Doubly Damped Stochastic Parallel Translations and Hessian Formulas

Xue-Mei Li

Abstract We study the Hessian of the solutions of time-independent Schrödinger equations, aiming to obtain as large a class as possible of complete Riemannian manifolds for which the estimate $C(\frac{1}{l} + \frac{d^2}{l^2})$ holds. For this purpose we introduce the doubly damped stochastic parallel transport equation, study them and make exponential estimates on them, deduce a second order Feynman-Kac formula and obtain the desired estimates. Our aim here is to explain the intuition, the basic techniques, and the formulas which might be useful in other studies.

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1 Introduction

The probability distribution of a Brownian motion or a Brownian bridge are reference measures with which we make L^2 analysis on the space of continuous paths (the Wiener space) and its subspaces of the pinned paths. On the Wiener space, these are Gaussian measures and are well understood. The theory of the probability distributions of Brownian motion and Brownian bridges on more general manifolds is less developed. These include elliptic and semi-elliptic diffusion in an Euclidean space, with non-constant coefficients. For example we would like to describe the tail behaviour of the measure, but how do we describe the set of path far away? Instead, we measure the size of the tails by checking whether a Lipschitz continuous function f is exponentially integrable, or whether $\mathbf{E}(e^{cf^2})$ is finite for a constant c. In fact, a theorem of Herbst states that if a probability measure μ on \mathbf{R}^n satis-

Department of Mathematics, Imperial College London, London SW7 2AZ. e-mail: xue-mei.li@Imperial.ac.uk. We dedicate this paper to Michael Röckner on the occasion of his 60's birthday.

fies the following logarithmic Sobolev inequality $\mathbf{E}(f^2\log(f^2)) \leq c_0\mathbf{E}|\nabla f|^2$, then $\mathbf{E}(e^{\varepsilon f^2}) \leq e^{\frac{\varepsilon}{(1-c_0\varepsilon)}}$ for any smooth function $f: \mathbf{R}^n \to \mathbf{R}$ with $\|f\|_{L^2(\mu)} = 1$ where ε is any number in $(0,1/c_0)$ and. Also if a mean zero function f is Lipschitz continuous with $\|f\|_{\mathrm{Lip}} = 1$, then $\mathbf{E}(e^{\alpha f}) \leq e^{c_0\alpha^2}$ holds for any number α , On a path or loop space, similar results hold [1]. There are many other applications of Logarithmic Sobolev inequalities, see e.g. M. Ledoux's Saint Flour notes [12] and the reference therein. To obtain such functional inequalities we make use of estimates on the fundamental solutions of Kolmogorov equations.

Let M denote a connected smooth manifold with a complete Riemannian metric g. Denote by (g^{ij}) the inverse of the Riemannian metric $g = (g_{ij})$. There exists on M a strong Markov process with Markov generator $\frac{1}{2}\Delta$ where Δ is the Laplace-Beltrami operator which in local coordinates takes the form

$$\Delta f(x) = \frac{1}{\sqrt{\det g(x)}} \partial_i \left(\sqrt{\det g} \, g^{ij} \partial_j f \right)(x).$$

This stochastic process is said to be a Brownian motion. If Z is a vector field we denote by L_Z Lie differentiation in the direction of Z so for a real valued function f, $L_Z f = df(Z)$. Observe that any second order elliptic differential operator is of the form $\frac{1}{2}\Delta + L_Z$ where Δ is the Laplacian for the Riemannian metric induced by the operator, so they are generators of Brownian motions (with possibly a nonzero drift). In this article we are mainly concerned with gradient drifts, $Z = 2\nabla h$ where $h: M \to \mathbf{R}$ is a smooth function. In coordinates, the operators are of the form $\sum_{i,j} a_{i,j}(x) \frac{\partial^2}{\partial x_i \partial x_j} + \sum_k \frac{\partial}{\partial x_k}$, which is (locally) elliptic and is in general not strictly elliptic.

Set $\Delta^h = \Delta + 2L_{\nabla h}$, this is called the Bismut-Witten or the weighted Laplacian. With respect to the weighted volume measure $e^{2h}dx$, Δ^h is like a Laplacian. In particular, $\Delta^h = -(d+\delta^h)^2$ where δ^h is the adjoint of d on $L^2(e^{2h}dx)$. All three operators, d, δ^h and Δ^h , extend to acting on differential forms. Then $d+\delta^h$ and all its powers are essentially self-adjoint on C_K^∞ , the space of smooth compactly supported differential forms, for the details see [14]. The densities of the probability distributions of the weighted Brownian motion are the weighted heat kernel. These are the fundamental solution to the equation $\frac{\partial}{\partial t} = \frac{1}{2}\Delta^h$. There is also a commutative relation with which one can obtain gradient estimates for the weighted heat kernels under conditions on the Ricci curvature (without involving their derivatives).

We introduce the notations. Let $\operatorname{Ric}_x : T_x M \times T_x M \to \mathbf{R}$ denote the Ricci curvature and let \mathscr{R} denote the curvature tensor. Let $\operatorname{Ric}_x^{\sharp} : T_x M \to T_x M$ denote the linear map defined by the relation: $\langle \operatorname{Ric}_x^{\sharp}(u), v \rangle = \operatorname{Ric}_x(u, v)$. One of the novelties is to introduce the symmetrised tensor Θ , see [13],

$$\langle \Theta(v_2)v_1, v_3 \rangle = \left(\nabla_{v_3} \operatorname{Ric}^{\sharp}\right) (v_1, v_2) - \left(\nabla_{v_1} \operatorname{Ric}^{\sharp}\right) (v_3, v_2) - \left(\nabla_{v_2} \operatorname{Ric}^{\sharp}\right) (v_1, v_3), (1)$$

where $v_1, v_2, v_3 \in T_{x_0}M$, and to impose growth conditions on a bilinear map Θ^h from $T_{x_0}M \times T_{x_0}M$ to **R** instead of imposing conditions on $|\nabla \operatorname{Ric}^{\sharp}|$. The bilinear form is

defined by the formula

$$\Theta^{h}(v_{2})(v_{1}) = \frac{1}{2}\Theta(v_{2}, v_{1}) + \nabla^{2}(\nabla h)(v_{2}, v_{1}) + \mathcal{R}(\nabla h, v_{2})(v_{1}), \qquad v_{1}, v_{2} \in T_{x_{0}}M.$$

We are particularly interested in adding a zero order potential and consider a time independent Schrödinger equation, which is a parabolic partial differential equation of the form

$$\frac{\partial u}{\partial t} = (\frac{1}{2}\Delta + L_{\nabla h} + V)u,\tag{2}$$

where $u : \mathbf{R} \times M \to \mathbf{R}$ is a real valued function. For simplicity the zero order potential function $V : M \to \mathbf{R}$ will be assumed to be bounded and Hölder continuous and so $\Delta + 2L_{\nabla h} + V$ is essentially self-adjoint on $C_K^{\infty} \subset L^2(M, e^{2h}dx)$, the space of smooth functions with compact supports.

Our objectives are to obtain global estimates for the Schrödinger semi-group $e^{(\frac{1}{2}\Delta^h+V)t}f$ where f is a bounded measurable function, for its gradient, and for its second order derivatives in terms of the geometric data of the Riemannian manifolds. We will be also interested in such estimates for its fundamental solutions, which we denote by $p^{h,V}(t,x,y,t)$, or $P^V(t,x,y)$ if h vanishes identically or $p^h(t,x,y)$ if V vanishes identically, and V vanishes identically, and V vanishes identically, with capital V, e.g. $V^{h,V}_t$, will be used to denote the corresponding semi-groups.

The commutative relation we mentioned earlier is as follows: the differential d and the semi-group $e^{\frac{1}{2}\Delta^h}$ commute on C_K^{∞} , and consequently $de^{\frac{1}{2}\Delta^h}$ solves the heat equation on differential 1-forms: $\frac{\partial}{\partial t}\phi = \frac{1}{2}\Delta^h\phi$. If $M = \mathbf{R}^n$ this equation on differential 1-forms is an equation on 'vector-valued' functions. Let us denote by $\nabla^{h,*}$ the adjoint of ∇ on $L^2(e^{2h}dx)$, then this equation becomes:

$$\frac{\partial}{\partial t}\phi = \frac{1}{2}\nabla^{h,*}\nabla\phi - \frac{1}{2}\phi(\operatorname{Ric}^{\sharp} - 2\nabla\nabla h). \tag{3}$$

To see this we observe that, if ∇^* is the adjoint of ∇ on $L^2(dx)$, then there is the Weitzenbock formula $\Delta^h = -\nabla^*\nabla\phi - \mathrm{Ric}^\sharp(\phi) + 2L_{\nabla h}\phi$, $L_{\nabla h}\phi = \nabla\phi(\nabla h)d(\phi(\nabla h) + d\phi(\nabla h, \cdot))$ where $\iota_{\nabla h}$ denotes the interior product. Also, for a differential 1-form ϕ we apply the identity

$$\nabla^{h,*}\phi = \nabla^*\phi - 2\iota_{\nabla h}\phi$$

to see that

$$\nabla^{h,*}\nabla_{\cdot}\phi = \nabla^*\nabla_{\cdot}\phi - 2\imath_{\nabla h}\nabla_{\cdot}\phi$$

Equation (3) inspired the study of the damped stochastic parallel translation

$$W_t: T_{x_0}M \to T_{x_t(\boldsymbol{\omega})}M$$

along a path $x_l(\omega)$ which solves the stochastic damped parallel translation equation

$$\frac{DW_t}{dt} = -\frac{1}{2} \operatorname{Ric}_{x_t}^{\sharp}(W_t) + \nabla_{W_t} \nabla h, \quad W_0 = Id$$
 (4)

Here Id denotes the identity map on $T_{x_0}M$ and

$$\frac{DW_t}{dt} := //_t(x_{\cdot}(\boldsymbol{\omega})) \frac{d}{dt} \left(//_t^{-1}(x_{\cdot}(\boldsymbol{\omega})) W_t \right)$$

is the covariant derivative along $x_t(\omega)$ and $//_t(x_t(\omega)): T_{x_0}M \to T_{x_t(\omega)}M$ denotes the standard stochastic parallel translation and $//_t^{-1}(x_t(\omega)): T_{x_t(\omega)}M \to T_{x_0}M$ is its inverse.

Stochastic parallel translations along the non-differentiable sample paths of a Brownian motion can be constructed by a stochastic differential equation on the orthonormal frame bundle. This goes back to J. Eells, K. D. Elworthy and P. Malliavin, earlier attempts go back to K. Itô and M. Pinsky. M. Emery and M. Arnaudon studied parallel translations along a general semi-martingale [9]. The damped parallel translation goes back to E. Airault [2]. The damped stochastic parallel translation takes into accounts of the effect of the Ricci curvature along its path and unwind it, leading to the magic well known formula: $de^{\frac{1}{2}\Delta^h_t}f(v) = \mathbf{E}df(W_t(v))$ for (x_t) a Brownian motion with the initial value x_0 . This holds for all compact manifolds and for more general manifolds.

The global estimates we are after are of the form

$$\left|\nabla^2 p(t, x_0, y_0)\right| \le C\left(\frac{1}{t} + \frac{d^2(x_0, y_0)}{t^2}\right), \quad t \in (0, 1], \ x, y \in M.$$
 (5)

Such estimates (for h=0, V=0 and for compact manifolds) were obtained in [23] and were generalised to other types of manifolds we refer to the references in [13]. We should remark that adequate care must be taken when generalising estimates from compact manifolds to non-compact manifolds. For example taking a localising sequence of stopping times may not come for free and any technique involving differentiating a stochastic flow with respect to its initial point will likely need the additional assumption the strong 1-completeness [15]. See also [4], [22], [25], [10], and [24].

Our goal is to establish these estimates for as large a class of manifolds as possible and extend them to the operators $\frac{1}{2}\Delta + L_{\nabla h} + V$. If both V and h vanish, these estimates are relevant for the study of the space of continuous loops and pinned paths using the Brownian bridge measure, e.g. the probability measure induced by a Brownian motion conditioned to return to a point y_0 at time 1. Naturally the Brownian motions with the symmetric drift ∇h , which we refer as an h-Brownian motion, are also candidates for such studies. We recall, that the h-Brownian bridge is a Markov process x_t on [0,1) with the Markov generator

$$\frac{1}{2}\Delta^h + \nabla \log p(1-t, x, y_0),$$

and $\lim_{t\to 1} x_t = y_0$ where y_0 is the terminal value. Observe that the corresponding damped parallel translation would be

$$\frac{DW_t}{dt} = -\frac{1}{2} \operatorname{Ric}_{x_t}^{\sharp}(W_t) + \nabla_{W_t} \nabla h + \nabla_{W_t} \nabla \log p(1-t, x_t, y_0), \quad t < 1.$$

Gradient estimates on the semi-group associated with the Brownian bridge will naturally involve the second order derivative of $\log p(1-t,x,y_0)$. It is clear that the small time asymptotics of the Hessian are relevant, and estimates of the type (5) appear to be essential for analysing the Brownian bridge measure and useful for the L^2 analysis of loop spaces. In this paper we explain the main formulas and constructions rom [13] that leads to these estimates.

2 Summary of Results

The following is summary of some results from [13].

- (a) We extend estimate (5) to a more general class of manifolds replacing the linear growth condition on $|\nabla \text{Ric}|_{op}$ by a linear growth condition on Θ , Θ being a symmetrised tensor obtained from $\nabla \text{Ric}^{\sharp}$ after taking into accounts of the effects of ∇h defined (1).
- (b) Our proof is based on an elementary Hessian formula which will then lead to an integration by parts type Hessian formula. For these formulas we introduce a doubly damped stochastic parallel transport equation which is defined using Θ . It is natural to call these solutions 'doubly damped stochastic parallel translations'. We denote the solutions by $W_t^{(2)}$, see Lemma 1. We have the formula:

$$\operatorname{Hess}(P_t^h f)(v_2, v_1) = \mathbf{E}\left[\nabla df(W_t(v_2), W_t(v_1))\right] + \mathbf{E}\left[df\left(W_t^{(2)}(v_1, v_2)\right)\right].$$

Here W_t denotes the damped stochastic parallel translation defined by (4). The Second Order Feynman-Kac Formula which does not involve the derivative of f is given in Theorem 1.

- (c) Such estimates will be also extended to the symmetric operators $\frac{1}{2}\Delta^h$ and $\frac{1}{2}\Delta^h + V$. Both operators are essentially self-adjoint on C_K^{∞} . By unitary transformations the drift term and the potential term can be treated almost exchangeably, however the drift and the zero order term do behave differently. For example we assume that h is smooth and pose no direct assumptions on its its growth at the infinity while the zero order term V is only Hölder continuous and is assumed to be bounded. We obtain a second order Feynman-Kac formula, see Theorem 1, the Hölder continuity of V is needed and is used to offset singularities in some of the integrals of the formulas. The modified doubly damped equation involves Θ^h instead of Θ itself.
- (d) These estimates are refined for a subclass of manifolds with a pole, for which we make use of and obtain some nice estimates in terms of the semi-classical Brownian bridges, a more careful study of the semi-classical Brownian bridge measure can be found in [17]. See also [18] for generalised Brownian bridges and [20] for gradient estimates.

3 Key Ingredient

Let X(e) be smooth vector field on M given by an isometric embedding $\phi: M \to \mathbf{R}^m$ and so X(e) is the gradient of the real valued function $\langle \phi, e \rangle$ where $e \in \mathbf{R}^m$. If $\{e_i\}$ is an o.n.b. of \mathbf{R}^m , this induces a family of vector fields $X_i(x)$ where $X_i(x) = X(x)(e_i)$. Let $X_i(x) = X_i(x)(e_i)$.

$$dx_t = \sum_{i=1}^m X_i(x_t) \circ dB_t^i + \nabla h(x_t) dt = X(x_t) \circ dB_t + \nabla h(x_t) dt$$

where \circ denotes Stratonovich integration and B_t , $B_t = (B_t^1, \dots, B_t^m)$, is an \mathbb{R}^m -valued Brownian motion on a filtered probability space with the usual assumptions. Then $F_t(x_0)$, the solution with the initial value $x_0 \in M$, is a Brownian motion with the initial value x_0 .

If x_t is a semi-martingale, the stochastic damped parallel translation $//t(x(\omega))$ along $x_t(\omega)$, which is also denoted by //t, allows us to bring a vector in the tangent space of a solution path at time 0 to its tangent space at time t, to differentiate it there and to bring it back to time t by the inverse parallel translation $//t^{-1}$. If (x_t) is a Brownian motion with the initial value x_0 , the damped stochastic parallel translation W_t along its sample paths, where

$$rac{DW_t}{dt} = -rac{1}{2}\mathrm{Ric}_{x_t}^\sharp(W_t) +
abla_{W_t}
abla h, \quad W_0 = Id,$$

compensates the effect of the Ricci curvature in equation (3) and unwind it, leading to the magic well known formula,

$$de^{\frac{1}{2}\Delta^{h_t}}f(v) = \mathbf{E}\left[df(W_t(v))\right],$$

which holds trivially for compact manifolds and for manifolds with Ric - 2Hess h bounded from below and for more general manifolds.

For the second order derivatives of the fundamental solution of the heat kernel, we ought to differentiate W_t with respect to its initial data, i.e. we differentiate $//(t^{-1}(x))W_t(x_0)$ which is a map from M to the space of linear maps which we denote by $\mathcal{L}(T_{x_0}M;T_{x_0}M)$.

We introduce the doubly damped stochastic parallel translation equation, whose solution we call doubly damped stochastic parallel transports/translations. Unlike damped parallel translations, the doubly damped ones involve genuine stochastic integrals (unless the curvature vanishes) and it is a challenge to obtain exponential estimates. We also recall that the damped parallel translations are conditional expectations of the spatial derivative of the solution to the gradient SDE. The doubly damped ones are obtained by differentiating the damped parallel translations, followed by taking conditional expectations. The beauty of it is that it satisfies the doubly damped stochastic parallel translation equation:

$$Dv_{t} = \left(-\frac{1}{2}\operatorname{Ric} + \operatorname{Hess} h\right)^{\sharp}(v_{t}) dt + \frac{1}{2}\Theta^{h}(W_{t}(v_{2}))(W_{t}(v_{1})) dt + \mathcal{R}(d\{x_{t}\}, W_{t}(v_{2}))W_{t}(v_{1}).$$
(6)

We also introduce the notation $d\{x_t\}$, by which we mean integration with respect to the martingale part of $\{x_t\}$, see [8, sec. 4.1] for detail. This allows us to give statements on a h-Brownian motion that is independent of its representation as a solution to a specific stochastic differential equation. In particular we may use any of the two canonical representations: (1) $x_t = \pi(u_t)$ where u_t is the solution to the canonical SDE on the orthonormal frame bundle

$$du_t = H(u_t) \circ dB_t + \mathfrak{h}_{u_t} (\nabla h(\pi(u_t))) dt$$

where \mathfrak{h}_u denotes the horizontal lift map at a given frame u, and $\pi: OM \to M$ takes a frame, a point of OM, to its base point. (2) x_t is the solution of a gradient SDE. If u_t is the solution to the canonical SDE on the orthonormal frame bundle, then $d\{x_t\}$ is interpreted as $u_t dB_t$. If (x_t) is the solution to to a gradient SDE driven by X then $d\{x_t\}$ is interpreted as $X(x_t) dB_t$.

Lemma 1. Suppose that the gradient SDE is strongly 1-complete and suppose that $v_1, v_2 \in T_{x_t}M$ and $x_0 \in M$ and (x_t) is the solution to the gradient SDE. Let $W_t^{(2)}(v_1, v_2)$ denote the solution to the following covariant differential equation (the doubly damped stochastic parallel translation equation):

$$Dv_{t} = \left(-\frac{1}{2}\operatorname{Ric} + \operatorname{Hess} h\right)^{\sharp}(v_{t}) dt + \frac{1}{2}\Theta^{h}(W_{t}(v_{2}))(W_{t}(v_{1})) dt + \mathcal{R}(d\{x_{t}\}, W_{t}(v_{2}))W_{t}(v_{1}),$$

$$v_{0} = 0.$$

Then $W_t^{(2)}(v_1, v_2)$ is the local conditional expectation of $\nabla_{v_2} W_t(v_1)$ with respect to the filtration $\mathscr{F}_t^{x_0} := \sigma\{x_s : s \leq t\}$. If furthermore the latter is integrable, then

$$W_t^{(2)}(v_1,v_2) = \mathbf{E}\left\{\nabla_{v_2}W_t(v_1) \mid \mathscr{F}_t^{x_0}\right\}.$$

The proof for the lemma consists of stochastic calculus involving $\frac{D}{ds}W_t(j(s))$ where j is a parallel field with $j(0) = v_2$, along the normalised geodesic γ with the initial condition x_0 and the initial velocity $\dot{\gamma}(0) = v_1$. Observe also that

$$\nabla_{v_2} W_t(v_1) = \frac{D}{ds}|_{s=0} W_t(j(s)).$$

Strong 1-completeness of an SDE is a concept that is weaker than strong completeness, by the latter we mean the existence of a global solution to the SDE which is continuous with respect to the initial value. Let p = 1, 2, ..., n where n is the dimension of the manifold. Roughly speaking, an SDE is strongly p-complete if for

a.s. every ω , and for all t, $F_t(x,\omega)$ is continuous with respect to the initial point x when x is restricted to a sub-manifold of dimension p (or to a smooth C^1 curve if p=1). The first example of an SDE which is complete (i.e. its solution from any initial point has infinite life time) and which is not strongly complete was given by K.D. Elworthy, prior to which it was generally believed that the two problems are equivalent. The concept of strong p-completeness was introduced in [14, 15] where we also give examples of strongly p-1-complete SDEs which are not strongly p-complete and n > 2 and $p \le n$. For n = 1 completeness is equivalent to strong completeness, similarly for n = 2, strong completeness is equivalent to strong completeness. In [19], a non-strongly complete SDE on \mathbb{R}^2 is given: it has one single driving linear Brownian motion and is driven a smooth bounded driving vector field. We emphasise that, due to the fact that the exit time of $F_t(x, \omega)$ from a geodesic ball (even one with smooth boundary) is not necessarily continuous with respect to the initial point x, and it is not trivial to solve the strong 1-completeness by localisation. The strong 1-completeness for gradient SDEs was specially studied in [13]. See also the books [5] and [11].

Remark 1. If the gradient SDE is strongly 1-complete, $s \mapsto |W_t(\dot{\gamma}(s))|$ is continuous in L^1 and $\mathbf{E}|T_{\gamma(s)}F_t|$ is finite, we know that for all $f \in BC^1$, $d(P_t^h f)(v_1) = \mathbf{E}df(W_t(v_1))$, [13, 15, 6]. From this we see immediately that

$$|d(P_t f)|_{L_\infty} \leq |df|_\infty \mathbf{E}\left(e^{\int_0^t \rho^h(x_s)ds}\right),$$

where $\rho^h(x) = \sup_{|v|=1, v \in T_x M} \{-\frac{1}{2}\operatorname{Ric}(v, v) + Hess(h)(v, v)\}$. A more relaxed condition for this to hold can be obtained, but most of the assumptions here will be needed later. If ρ^h is bounded by -K then we see immediately on direction of the characterisation for the Ricci curvature to be bounded below by K, by taking h=0 in the earlier estimate,

$$|d(P_t f)|_{I_{\infty}} < |df|_{\infty} e^{-Kt}$$
.

We give below the second order analogue. Denote by $T_{x_0}F_t(v_0)$ the derivative flow of $F_t(x)$, it solves the equation

$$dV_t = (\nabla X)_{x_t}(V_t) \circ dB_t + (\nabla^2 h)_{x_t}(V_t)dt,$$

$$V_0 = v_0.$$

Useful moment estimates on the derivatives flows for non-compact manifolds can be found in [16]. Recall that j(s) is a parallel field along the geodesic γ with $\dot{\gamma}(0) = v_1$. with the initial value $j(0) = v_2$.

Lemma 2. Suppose that Ric - 2Hess(h) is bounded from below and that the gradient SDE is strongly 1-complete. Suppose also the statements (a) and (b) below hold.

(a) for every s, $\mathbf{E}|T_{\gamma(s)}F_t|$ and $\mathbf{E}|\nabla_{TF_t(\gamma(s))}W_t|$ are finite.

(b)
$$s \mapsto \mathbb{E}\left\{\frac{D}{ds}W_t(j(s)) \middle| \mathscr{F}_t^{\gamma(s)}\right\}$$
 is continuous in $L^1(\Omega)$;

Then for all $f \in BC^2$,

$$\operatorname{Hess}(P_t^h f)(v_2, v_1) = \mathbf{E}\left[\nabla df(W_t(v_2), W_t(v_1))\right] + \mathbf{E}\left[df\left(W_t^{(2)}(v_1, v_2)\right)\right]. \tag{7}$$

From this lemma we immediately obtain the following estimate:

$$\left| \operatorname{Hess}(P_t^h f) \right|_{L_{\infty}} \leq |\nabla df|_{\infty} \operatorname{\mathbf{E}}\left(e^{2\int_0^t \rho^h(x_s) ds} \right) + |df|_{L_{\infty}} \operatorname{\mathbf{E}}\left| W_t^{(2)} \right|.$$

It is clear that estimation on $\mathbf{E} \left| W_t^{(2)} \right|$ will be useful, this is given in [13] and which we do not include here. As we shall see, to obtain Hessian estimates of the form (5), we will need to obtain exponential integrability of $\left| W_t^{(2)} \right|^2$, such estimates will be given shortly after we explain why this is so.

The following are the basic assumptions. [C1.]

- (a) Ric -2 Hess $(h) \ge -K$;
- (b) $\sup_{s\leq t} \mathbf{E}(\|W_s^{(2)}\|^2) < \infty;$
- (c) for all $f \in BC^2(M; \mathbf{R})$, $v_1, v_2 \in T_{x_0}M$, the elementary Hessian formula (7) holds.

For the Schrödinger equation we first set, for r < t and $x_0 \in M$,

$$\mathbb{V}_{t-r,t} = (V(x_{t-r}) - V(x_0))e^{-\int_{t-r}^{t} [V(x_s) - V(x_0)]ds}.$$
 (8)

Set also.

$$N_t = \frac{4}{t^2} \int_{\frac{t}{2}}^t \langle X(x_s) dB_s, W_s(v_1) \rangle \int_0^{\frac{t}{2}} \langle X(x_s) dB_s, W_s(v_2) \rangle. \tag{9}$$

Theorem 1 (Second Order Feynman-Kac Formula). Suppose that $\underline{C1}$ holds. Let V be a bounded Hölder continuous function. Then for any $f \in \mathcal{B}_b(M; \mathbf{R})$,

$$\operatorname{Hess} P_{t}^{h,V} f(v_{1}, v_{2}) = e^{-V(x_{0})t} \mathbf{E} \left[f(x_{t}) N_{t} \right] + e^{-V(x_{0})t} \mathbf{E} \left[f(x_{t}) \frac{2}{t} \int_{0}^{t/2} \langle X(x_{s}) dB_{s}, W_{s}^{(2)}(v_{1}, v_{2}) \rangle \right]$$

$$+ e^{-V(x_{0})t} \int_{0}^{t} \mathbf{E} \left[f(x_{t}) \frac{2 \mathbb{V}_{t-r,t}}{t-r} \int_{0}^{(t-r)/2} \langle X(x_{s}) dB_{s}, W_{s}^{(2)}(v_{1}, v_{2}) \rangle \right] dr$$

$$+ e^{-V(x_{0})t} \int_{0}^{t} \mathbf{E} \left[f(x_{t}) \mathbb{V}_{t-r,t} N_{t-r} \right] dr.$$

$$(10)$$

For h = V = 0, a version of the Hessian formula was first given in [6] followed by another in [3]. A version of the Hessian formula for $h \equiv 0$ and $V \neq 0$ was also given in [6], however no proof was given. The doubly damped stochastic parallel translation equations were not present in either papers, nor were any extensive estimates given. Hessian formula and estimates for non-linear potential, on linear space, were given in [21]. A formula for the Laplacian of the semigroup $P_t f$ can be found in [7].

Corollary 1. We assume $V(x_0) = 0$. Then,

$$\begin{split} &\operatorname{Hess} p^{h,V}(t,x_0,y) \\ &= \operatorname{Hess} p_t^h(x_0,y) + \int_0^t \int_M V(z) \operatorname{Hess} p^h(t-r,x_0,z) p_r^h(z,y) \mathbf{E}[e^{-\int_0^r V(Y_s^{r,z,y})}] dz dr, \end{split}$$

where $Y_s^{r,z,y}$ is the h-Brownian bridge with terminal value r, initial value z and terminal value y.

Finally we indicate how to obtain estimates from these formulas. Let us take V = 0 for simplicity, so the formula reads:

$$\operatorname{Hess} P_t^h f(v_1, v_2) = \mathbf{E} \left[f(x_t) N_t \right] + \mathbf{E} \left[f(x_t) \frac{2}{t} \int_0^{t/2} \langle X_i(x_s) dB_s, W_s^{(2)}(v_1, v_2) \rangle \right]. \tag{11}$$

We then choose f(x) to be the fundamental solution $p^h(t,x,y_0)$, so $P_t^h f(x_0,y_0) = p(2t,x_0,y_0)$. In particular,

$$\frac{\operatorname{Hess} p^{h}(2t, x_{0}, y_{0})(v_{1}, v_{2})}{p^{h}(2t, x_{0}, y_{0})} = \mathbf{E} \left[\frac{p(t, x_{t}, y_{0})}{p(2t, x_{0}, y_{0})} N_{t} \right] + \mathbf{E} \left[\frac{p(t, x_{t}, y_{0})}{p(2t, x_{0}, y_{0})} \frac{2}{t} \int_{0}^{t/2} \langle X(x_{s}) dB_{s}, W_{s}^{(2)}(v_{1}, v_{2}) \rangle \right].$$
(12)

The right hand side can then be estimated. Since $|W_t|$ is bounded by a deterministic function (when ρ^h is bounded above), the first term of the right hand side is easier to estimate. Let us work with the second term,

$$\frac{1}{p(2t,x_{0},y_{0})} \mathbf{E} \left(p(t,x_{t},y_{0}) \frac{2}{t} \int_{0}^{t/2} \langle X(x_{s}) dB_{s}, W_{s}^{(2)}(v_{1},v_{2}) \rangle \right) \\
\leq \frac{2}{t} \mathbf{E} \left(\frac{p(t,x_{t},y_{0})}{p(2t,x_{0},y_{0})} \log \frac{p(t,x_{t},y_{0})}{p(2t,x_{0},y_{0})} \right) + \frac{2}{t} \log \mathbf{E} \left(\exp \left(\int_{0}^{t/2} \langle X_{i}(x_{s}) dB_{s}, W_{s}^{(2)}(v_{1},v_{2}) \rangle \right) \right) \\
\leq \frac{2}{t} \sup_{y \in M} \log \frac{p(t,y,y_{0})}{p(2t,x_{0},y_{0})} + \frac{2}{t} \log \mathbf{E} \left(\exp \left(\int_{0}^{t/2} \langle X_{i}(x_{s}) dB_{s}, W_{s}^{(2)}(v_{1},v_{2}) \rangle \right) \right).$$

This can then be refined by heat kernel estimates and by estimates on $\mathbf{E}\left(e^{\left|W_s^{(2)}(\nu_1,\nu_2)\right|^2}\right)$, we illustrate the latter below. The other terms can be treated similarly.

Lemma 3. Suppose that $|\rho^h| \le K$, $||\mathcal{R}_x|| \le ||\mathcal{R}||_{\infty}$, and $||\Theta^h||^2 \le c + \delta r^2$ for δ sufficiently small. Set $C_1(T,0) = 1$,

$$C_1(T,K) = \sup_{0 < s < 3KT} \frac{1}{s} (e^s - 1), \quad \alpha_2(T,K,\|\mathscr{R}\|_{\infty}) = \frac{1}{49n^2 \|\mathscr{R}\|_{\infty}^2 C_1(T,K)}.$$

Then there exists a universal constant c such that for unit vectors $v_1, v_2 \in T_{x_0}M$, and for any $\alpha \leq \alpha_2(T, K, ||\mathcal{R}||_{\infty})$,

$$\begin{split} \mathbf{E} \exp\left(\alpha \gamma |W_t^{(2)}(v_1, v_2)|^2\right) &\leq c e^{2\alpha \gamma} \sqrt{\mathbf{E} \exp\left(4t \ \gamma \alpha \int_0^t e^{3Ks} \|\Theta^h\|_{x_s}^2 ds\right)} \\ &\leq c e^{\frac{2\gamma}{49n^2 \|\mathscr{R}\|_{\infty}^2}} \sqrt{\mathbf{E} \exp\left(\frac{4t \gamma}{49n^2 \|\mathscr{R}\|_{\infty}^2 C_1(t, K)} \int_0^t e^{3Ks} \|\Theta^h\|_{x_s}^2 ds\right)}. \end{split}$$

At this stage we must choose an optimal condition on the growth of $|\Theta^h|$ at the infinity and estimate the exponential integrability of the radial distance function. The condition we imposed is linear growth, with the linear part sufficiently small (or we may compensate the size of the linear part by taking t in a small interval $[0,t_0]$). With these estimates we conclude this paper, and invite the interested reader to consult [13] for technicalities and further results. There we also studied the class of manifolds with a pole using semi-classical bridges. The use of semi-classical bridge for derivatives estimates is novel. See also [17]. Finally we pose the following open question. We know that a bound of the form e^{-Kt} on P_t characterises the lower boundedness of a Ricci curvature. Elworthy asked me whether I can use equation (7) and obtain some characterisation for manifolds. Let me make precise a question here.

Open Problem. For $f \in BC^2$, by Lemma 2,

$$\left| \operatorname{Hess}(P_t^h f) \right|_{I_{\infty}} \leq |\nabla df|_{\infty} \mathbf{E} \left(e^{2 \int_0^t \rho^h(x_s) ds} \right) + |df|_{L_{\infty}} \mathbf{E} \left| W_t^{(2)} \right|.$$

If the Ricci curvature is bounded from below by K, we have

$$\left| \operatorname{Hess}(P_t^h f) \right|_{L_{\infty}} \leq |\nabla df|_{\infty} e^{-2Kt} + |df|_{L_{\infty}} \mathbf{E} \left| W_t^{(2)} \right|.$$

Can we characterise the class of complete Riemannian manifolds, among those whose Ricci curvature is bounded from below by K and whose sectional curvature and symmetrised tensor Θ^h are bounded?

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