters. The specific

529 loading condition ignoring the self-weight contribution of masonry has been chosen to activate  
 530 both flexural damage and shear sliding between adjacent the rings, thus providing suitable  
 531 information to the optimisation algorithm.

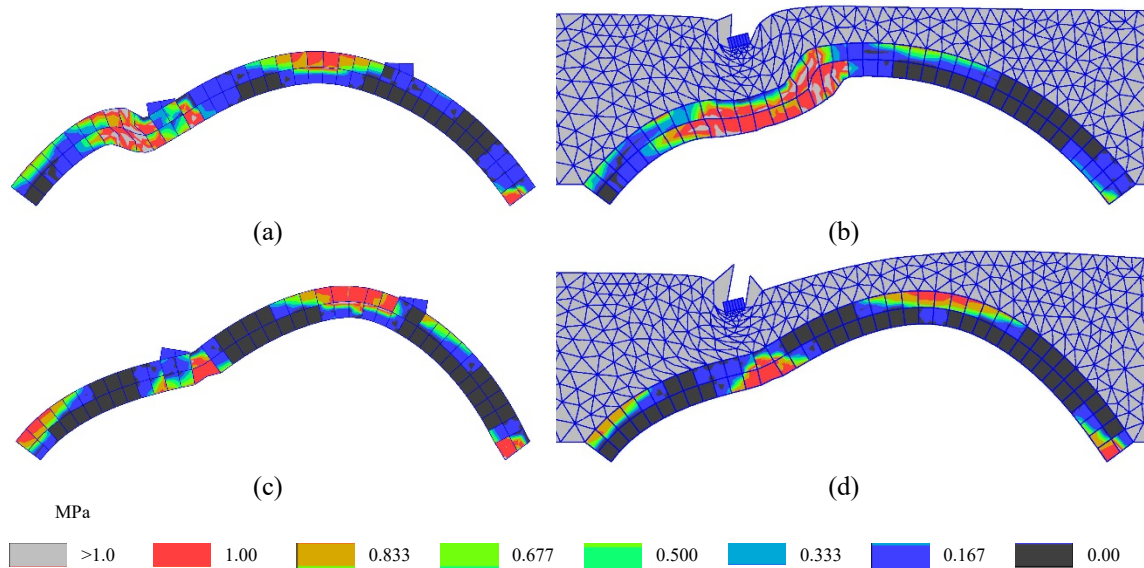


Figure 11. Deformed shape and damage in tension contours at failure for the weak bare (a) and confined (b) arches, and for strong bare (c) and confined (d) arch specimens.

532  
 533 It should be pointed out that the selection of appropriate virtual tests which should activate the  
 534 most critical failure modes of the investigated masonry specimens is the critical step for a  
 535 successful application of the proposed calibration strategy. For instance, in the case of multi-  
 536 ring arches, the failure mechanism of the virtual test should be characterised by flexural  
 537 damage, namely the activation of one or more plastic hinges and shear sliding along the rings.  
 538 In the alternative case, if it is not possible to identify a virtual test with these characteristics,  
 539 multiple virtual tests may be considered, and the multi-objective optimisation procedure should  
 540 consider error functions for each of these.  
 541 Some parametric analyses, not included in the paper for the sake of brevity, have been  
 542 performed to identify the parameters that affect most significantly the arch response. As a result  
 543 of these parametric analyses, and to limit the computing time associated with the model



544 calibration, six model parameters are considered as unknowns in the optimisation procedure:  
 545 the Young modulus ( $E$ ), the tensile strength ( $f_t$ ) and fracture energy ( $G_f$ ), the ratio ( $\tilde{f}_y$ ) between  
 546 the yielding and ultimate compression strength, the parameter governing stiffness recovery  
 547 from tension to compression ( $w_c$ , see Section 2.1) and the angle of dilatancy ( $\psi$ ). The range for  
 548 each parameter is reported in the Table 4. The remaining parameters of the solid elements are  
 549 fixed equal to the default values in Table 3.

550 A load-based partitioning strategy is used for the solution of the calibration problem based on  
 551 two objectives as defined in Eqs (16), (18) with  $\omega_1$ ,  $\omega_2$  the errors due to the loads at L/4  
 552 ( $F_1$ ) and 3/4L ( $F_2$ ), normalised with respect to a reference value with the same units (final  
 553 squared strain energy, divided by the time interval of the virtual test, [J<sup>2</sup>/s]).

554 The evaluated PF (Figure 12) appears discontinuous; the minimum of  $\omega_2$  ( $4 \cdot 10^{-3}$ ) is reached  
 555 for  $\omega_1 = 0.10$ . Conversely, the minimum of  $\omega_1$  ( $6 \cdot 10^{-4}$ ) corresponds to a much larger value  
 556  $\omega_2 = 1.19$ . This implies that calibrating the response on the force-displacement curve of the  
 557 variable load at L/4 (minimum  $\omega_1$ ) may entail significant error in the total energy.

558 In **Error! Reference source not found.a**, the solutions on the PF are shown using the  
 559 normalised errors  $\omega_1^* = (\omega_1 - \omega_1^{min})/(\omega_1^{max} - \omega_1^{min})$  and  $\omega_2^* = (\omega_2 - \omega_2^{min})/(\omega_2^{max} -$   
 560  $\omega_2^{min})$  ranging from 0 to 1. In the following, the solution corresponding to the minimum global  
 561 error  $\omega_{min}^* = \min \left\{ \sqrt{\omega_1^{*2} + \omega_2^{*2}} \right\} = 0.078$ , highlighted in the graph, is considered as the  
 562 reference solution.

563 **Error! Reference source not found.b** shows the displacement  $d_l$  at L/4 against the load  $F_l$   
 564 associated with all the solutions of the PF compared to the response of the mesoscale model.  
 565 Three families of curves can be observed: the curves that minimise the error related to  $F_1$ ,  
 566 which fit very well the response of the mesoscale; the curves that minimise  $\omega_2$ , i.e., the error  
 567 related to  $F_2$ , which provide a significant overestimation of the peak-strength of the arch (from  
 568 35 kN to 45 kN) and the curves that minimise the global dimensionless error allowing for the

569 two objectives of the optimisation procedure. These solutions and in particular the solution  
570 associated with the minimum error provide a good prediction in terms of initial stiffness and  
571 peak load, but they show some differences regarding the post-peak stage. In **Error! Reference**  
572 **source not found.c**, the vertical displacement at three-quarter  $d_2$  span is plotted against the  
573 load  $F_1$ . It is possible to see that in no case the final uplift shown by the mesoscale model is  
574 attained by the macroscale models of the PF. However, the dashed line, identifying the  
575 compromise solution of the multi-objective optimisation, shows a general good agreement with  
576 the mesoscale curve.

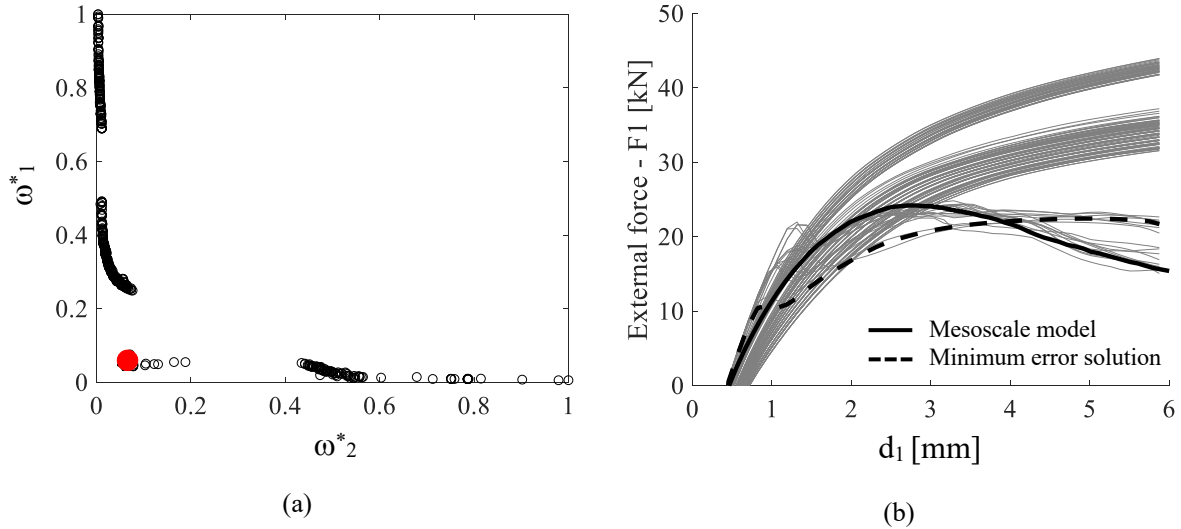
577 It is important to point out that even when the minimum error solution provides a satisfactory  
578 approximation of the load-displacement curves, it can lead to a failure mechanism inconsistent  
579 with that predicted by the virtual test. Therefore, a comparison in terms of failure mechanism  
580 is crucial to choose the best solution from the PF.

581 The ultimate deformed shapes corresponding to the minimum error solution and the two ends  
582 of the PF, each providing the minimum of one error function, are reported in Figure 12d. It can  
583 be noticed that the solution corresponding to  $\omega_{2,\min}^*$  (best fit of the load-displacement at the  
584 loaded point) does not indicate shear failure of the arch, leading to a wrong failure mode  
585 prediction.

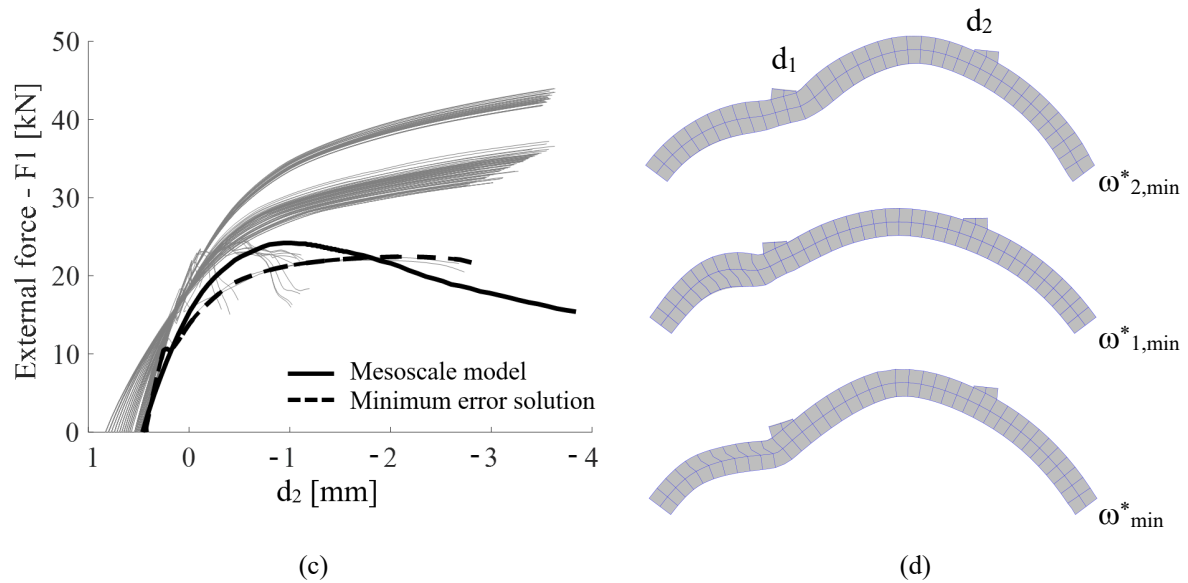
586 Conversely, the two other solutions predict different shear failure mechanisms. In particular,  
587 the minimum error solution ( $\omega_{\min}^*$ ) represents well the failure mode of the virtual test by  
588 predicting circumferential shear deformations, in accordance with the shear sliding among  
589 adjacent rings determined by the mesoscale model.

590 In Table 4, the calibrated parameters for the minimum error solution are shown, together with  
591 the range identified by the solutions in its surroundings ( $\omega^* < 2.5\omega_{\min}^*$ ). These latter values  
592 allow investigating the sensitivity of the calibrated response to each parameter, compared to  
593 the initial variation range. Analysing the results, all parameters seem rather univocally

594 determined as a significant reduction of the ranges is observed, with the only exception of  
 595  $\tilde{f}_y$ . It is interesting to see that the calibrated  $w_c$  is close to zero, unlikely the typical assumption  
 596 considered in the simplified calibration.



597



598 Figure 12. Calibration of the continuum weak model: (a) Pareto Front solutions; load-displacement curve at  
 599 (b)  $L/4$  and (c)  $3/4L$ ; (d) failure mechanisms of the minimum error and PF solutions.

600 The results of the model calibration are validated by analysing the confined and bare arches  
 601 subjected to the loading condition described in the initial mesoscale simulation in Section 4.2.  
 602 Moreover, two additional load conditions for the bare arch are investigated in which the arch  
 603 is subjected to a concentrated force alternatively applied at the mid span and at one eight span

604 without initial symmetric forces. The load-displacement curves for the weak masonry models  
 605 are shown in Figure 13, where the continuum calibrated model is compared against the  
 606 mesoscale model and the continuum model with simplified calibration procedure for  $w_c=0$ .

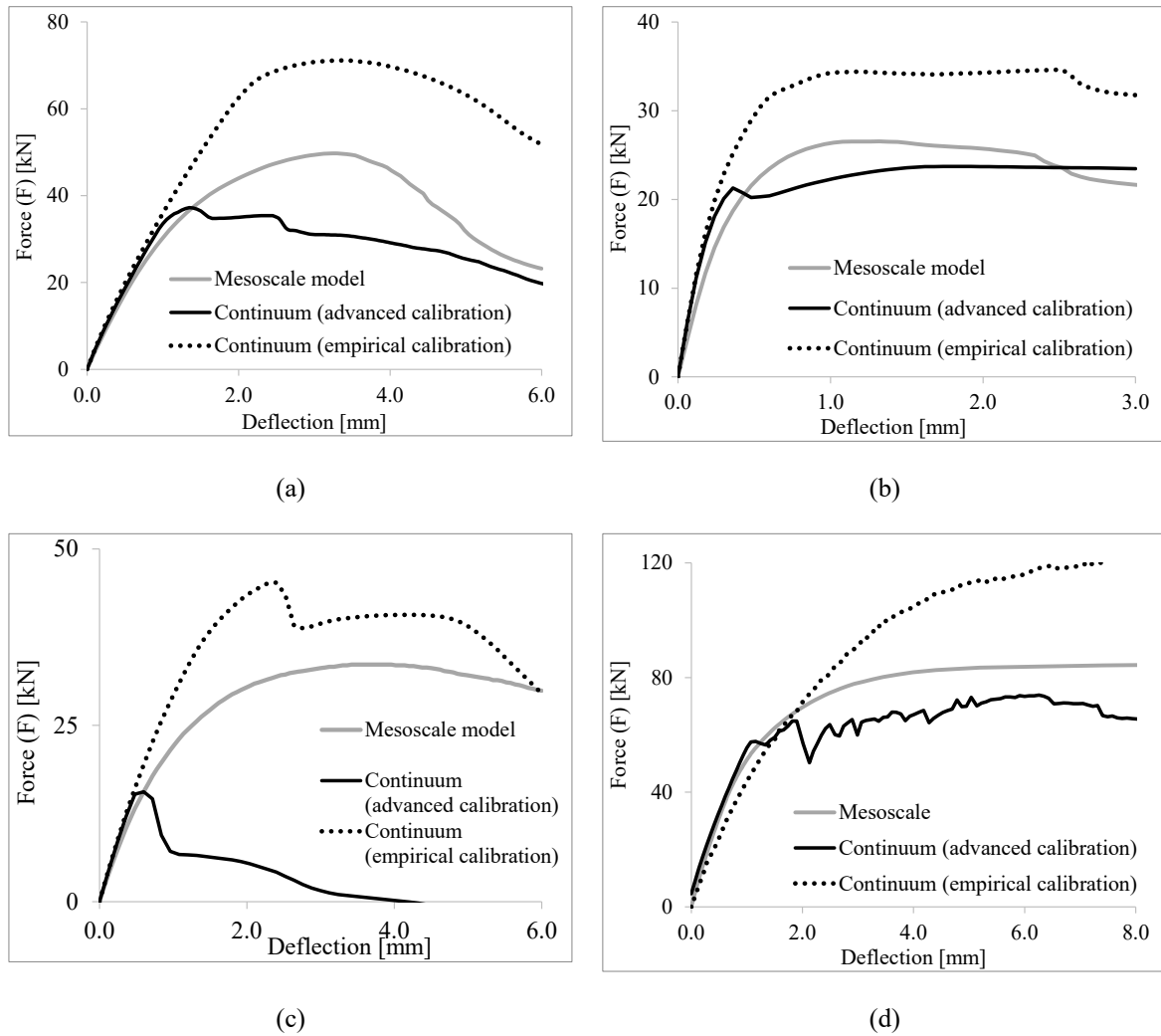
607 **Table 4:** Input parameters and results of the calibration procedure for the continuum model (weak masonry)

Parameter	Unit	Initial range		Solutions with		Minimum error Solution
		Lower bound	Upper bound	$\omega^* \leq 2.5\omega_{min}^*$		
				Upper bound	Upper bound	
$E$	MPa	1000	6000	2050	2550	2500
$\psi$	°	0	90	18	26	25
$f_t$	N/mm	0.01	1.0	0.024	0.041	0.025
$G_t$	N/mm	0.001	0.5	0.021	0.028	0.022
$\tilde{f}_y$	-	0.01	1.0	0.40	0.83	0.72
$w_c$	-	0.0	1.0	0.00	0.16	0.02

608  
 609 Generally, the model calibrated by the advanced procedure exhibits a much-improved  
 610 agreement with the mesoscale model. However, in the case of the bare arch loaded at the  
 611 quarter-span, the continuum model shows a premature shear failure which leads to a significant  
 612 underestimation of the maximum load and displacement capacity of the arch (Figure 13c). In  
 613 the case of bare arch loaded at the mid span (Figure 13a), the macromodel calibrated by the  
 614 advanced procedure underestimates the peak-load value. It provides, however, an adequate  
 615 prediction of the residual strength and pre/post peak response. Finally, a satisfactory  
 616 comparison can be observed in the cases of bare arch loaded at one eighth span and the confined  
 617 arch specimen (Figure 13b and 13d).

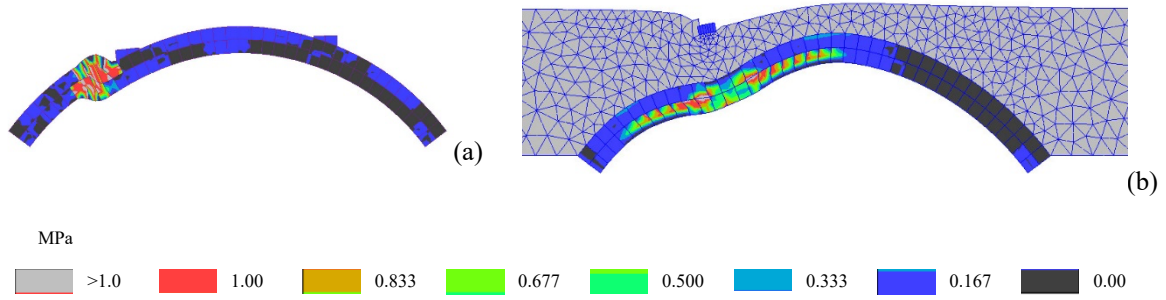
618 The failure mechanism of the bare arch, displayed in Figure 14a, is characterised by an evident  
 619 punching effect due to the shear failure of the element of the arch underneath the load, which

620 is not observed in Figure 7a. On the contrary, the failure mechanism of the confined arch  
 621 (Figure 14b) is rather consistent with the mechanism obtained by the mesoscale model (Figure  
 622 7b). The sliding between adjacent rings is represented by shear failure of the solid elements;  
 623 however, unlikely the continuum model calibrated by means of the simplified procedure  
 624 (Figure 11b), the punching effect is not observed.



625 Figure 13. Load-displacement capacity curves for the bare arch loaded with a concentrated force at (a)  
 626 mid-span, (b) one eighth span, (c) one quarter span with initial forces, and (d) for the confined arch.

627  
 628  
 629



630 Figure 14. Failure mechanism and damage in tension contours predicted by the calibrated continuum  
 631 model for the (a) bare arch and (b) the confined arch.

632 In conclusion, it can be affirmed that the rigorous calibration procedure allows for a substantial  
 633 improvement of the continuum model predictions in representing the confined arch specimen.  
 634 However, the calibrated continuum model still shows evident limits in simulating the bare arch  
 635 response, as it provides an unrealistic failure mode underestimating both the global strength  
 636 and ductility of the arch.

637 It is worth pointing out that despite the fact that the adopted calibration strategy is based on the  
 638 use of the entire arch specimen for the virtual test, its applicability is still computationally  
 639 efficient as it applied on a simple 2D strip model neglecting the backfill and its interaction with  
 640 the arch.

#### 641 4.4 Hybrid model simulations

642 In this section, the proposed hybrid model described in Section 2 is employed to simulate the  
 643 bare and the confined arch of Figure 5. The same mesh with quadratic solid elements  
 644 considered in the previous macroscale continuum simulations has been employed, but  
 645 circumferential nonlinear interfaces (Section 2) have also been introduced to connect each pair  
 646 of adjacent solid elements along the thickness of the arch according to the proposed hybrid  
 647 macroscale representation (Figure 1b). The model material parameters of the solid elements  
 648 and the circumferential interfaces are calibrated through the rigorous procedure described in  
 649 Section 3 and applied before to determine the material properties for the continuum macroscale

650 model. Nine model parameters are calibrated by the optimisation procedure. Namely, the same  
651 six parameters for the solid elements, already considered in Section 4.3.2 with the addition of  
652 three parameters characterising the interfaces: the shear stiffness ( $k_t^M$ ), the cohesion ( $c^M$ ) and  
653 the shear fracture energy ( $G_s^M$ ) which have been identified as the most significant interface  
654 parameters affecting the response of multi-ring arches [49]. The remaining interface parameters  
655 are assumed either coincident to the parameters of the mesoscale model ( $f_c, G_c, tg\phi, tg\phi_g$ ) or  
656 proportional to the parameters assumed as unknown in the optimisation ( $k_n^M, f_t^M, G_t^M$ ) as  
657 indicated below:

$$\begin{aligned}
k_n^M &= k_t^M \cdot \frac{k_n^m}{k_t^m} \\
f_t^M &= c^M \cdot \frac{f_t^m}{c^m} \\
G_t^M &= G_s^M \cdot \frac{G_t^m}{G_s^m}
\end{aligned} \tag{19}$$

658 where the superscripts  $M$  and  $m$  refer to the macroscale and mesoscale representation,  
659 respectively.

660 In the following, the interface stiffness and cohesion parameters are represented by the non-  
661 dimensional coefficients:

- 662 -  $k^* = 2k_t^M t(1 + \nu)/E$ , where  $E$  and  $\nu$  are the Young modulus and the Poisson's  
663 coefficient of the solid elements and  $t$  a fictitious thickness equal to 1 mm.
- 664 -  $c^* = c^M/f_{v0}$ , where  $f_{v0}$  is the initial shear strength of the solid elements evaluated  
665 through Eq. (5) with  $\boldsymbol{\kappa} = \mathbf{0}$ .

666 The variation ranges of the unknown parameters are reported in Table 5. The calibration was  
667 performed for both weak and strong masonry by using the same procedures and objective  
668 functions as for the continuum model.

669 Figure 15a displays the solutions belonging to the PF ( $\omega_1^* - \omega_2^*$ ) for the calibration of the weak  
670 masonry model, while the corresponding load-displacement curves are reported in Figure  
671 15b,c. Comparing these results to the solutions of the optimisation for the continuum model in  
672 Section 5.1, it can be observed that:

- 673 - The solutions of the PF are uniformly distributed in the space  $\omega_1$ - $\omega_2$  and are  
674 associated with much lower errors.
- 675 - The load-displacement curves are less dispersed and much closer to the mesoscale  
676 predictions.

677 These remarks confirm that the further free parameters included in the optimisation algorithm  
678 ( $k^*, c^*, G_t^M$ ) effectively improve the quality of the results. In this case, the absolute minimum  
679 of  $\omega_2$  ( $4.88 \cdot 10^{-3}$ ) is reached for a value of  $\omega_1 = 4.97 \cdot 10^{-3}$  while the minimum of  $\omega_1$   
680 ( $4.66 \cdot 10^{-5}$ ) corresponds to  $\omega_2 = 7.57 \cdot 10^{-3}$ . Following the procedure in Section 5.1, the  
681 two errors are normalised leading to a minimum solution error  $\omega_{\min}^* = \min\left(\sqrt{\omega_1^{*2} + \omega_2^{*2}}\right) =$   
682 0.2284 which corresponds to the set of model parameters reported in Table 5.

683 Contrary to what is observed in the case of the continuous model (Figure 12d), here the  
684 minimum and PF solutions provide the same failure mechanism (Figure 15d), confirming the  
685 ability of the hybrid model to represent the sliding mechanism among adjacent rings.

686 The PF for the model with strong masonry is shown in Figure 16a. The minimum error of  $\omega_2 =$   
687 0.011 is reached for  $\omega_1 = 0.0056$ , while the minimum error  $\omega_1 = 0.0012$  is associated with  
688  $\omega_2 = 0.337$ . The minimum solution error corresponds to  $\omega_{\min}^* = 0.3207$ , ( $A$  in Figure 16a)  
689 and the matching set of parameters are indicated in Table 5. It can be noticed that the latter  
690 solution is characterised by a low value of the dimensionless cohesion ( $c^*=0.69$ ). This  
691 circumstance may potentially lead to a response characterised by sliding between the ring  
692 which is in disagreement with the mesoscale results. For this reason, another solution ( $B$  in



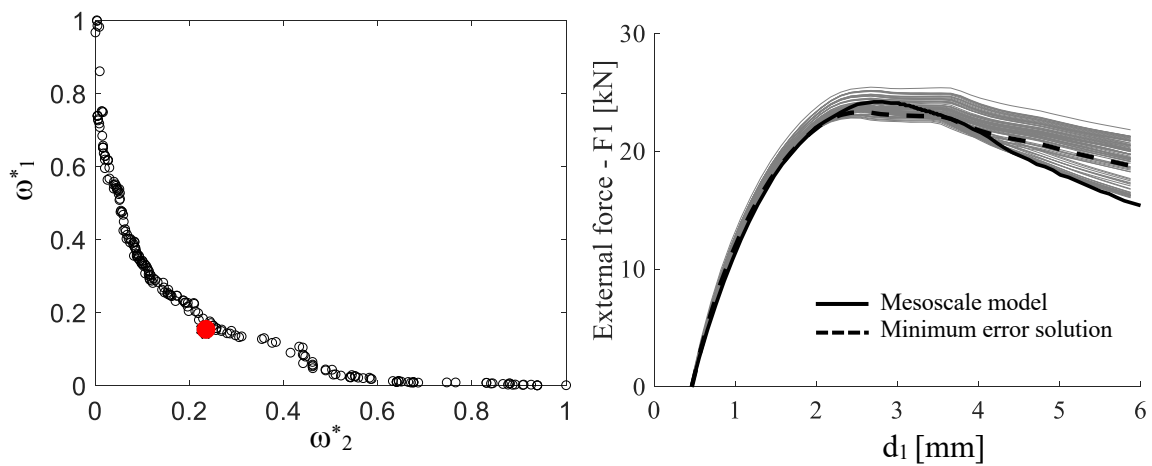
693 Figure 16a), corresponding to  $\omega_B^* = 0.4673$ , is also considered. This solution has been chosen  
694 as the solution with minimum error among those characterised by  $c^* > 1$ .

695 The response curves of the PF solutions are shown in Figure 16b and 16c with the two reference  
696 solutions reported with dashed lines. Two families of curves, which tend to minimise the two  
697 objectives separately are visible. Although the PF is less regular than the previous case, the  
698 selected solutions confirm a good match with the mesoscale curves. Finally, Figure 16d depicts  
699 the failure mechanisms corresponding to the two selected solutions, A and B. It is observed  
700 that the two failure mechanisms are relatively consistent with each other. However, the failure  
701 mechanism associated with solution A comprises sliding along the circumferential interface,  
702 evidencing ring separation not observed in the virtual test. On the contrary, Solution B provides  
703 a better approximation of the flexural failure mechanism of the virtual test, thus it corresponds  
704 to the best solution to calibrate the macroscale hybrid description for the strong-masonry arch.

705 Analogously to the procedure for the continuum model in Section 4.3.2, the results of the  
706 calibration are validated considering the bare and confined specimens, plus two further models  
707 representing the bare arch subjected to a concentrated force at mid span and at one-eighth span.  
708 The load-displacement curves of the hybrid model are displayed in Figure 17, where they are  
709 compared against the mesoscale curves and the predictions of the continuum macroscale model  
710 calibrated by the advanced procedure in Section 4.3.2.

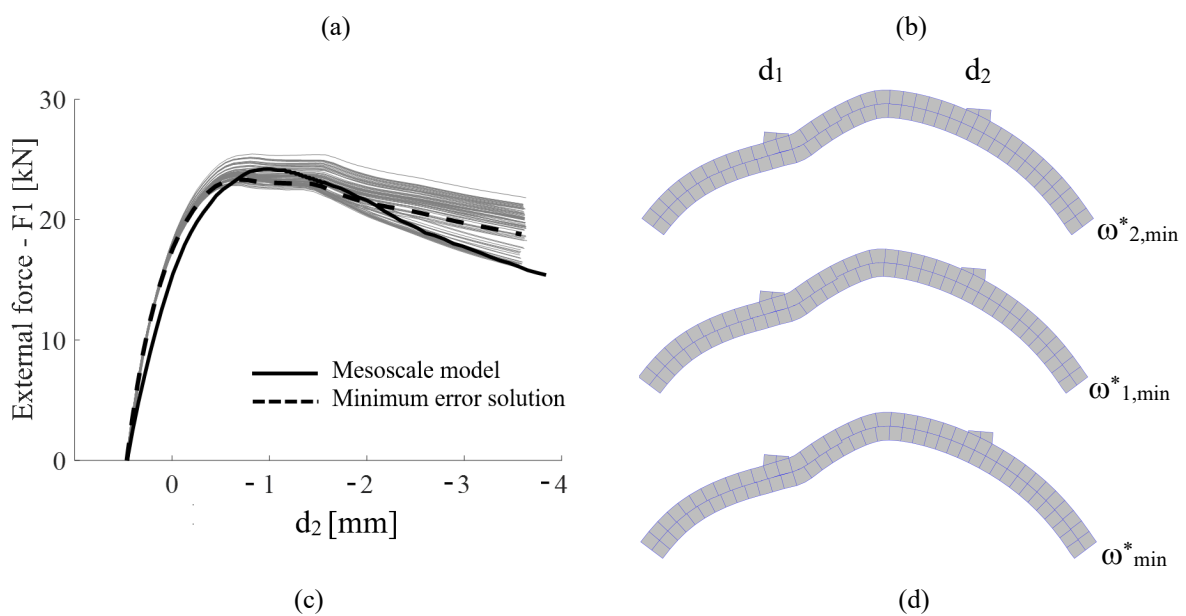
711 The results of the hybrid macromodel are in a good agreement with those obtained by the  
712 mesoscale model confirming a generally improved prediction compared to the results provided  
713 by the continuum macromodel. The only exception is represented by the load condition with  
714 the force at  $L/8$ , where the continuum model provides a better prediction of the peak-loads.  
715 However, the curve of the hybrid model, also in this case, is more consistent to the mesoscale  
716 curve in the pre- and post-peak stages. It appears that the presence of the backfill reduces the  
717 differences between the results, with mesoscale and hybrid model curves almost coincident.

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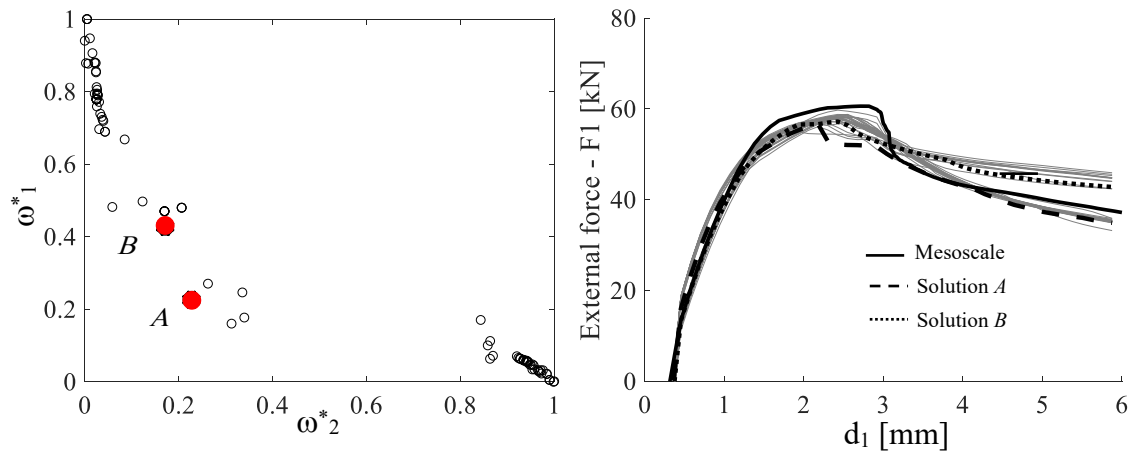
Figure 15. Calibration for the weak masonry model: Pareto Front solutions (a); capacity curves at one quarter (b) and three quarter span (c); failure mechanisms of the minimum error solution and the frontier solutions of the PF (d).

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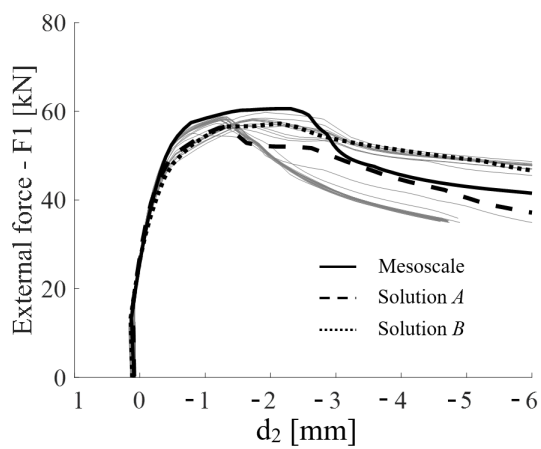


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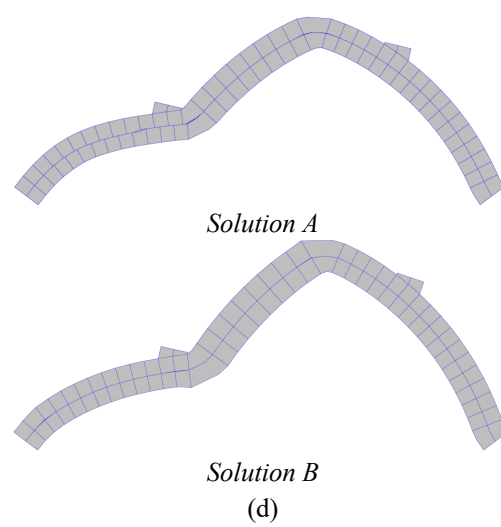
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(a)

(b)



(c)



(d)

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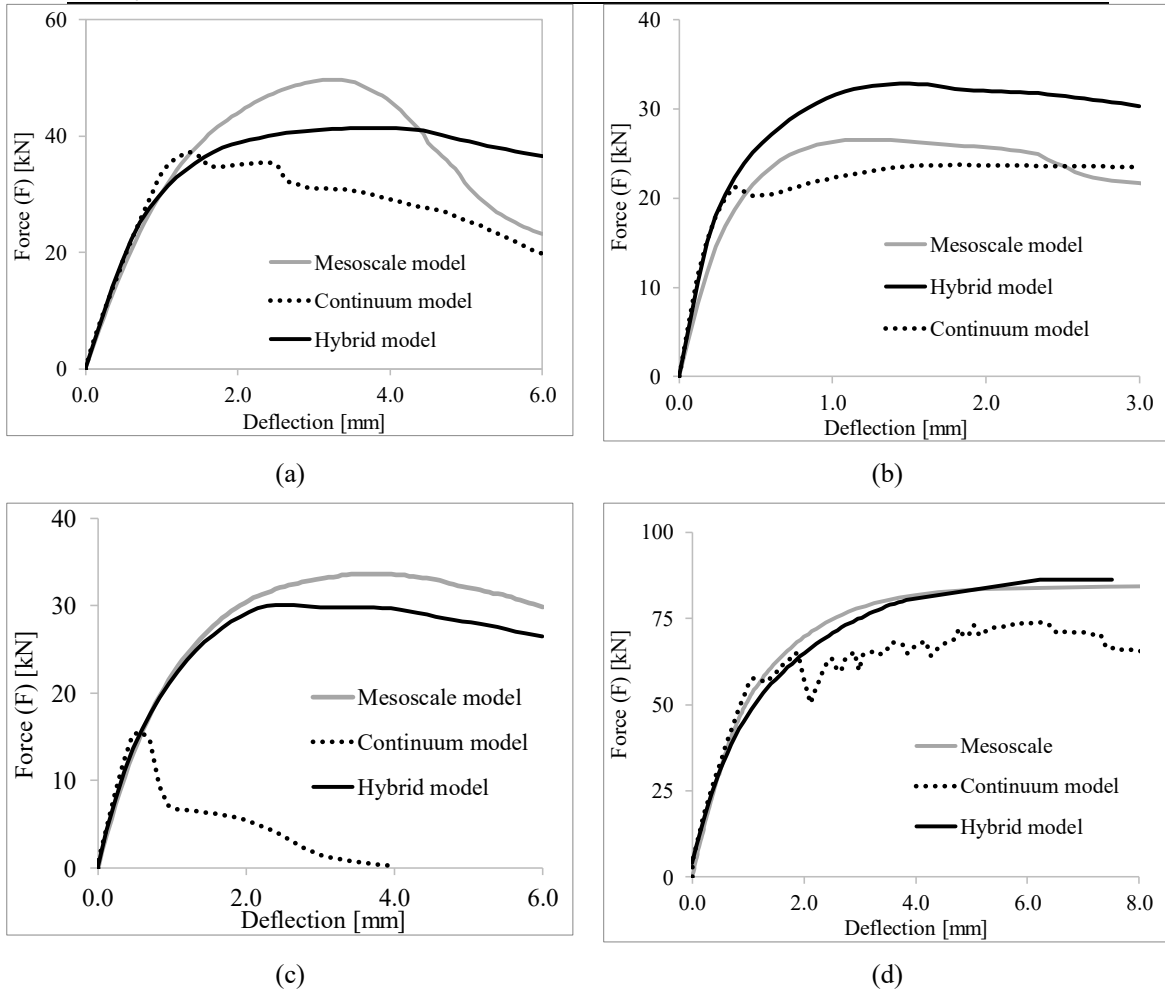
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Figure 16. Calibration for the strong masonry model: Pareto Front solutions (a); capacity curves at one quarter (b) and three quarter span (c); failure mechanisms of the two selected solutions A and B (d).

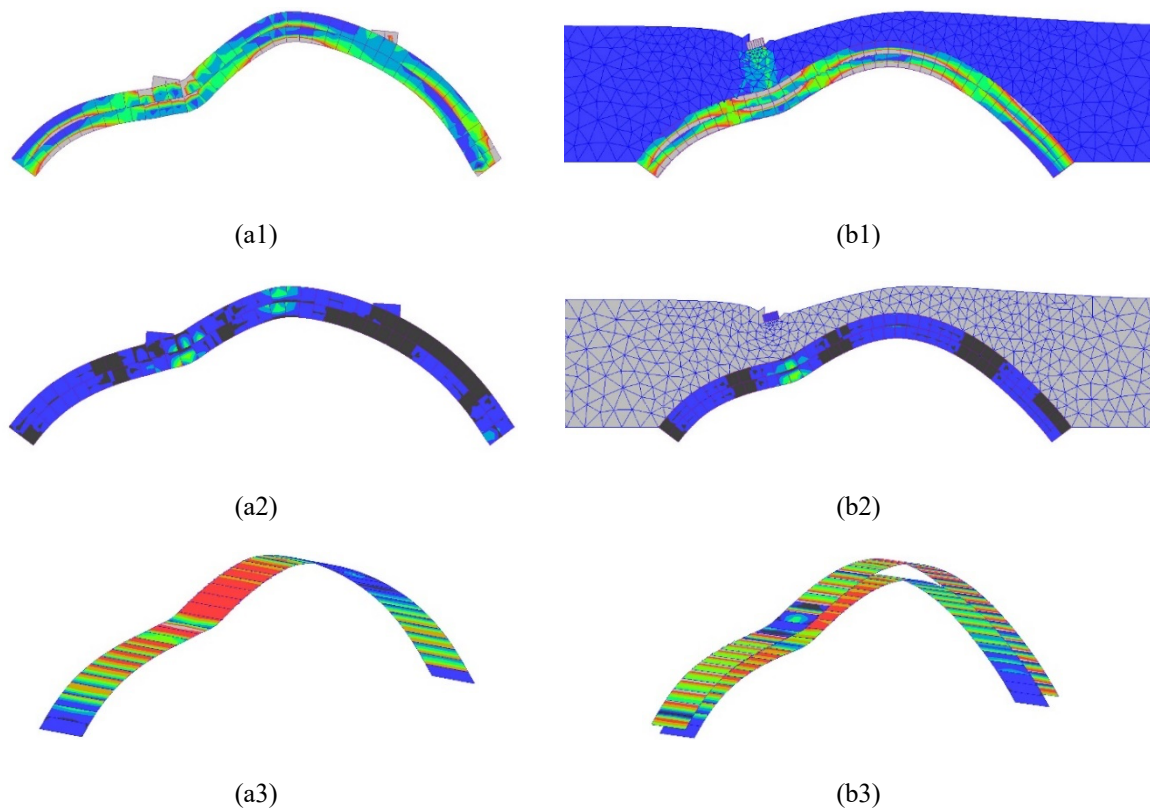
**Table 5:** Input parameters and results of the calibration procedures for the hybrid model.

Parameter	Unit	Initial range		Minimum error solution	Minimum error solutions	
		Lower bound	Upper bound		strong masonry	
					Solution A	Solution B
$E$	MPa	1000	6000	2350	4700	4200
$\psi$	°	0	90	63	23	59
$f_t$	N/mm	0.01	1.0	0.083	0.164	0.137
$G_t$	N/mm	0.001	0.5	0.089	0.042	0.03
$\tilde{f}_y$	-	0.01	1.0	0.79	0.71	0.68
$w_c$	-	0.0	1.0	0.10	0.48	0.00
$k^*$	-	0.01	1.5	0.01	0.01	0.04
$c^*$	-	0.0	2.0	0.62	0.69	1.44
$G_t^M$	N/mm	0.01	0.25	0.096	0.167	0.213



741 Figure 17. Weak masonry model: load-displacement capacity curve for the bare arch loaded with a  
 742 concentrated force at (a) mid-span, (b) at one eighth span and (c) at one quarter span with initial forces,  
 743 and (d) for the arch interacting with backfill loaded at one quarter span.

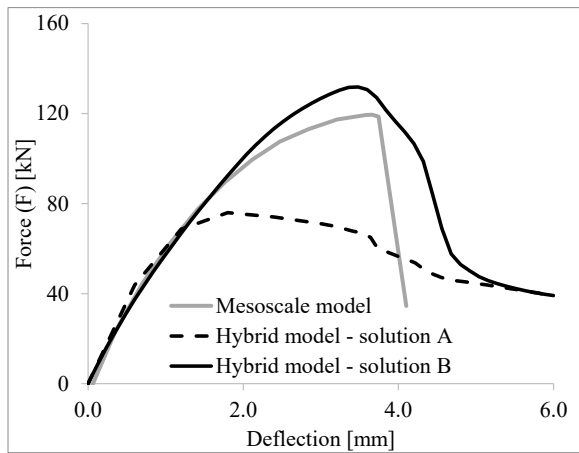
744 The failure modes of the bare arch and the arch interacting with backfill are displayed in Figure  
 745 18, where the von-Mises equivalent stress distribution in the solid elements and the damage  
 746 contours on the interface elements are also shown. A good agreement between the failure  
 747 modes of the hybrid model and those obtained by the mesoscale model, both in terms of stress  
 748 distribution (Figure 7) and damage index distribution along the ring-to-ring and arch-to-  
 749 backfill interfaces (Figure 8). Importantly, the ring sliding mechanism, which could not be  
 750 predicted by the continuum model, is well described by the proposed hybrid macroscale  
 751 representation.



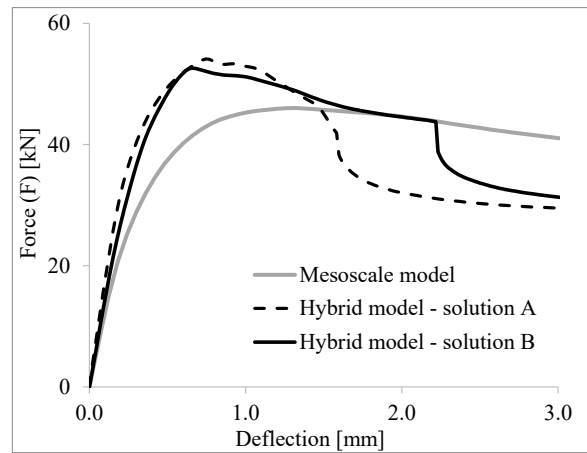
753 Figure 18. Weak masonry model: deformed shape at collapse for the (a) bare arch and (b) the arch  
 754 interacting with backfill loaded at quarter span showing (a1, b1) von Mises stresses, (a2, b2) damage in  
 755 tension contours for the solid elements and (a3, b3) the interface elements.

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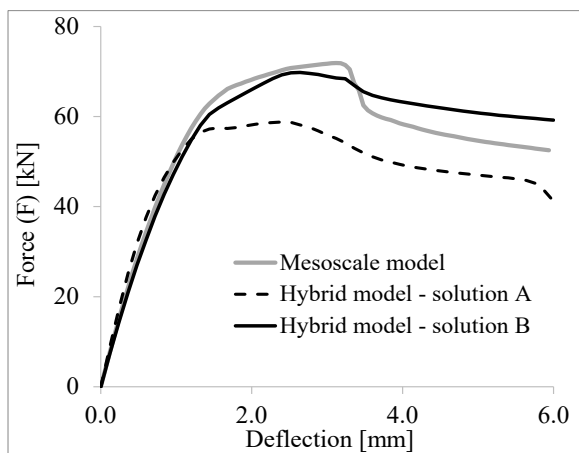
758 The results of the calibration analyses in term of load-displacement curves and failure  
759 mechanisms for the strong masonry model are shown in Figure 19 and in Figure 20,  
760 respectively. The curves obtained using the solution *B* parameters indicate a good agreement  
761 with the mesoscale results for all the considered models and loading conditions. Conversely,  
762 solution *A* leads to a significant underestimation of the ultimate strength of the structure in the  
763 cases of the bare arch loaded at  $L/2$  (Figure 19a) and the arch interacting with backfill (Figure  
764 19d). The failure mechanisms obtained by the macromodel calibrated with the solution *B*  
765 (Figure 20) result in a good agreement with those obtained by the mesoscale model (Figure 7).  
766 In conclusion, it can be stated that the advanced calibration procedure leads to a realistic set of  
767 mechanical parameters to describe the global response of the arches, including strong masonry  
768 arches, under a wide range of boundary and loading conditions. The adopted approach to select  
769 the reference solution from the results of the multi-objective optimisation procedure, based on  
770 the analysis of the Pareto Front, appears to be sufficiently accurate and robust. However, the  
771 circumstance by which the minimum error solution (*A*) provided less satisfactory results  
772 compared to those obtained by another solution belonging to the PF solution (solution *B*)  
773 denotes that further improvements to the selection strategy from the PF and/or the definition  
774 of the multi-objective optimisation problem may be needed. This open issue will be  
775 investigated by the authors in future studies.



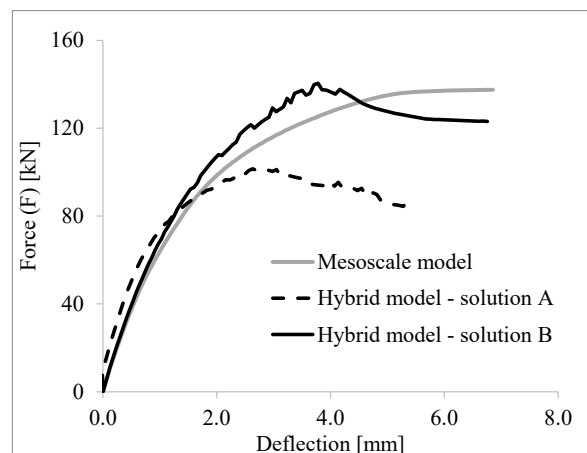
(a)



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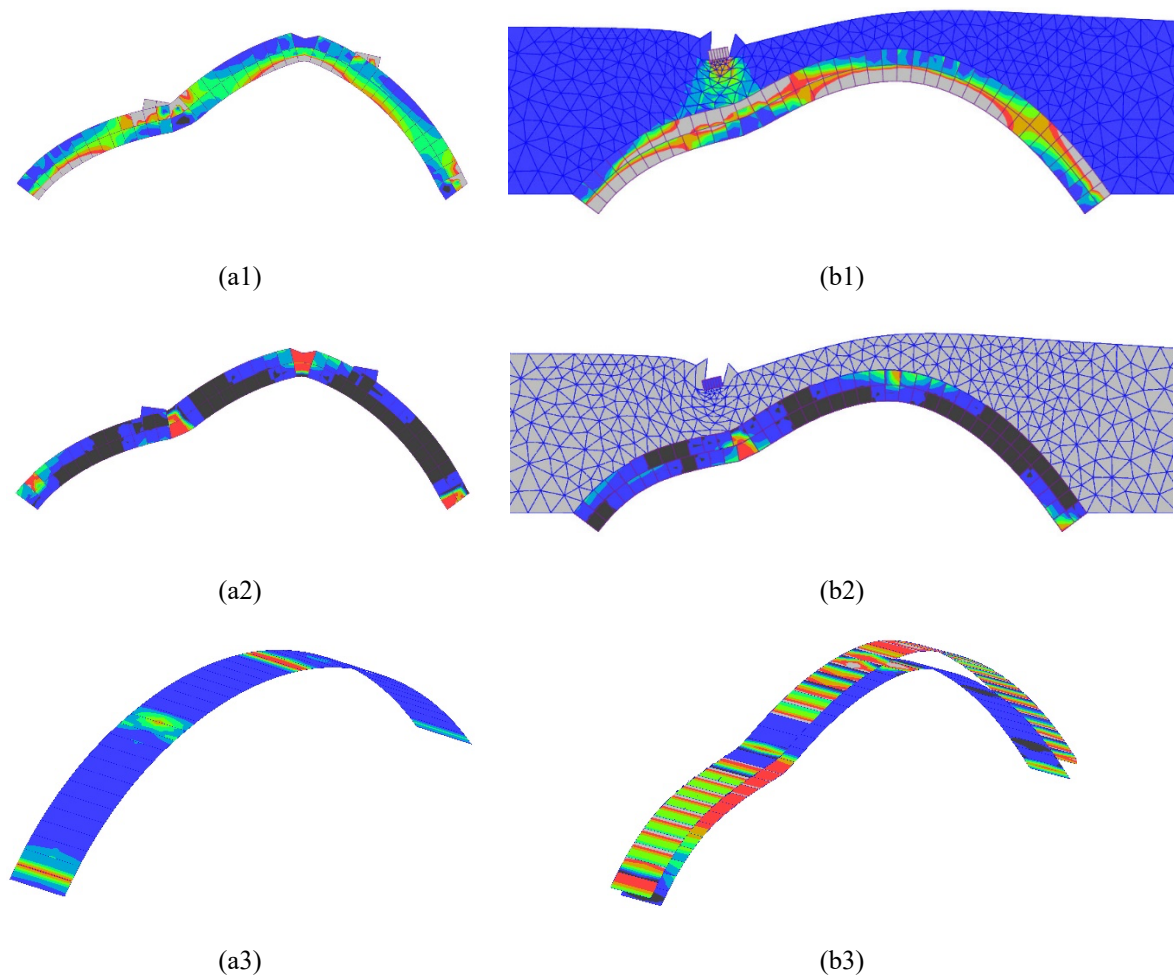
(c)



(d)

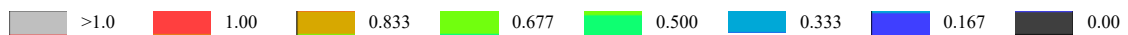
776 Figure 19. Load-displacement curves for the strong masonry bare arch loaded at mid-span (a), one eighth  
777 span (b), one quarter span (c) and for the arch interacting with backfill loaded at quarter span (d).

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Colour map for stress (MPa) and damage index (-)



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Figure 20. Strong masonry model: deformed shape at collapse for (a) the bare arch and (b) the arch

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interacting with backfill loaded at quarter span showing (a1, b1) von Mises stresses, (a2, b2) damage in

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tension contours for the solid elements and (a3, b3) the interface elements.

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## 5 NUMERICAL SIMULATIONS OF LONG SPAN MASONRY ARCHES AND

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### BRIDGES

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In this section, the hybrid model and the advanced calibration procedure are employed to

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analyse a masonry arch with seven rings and a 16 m span. The numerical results, obtained using

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ADAPTIC [31], enable investigating the computational efficiency of the proposed macroscale

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strategy and the improved accuracy due to the use of various nonlinear circumferential interface

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layers for modelling large masonry arches typical of realistic masonry bridges. Similar to the



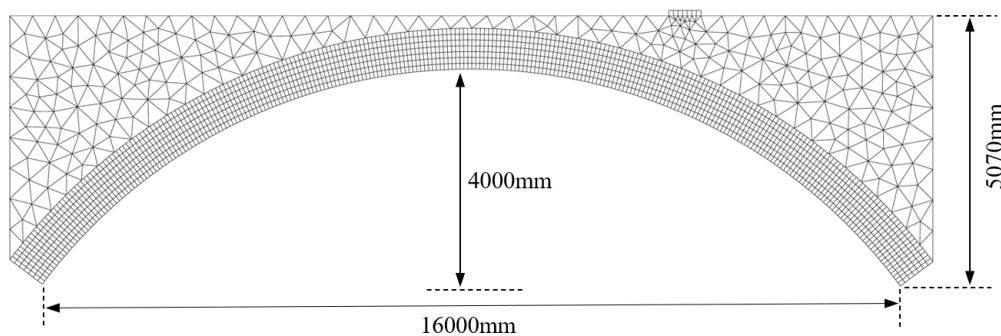
800 discussion in Section 4, the performance of the macroscale hybrid description is assessed based  
801 on load-displacements curves and predicted failure mechanisms which are compared against  
802 computationally expensive mesoscale simulations.

803 The considered masonry arch specimen with stretcher bond is characterised by a 770 mm  
804 thickness and a span-to-rise ratio of 4.0. The considered material properties for the masonry  
805 constituents are those for the weak masonry material reported in Tables 1 and 2. Figure 21  
806 shows the geometrical layout of the mesoscale model for the specimen with a representative  
807 1 m width. The vertical loads are applied at the top of the backfill at quarter span of the arch,  
808 assuming the same boundary conditions described in Section 4, where the bases of the arch and  
809 the backfill are fixed, and the vertical ends of the backfill are restrained against horizontal  
810 displacement. The mesoscale mesh for the arch with only one element along the width  
811 comprises 35672 nodes with a total of 106680 DOFs.

812 The virtual test requires loading the bare arch by the masonry self-weight and two constant  
813 symmetric forces of 22 kN applied at one-quarter and three-quarter span, after which an  
814 increasing force is applied at one-quarter span until the failure of the arch. Three hybrid  
815 macroscale models are adopted, utilising one, two and three circumferential interface layers  
816 equally spaced along the thickness of the arch. Table 6 reports the number of elements and  
817 DOFs for the three arch macromodels (*1-Interface*, *2-Interface* and *3-Interface*), which require  
818 3.3%, 7.7% and 10.3% of the total DOFs of the detailed mesoscale model.

819 The macroscale mechanical parameters resulting from the calibration procedure are  
820 summarised in Table 7. The minimum error solution is considered for each model, while the  
821 parameters corresponding to the two ends of the Pareto Front are also reported only for the  
822 basic mesh with one circumferential interface layer (*1-Interface*). It can be observed that the  
823 model parameters for the three solutions of the *1-Interface* model are very close highlighting  
824 the robustness of the calibration algorithm.

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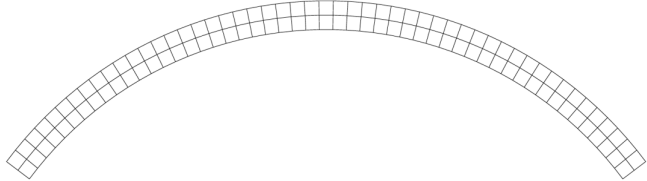
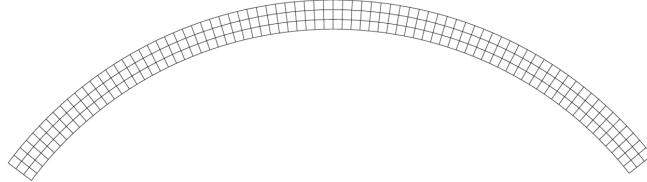
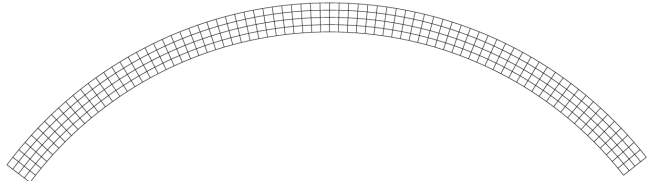


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Figure 21. Large-span arch bridge specimen.

**Table 6:** Characteristics of the three macroscale hybrid meshes.

<i>1-Interface</i>		nodes	1244
		solid FE	204
		interfaces	102
		DoFs	3624
<i>2-Interface</i>		nodes	2804
		solid FE	231
		interfaces	154
		DoFs	8244
<i>3-Interface</i>		nodes	3740
		solid FE	308
		interfaces	231
		DoFs	10992

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**Table 7:** Input parameters and results of the calibration procedures for the hybrid model.

Parameter	Unit	<i>1-Interface</i>			<i>2-Interface</i>	<i>3-Interface</i>
		$\omega_{1,min}^*$	$\omega_{min}^*$	$\omega_{2,min}^*$	$\omega_{min}^*$	$\omega_{min}^*$
$E$	MPa	2500	2500	2500	2600	2450
$\psi$	°	53	54	54	56	55
$f_t$	N/mm	0.053	0.036	0.053	0.051	0.049
$G_t$	N/mm	0.046	0.036	0.036	0.029	0.076
$\tilde{f}_y$	-	0.11	0.17	0.18	0.93	0.43
$w_c$	-	0.92	0.94	0.94	0.88	0.93
$k^*$	-	0.01	0.02	0.02	0.01	0.01
$c^*$	-	1.65	1.64	1.63	1.52	1.71
$G_t^M$	N/mm	0.221	0.250	0.217	0.166	0.168

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832 The load-displacement capacity curves and the deformed shape at failure predicted by

833 macroscale *1-Interface* model are shown in Figures 22a-b, where these are compared against

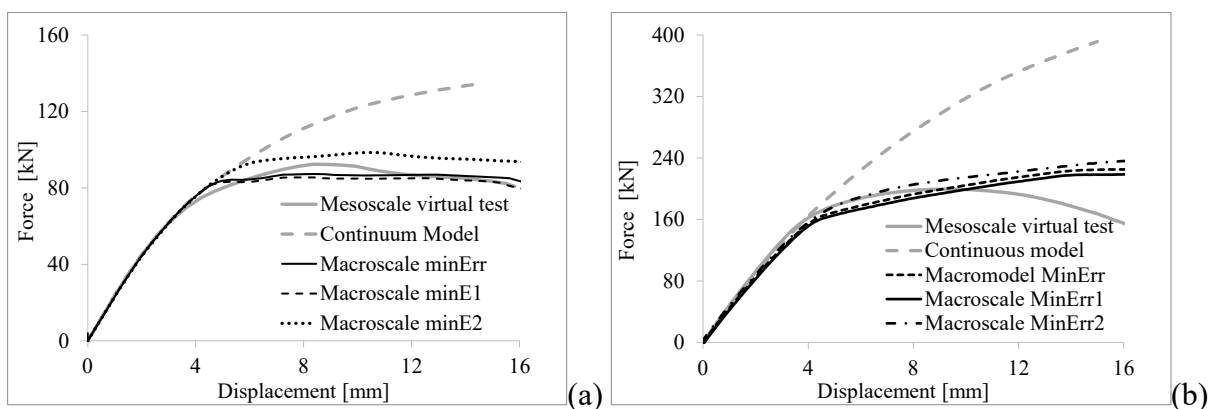
834 the mesoscale results. Figure 22a presents the response curves of the virtual test used for model

835 calibration, whereas Figure 22b shows the results for the arch confined by backfill (bridge

836 specimen) which have been considered for model validation. In the graphs, the response curves

837 of the continuum model without circumferential interface layers and calibrated by the empirical  
838 procedure described in Section 4.1 are also reported for comparison.

839 In the case of the bare arch, the response obtained using the proposed macroscale hybrid model  
840 is very close to the mesoscale results confirming the effectiveness of the developed calibration  
841 procedure. Conversely, the continuum model leads to a significant overestimation of the arch  
842 load-bearing capacity. The failure mechanisms of the macroscale and mesoscale descriptions  
843 for the bare arch are shown in Figure 23, where a good agreement can be observed. Both models  
844 predict a mixed failure mechanism characterised by the activation of three flexural plastic  
845 hinges and ring sliding in the two portions of the arch close to the skewbacks (Figures 23a and  
846 23c). The damage contours at the mesoscale interfaces and those at the circumferential  
847 interface layer of the macroscale *1-Interface* model are shown in Figures 23b and 23d,  
848 respectively. Comparing the load-displacement curves of the arch interacting with backfill  
849 (Figure 22b), it is observed that the hybrid model predicts well the response of the system until  
850 about 10 mm. For larger vertical displacements, the mesoscale model shows softening  
851 behaviour, though the response predicted by the macroscale *1-Interface* models with selected  
852 calibrated material properties exhibits slight hardening. Such discrepancy translates to some  
853 differences in the characteristics of the failure mode shown by the deformed shapes and stress  
854 contours in Figure 24. More specifically, the mesoscale model predicts distributed ring sliding  
855 leading to local shear failure associated with a drastic reduction of the arch resistance. On the  
856 other hand, the macroscale model shows sliding and flexural plastic deformations on the two  
857 continuum portions of the FE mesh separated by the single damaged circumferential interface  
858 layer which is not associated with sudden degradation of the load-bearing capacity of the arch.  
859



860  
861 Figure 22. Load-displacement capacity curves for the (a) bare arch (a) and (b) arch with backfill.

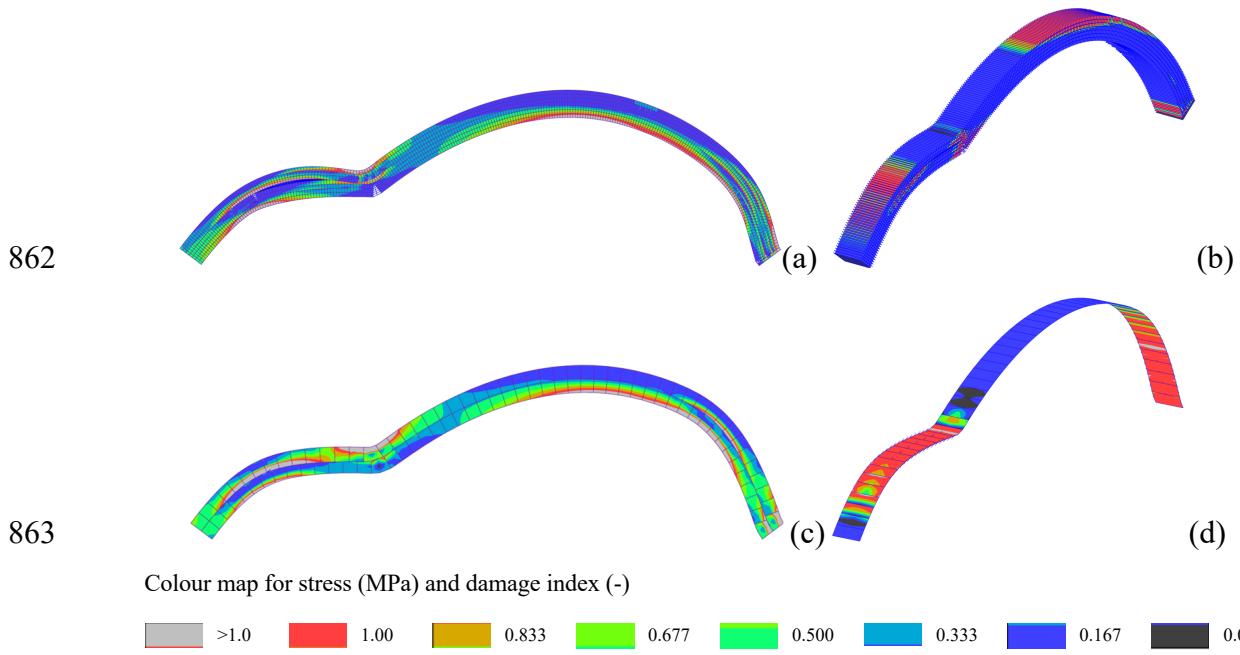


Figure 23. Ultimate deformed shapes of the bar arch (amplificated 50 times) predicted by the (a, b) mesoscale model and (c, d) the macroscale model with (a,c) von-Mises stresses in the solid elements and (b,d) damage in tension contours on the nonlinear interfaces.

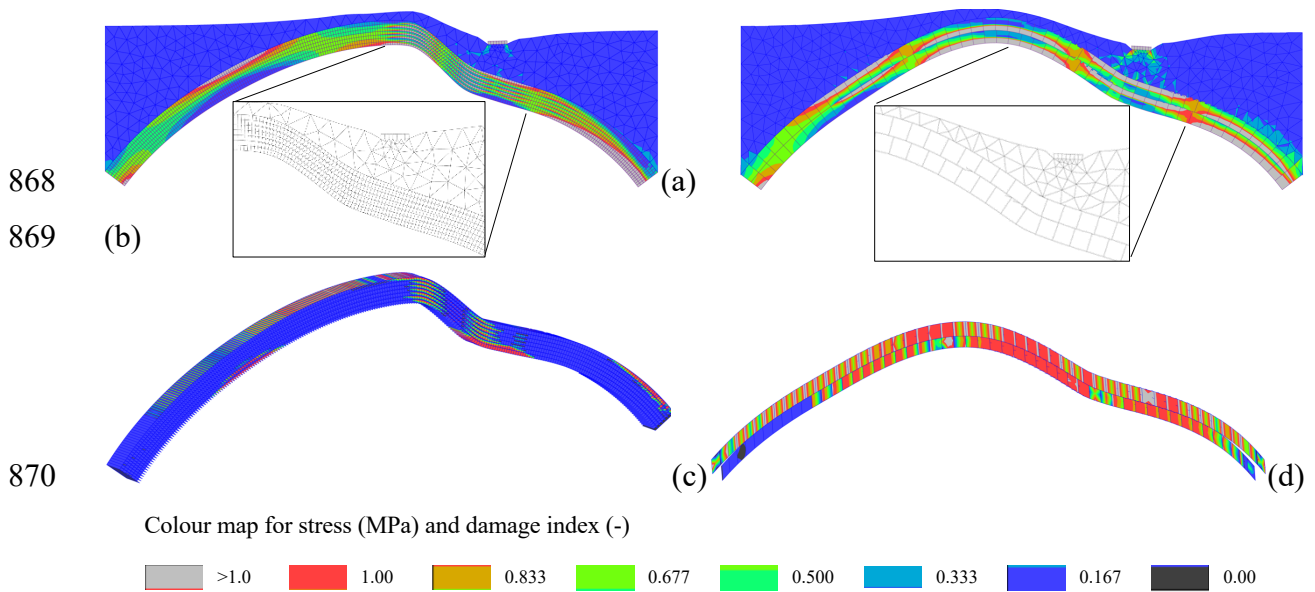
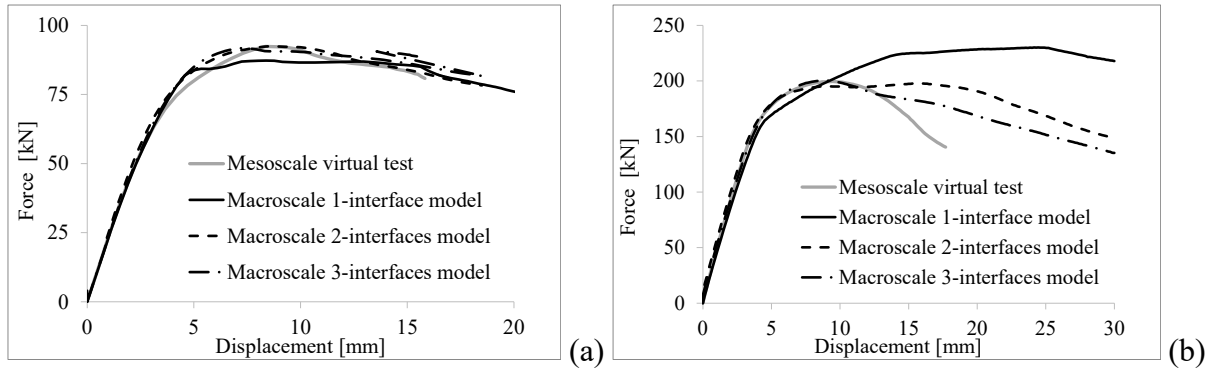


Figure 24. Ultimate deformed shape (amplificated 30 times) of the arch interacting with backfill with von-Mises stresses in the solid elements of the (a) mesoscale model and (b) the macroscale model; and damage in tension contours at the (c) mesoscale interfaces and (d) the nonlinear interfaces of the macroscale hybrid model.

875 Figure 25 shows the load-displacement mesoscale and macroscale curves for the bare arch and  
 876 the bridge specimen. It can be seen that each calibrated macroscale model reproduces very well  
 877 the arch response as determined by the baseline mesoscale model (Figure 25a). Furthermore,  
 878 the models with two or three circumferential interface layers (*2-Interface* and *3-Interface*  
 879 models) lead to improved predictions in terms of ultimate load and post-peak response  
 880 compared to the model with a single interface layer.

881



882

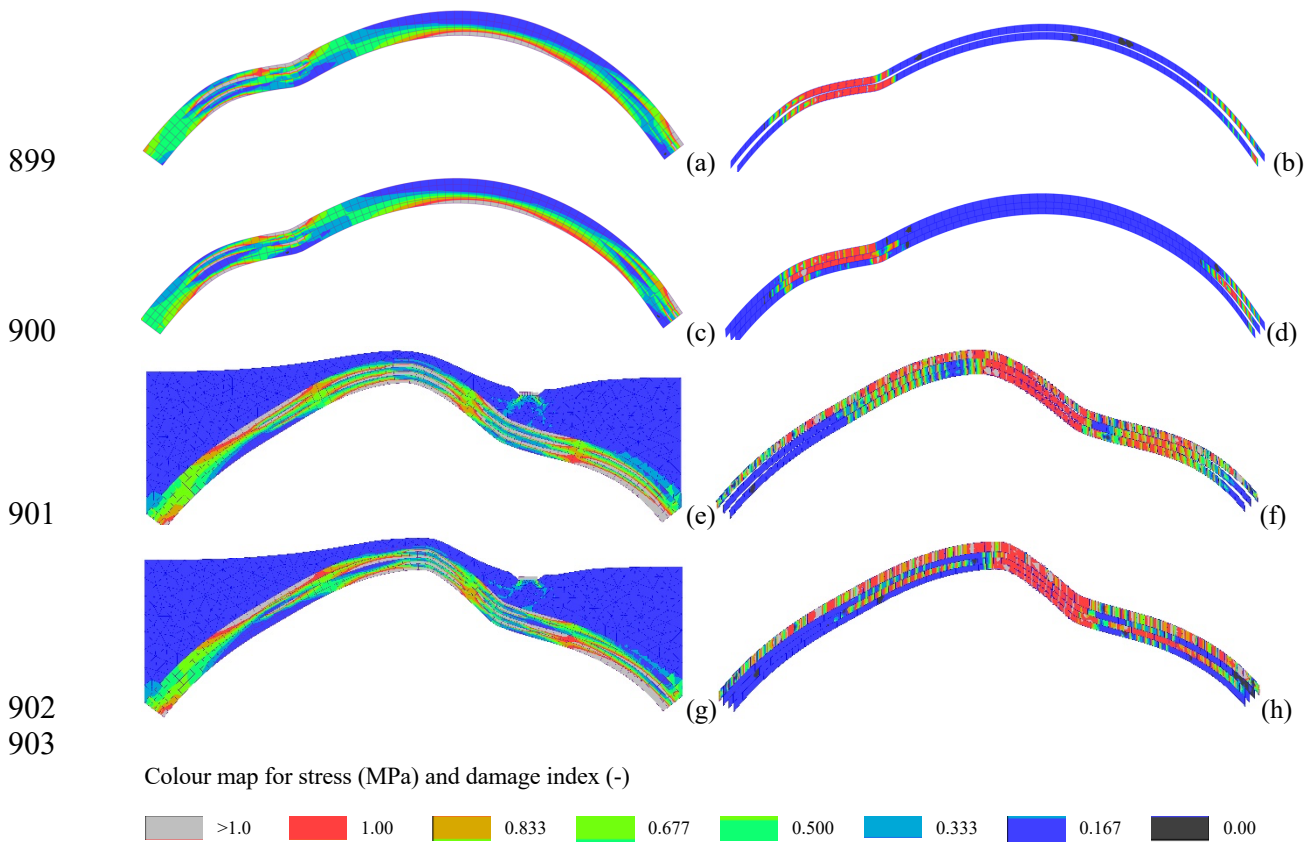
883 Figure 25. Load-displacement curves of the hybrid models with single and multiple interfaces: (a) bare

884

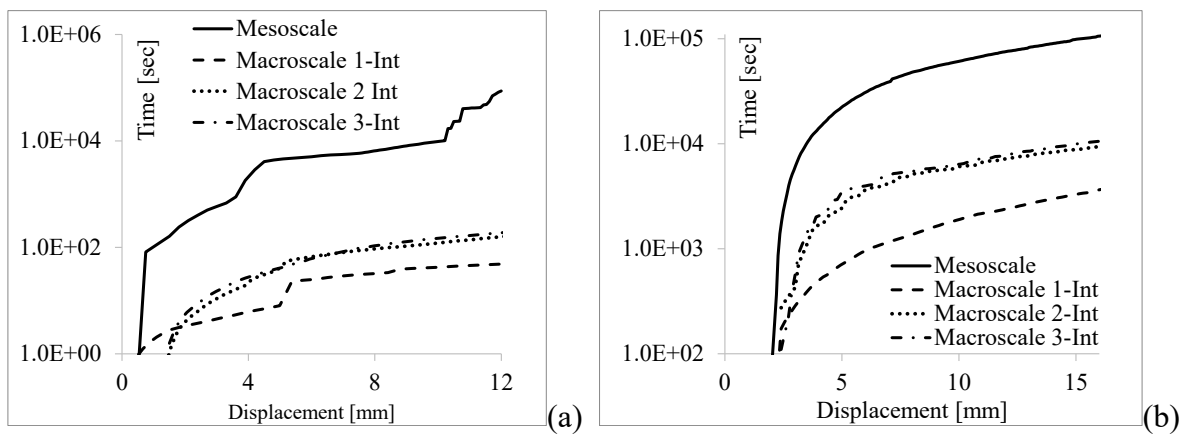
arch virtual test and (b) arch bridge specimen.

885 The failure mechanisms depicted in Figure 26 confirm that the use of multiple interface layers  
 886 allows a more realistic representation of the distributed shear-sliding failure mode, where no  
 887 notable differences between *2-Interface* and *3-Interface* models can be observed.

888 Finally, the efficiency of the proposed hybrid modelling strategy is evaluated against the  
 889 mesoscale results. In the case of the bare arch, the computing times required by the macroscale  
 890 *1-Interface*, *2-Interface* and *3-Interface* models are respectively 0.55%, 0.17% and 0.22% of  
 891 the wall clock time required by the mesoscale model (Figure 27a). In the case of the arch bridge  
 892 specimen, the corresponding computing times of the hybrid models are respectively 3.4%,  
 893 9.8% and 10.0% of the mesoscale time. These results confirm the much enhanced efficiency  
 894 guaranteed by the proposed modelling strategy for realistic simulations of large arch bridges,  
 895 where the use of the mesoscale modelling approach may become computationally prohibitive.  
 896 In the main, the introduction of multiple interface layers does not lead to a significant increase  
 897 of the computing time which is, for the analysed cases, always less than 10% of that required  
 898 by the mesoscale model.



904 Figure 26. Deformed shape at collapse with von-Mises stresses in the solid elements and in  
 905 tension contours on the interface elements for the bare arch (a,b) *2-Interface* and (c,d) *3-Interface*  
 906 models, and for the (e,f) *2-Interface* and (g,h) the *3-Interface* models for the arch bridge specimen.



908 Figure 27. Computing time for the simulation of the (a) bare arch and (b) the arch interacting with  
 909 backfill.

## 910 6 CONCLUSIONS

911 In this study, a hybrid continuum-discrete macro-modelling description for brick-masonry  
912 multi-ring arches and arch bridges is proposed. According to this modelling approach, the arch  
913 and backfill domain are modelled by 3D continuum solid elements, while layers of 2D zero-  
914 thickness nonlinear interfaces arranged along the circumferential direction of the arch simulate  
915 potential ring separation and the interaction between the arch and backfill.

916 Two advanced damage-plasticity constitutive models are employed for the 3D solid and  
917 interface elements. An effective and robust multi-level calibration procedure, based on  
918 minimisation of stress power error solved by means of Genetic Algorithms is developed to  
919 evaluate the model mechanical parameters employing the results from detailed mesoscale  
920 models on virtual experiments. These parameters can be easily calibrated from non-destructive  
921 or low-destructive in-situ tests on masonry units and mortar joints, which renders the proposed  
922 calibration procedure suitable also for practical assessment of historical bridges.

923 The accuracy and potential of the proposed modelling strategy and calibration procedure is  
924 demonstrated by analysing 2D-strip medium and long span masonry arch specimens, also  
925 interacting with backfill and characterised by different failure mechanisms. The results of the  
926 proposed hybrid model are compared to those obtained by detailed mesoscale and continuum  
927 finite element macroscale descriptions. It has been found that the proposed modelling strategy  
928 provides accurate predictions of the ultimate strength and displacement capacity of multi-ring  
929 masonry arches and the corresponding failure mechanisms. It allows for potential shear sliding  
930 between adjacent rings, where the use of only one circumferential interface layer is suitable for  
931 medium span arches but at least two layers are required for long span arches and bridges.  
932 Additionally, the proposed modelling strategy guarantees superior computational efficiency  
933 especially for the analysis of long span arches and bridges, where the use of detailed mesoscale  
934 modelling can become infeasible.

935 Importantly, the numerical results identified some drawbacks associated with the use of  
936 conventional isotropic finite element macromodels, which can be summarised as following:

- 937 - The use of continuum finite element macromodels, without a rigorous calibration  
938 of the mechanical parameters, can lead to an inaccurate and non-objective  
939 prediction of the arch response.
- 940 - Continuum finite element macromodels, even when calibrated by means of rigorous  
941 procedures, can fail in simulating the ultimate arch behaviour when it is driven by  
942 sliding between adjacent rings.

943 Both these limitations may significantly affect the results of safety assessments of masonry  
944 arch bridges, leading to a crudely approximated or completely misleading prediction of the  
945 effective safety level of the bridge and its mode of failure. In this regard, the proposed hybrid  
946 modelling strategy offers the possibility to significantly improve the accuracy of the numerical  
947 predictions without requiring a significant increase in the computational effort associated with  
948 the nonlinear analysis.

949

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