Chemical Principles for Electroactive Metal-Organic Frameworks

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Metal-organic frameworks (MOFs) are porous ordered arrays of inorganic clusters connected by organic linkers. The compositional diversity of the metal and ligand, combined with varied connectivity, has yielded over 20,000 unique structures. The application of electronic structure theory can provide deep insights into the fundamental chemistry and physics of these hybrid compounds and identify avenues for design of new multi-functional materials. In this article a number of recent advances in materials modelling of MOFs are reviewed. We present the methodology for predicting the absolute band energies (ionization potentials) of porous solids in comparison to standard semiconductors and electrical contacts. We discuss means of controlling the optical band gaps by chemical modification of the organic and inorganic building blocks. Finally, we outline the principles for achieving electroactive MOFs and outline the key challenges to be addressed in the coming years.

Keywords: metal-organic frameworks, density functional theory, workfunction, bandgap engineering, semiconductor

Introduction

Atomistic materials modelling has become a valuable tool in contemporary materials science. The accelerated characterization of known materials and the assessment of hypothetical systems is being supported by developments in software and hardware, including an international supercomputing infrastructure and a growing number of reliable simulation packages¹. The predictive power of numerical simulation approaches, based on a quantum mechanical description of solids, underpins the emerging field of computational materials design².

Whilst first-principles methodologies, e.g. based on density functional theory (DFT), were once limited to simple structures and compositions with 10s

of atoms in a crystallographic unit cell, modern computer architectures can support the direct simulation of 1000s of atoms. The chemical and physical properties of metal-organic frameworks (MOFs) are now accessible to highquality quantum mechanical simulations. In the past, simple empirical potentials proved useful for screening compositions and topologies in the context of gas storage³. As the interest in MOFs extends towards their physical properties and chemical reactivity, knowledge of electronic structure is essential.

In this article, following a brief description of the common workflow for calculating the physical response functions of hybrid solids, we review recent progress in our understanding of the chemical bonding underpinning the electronic structure of MOFs. This includes the atomic and molecular orbitals that overlap to form the valence and conduction bands, the electron addition and removal energies, and set of design principles for tailoring the electronic and optical activity towards functional devices.

From computer to properties

Input: crystal structure

A reliable structure is an essential starting point for quantum mechanical calculations of crystalline solids. The ground-state distribution of electrons is determined for a particular arrangement of ions in a lattice that extends infinitely across three spatial dimensions. The crystallography of MOFs is challenging: many reported structures contain solvent molecules, lattice sites with partial occupancy or missing hydrogen atoms. In addition, high symmetry space groups are usually assigned in the absence of hydrogen (due to their low electron density and weak diffraction intensity). The symmetry is commonly lowered when hydrogens are included, for example, when carbon atoms are replaced by the appropriate CH, $CH₂$ or CH₃ groups. One useful resource is the CORE-MOF database: http://gregchung.github.io/CoRE-MOFs that provides 'simulation ready' structures and has been employed in high-throughput screening projects⁴.

A further concern, which has not been fully explored for MOFs, is that Bragg diffraction provides insights into the (spatial and temporal) average crystal structure, but the local environment may be quite different⁵. For organic-inorganic

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solids, there is the possibility for vibrations, librations and rotations of atoms or molecular units^{6–9}. Structural disorder has been recognized as a potentially useful trait in framework materials¹⁰, but from the point of view of atomistic simulation it has been largely neglected up to this point.

Output: physical properties

The crystal structures taken from experiment are usually subject to local optimization for a chosen description of the interatomic interactions. The three lattice vectors and all lattice sites are minimized with respect to the external pressure and internal forces on the system. For a modern DFT exchangecorrelation functional (e.g. $PBEso1^{11}$), experimental and measured structural parameters usually agree to within several percent.

On the UK supercomputer ARCHER, a high-quality electronic structure calculation of a MOF with several hundred atoms in the unit cell will require up to 48 hours on 24 compute nodes (1152 cores). The high computational cost of performing a quantum mechanical calculation is balanced by the wealth of information that can be accessed, including structural, mechanical, magnetic and optical properties. For example, heats of formation can be used to screen hypothetical compositions¹², the strength of electron-exchange interactions can be used to predict magnetic critical temperatures¹³, and the composition of the frontier orbitals can be used to explain catalytic activity¹⁴.

Figure 1. Zr-UiO-66, Zn-MOF-5 and Ti-MIL-125 are composed of closed-shell metal oxide clusters (M) linked by benzene dicarboxylate (L). Their topologies are varied, and are determined by the geometry of the inorganic nodes. The electronic properties and frontier orbital compositions are determined by the interface of the inorganic and organic regions. Most MOFs have poor electronic communication between neighboring motifs, resulting in electronic band structures with low dispersion with charge carrier localization. The electronic band structure (density functional theory with HSE06 exchange and correlation) along the high symmetry lines of of Zn-MOF-5 is presented in the lower right panel with the atom-projected electronic density of states.

Electronic structure: more than a sum of parts

A wide variety of metal-organic frameworks have been reported with organic and inorganic networks ranging from $0 - 3$ dimensions of connectivity¹⁵. For example, the hybrid halide perovskites that are being intensively studied for applications in solar energy conversion can be considered as the combination of a 3D anionic inorganic framework with a 0D cationic molecular sublattice¹⁶. For conciseness, we restrict our present discussion to the standard porous MOFs with 3D inorganic-organic-inorganic connectivity, e.g. as found in the ubiquitous MOF-5 that combines a cationic zinc oxide sub-unit connected by anionic benzene dicarboxylate linkers to form a three-dimensional porous framework (see Figure 1).

The orbital composition and spatial location of the valence and conduction bands determines most equilibrium properties of materials. The electronic structure of most MOFs can be described in a manner similar to molecular orbital diagrams common to organometallic chemistry. However, the diverse chemistry of both the organic linkers and the inorganic nodes – for example, whether the node includes an inorganic cation like the Zn_4O^{6+} cluster of the IRMOF series – can require more detailed considerations. In particular, interesting chemistry occurs at the inorganic/organic interface. Whilst studied in detail in the context of surface science and molecular catalysis, the concepts of matching electronic energy levels and orbital symmetry are rarely invoked in MOF chemistry.

An electronic structure calculation can provide insights into the composition, energy and distribution of the frontier extended orbitals (electronic band edges) of any compound. Consider three champion MOFs: Zr-UiO-66,¹⁷ Zn-MOF-5,¹⁸ and Ti-MIL-125.¹⁹ Each framework features a closed-shell s⁰ or d⁰ metal connected by benzene dicarboxylate linkers. The band edges, however, are the product of the chemistry involving the ligand, metal, and their interface as shown in Figure 1. In the case of Zr-UiO-66 and Zn-MOF-5, both the valence and conduction bands are defined by organic orbitals. Thus, organic functionalization can be used to modify the chemistry of the band edges^{20,14} and the physical properties of the material²¹. Alternatively, the chemically inert $ZrO₂$ -based node found in UiO-66 and compositionally similar analogues such as $NU-1000^{22}$ can be used to anchor catalytically active metals, providing access to heterogeneous catalysts²³. Similar metal anchoring is possible through ligand substitutions that boast coordinating functionality (*e.g.* amines, thiols and alcohols)²⁴.

Owing to its wide optical band gap in the UV range ($E_g \sim 4.6$ eV) and high binding energy Zn_4O^{6+} derived bands (-8.1 eV below vacuum level – significantly

deeper than the valence band edge), Zn-MOF-5 offers a further level of chemical modularity allowing for band tuning through both organic funcationalisation, as well as through metal exchange at the Zn-sites. Substitutions for other d-block metals can install mid-gap metal-centered states that result in the MOF featuring a metal-to-ligand transition²⁵. Such cationic substitutions have been proposed in the UiO-type materials, although it remains an open question whether the metals exchange or are anchored to the node.

Ti-MIL-125 provides interesting avenues for chemical functionalization because the excitation from the ligand to metal cluster creates a transient Ti(III) center, which is stable for up to 900 ps^{14} . Moreover, unmodified Ti-MIL-125 features a band gap ($E_g \sim 3.8$ eV) that is tunable with simple organic functionalization²⁶.

The chemical modularity of these frameworks reflects their metastability and the chemical softness of the ligand-metal interface. The downside is that many frameworks decomposition in the presence of nucleophiles. Furthermore, weak electronic interaction at their interface results in the flat electronic bands (large effective masses for electrons and holes with $m^* \gg 1$ m_e) that are localized in real space.

Figure 2. Alignment of the electronic energy levels of metal-organic frameworks adapted from Ref. 27 . (a) Illustration of the internal vacuum level sampling procedure for Ti-MIL-125 along with 2D slices of the electron density (b) and electrostatic potential (c) inside the pore. (d) The calculated valence and conduction band positions for a range of frameworks that are discussed in the main text.

From electron energies to band diagrams

Knowing the electronic structure of a material can provide useful guidelines for design. The next step, in order to be able to consider a material in a device context, is to place the electron energies on an absolute scale, thus facilitating the construction of energy band diagrams, Figure 2. The construction of band diagrams is a cornerstone of semiconductor device design, as one of the pioneers of heterojunction design Herbert Kroemer famously stated that "If, in discussing a semiconductor problem, you cannot draw an Energy Band Diagram, this shows that you don't know what you are talking about"²⁸.

While ionization potential (electron removal energy) and electron affinity (electron addition energy) are properties that are easily defined for solids, they are

difficult to quantify. It was recognized as early as the 1930s that the photoelectric threshold of solids is influenced by two factors: the bulk band energies and the surface electrostatic double layer²⁹. Due to the sensitivity of the surface term to the history, environment and morphology of a particular sample, there is a large variation in reported values for a given compound.

Similarly for theorists, the use of periodic boundary conditions introduces problems in predicting reliable electron energies. Whilst periodic boundaries offer a convenient and elegant route to representing an infinite crystal from a finite repeating unit, the solution of the Coulomb interactions introduces an arbitrary reference point for the electrostatic potential, meaning that band energies cannot be compared directly between systems^{30,31,32}. In another approach the crystal can be modelled by considering a representative cluster of finite size, this provides an external vacuum level and facilitates comparison of band energies between systems^{33,34}; however, in practice the models are challenging to construct for multi-component systems.

In this context, we formulated a new procedure to calculate the bulk band energies of porous solids²⁷. The method is based on a periodic DFT approach, thereby avoiding truncation problems, and samples the electrostatic potential at the centre of a pore where the electron density has decayed to zero and the electrostatic potential has a plateau (see Figure 2). The procedure, which is applicable to porous solids, is analogous to the use of the external vacuum level above a surface slab as a reference in solid-state calculations. The center of large pores is also a vacuum that provides a reference level for alignment between different materials.

Practically, the alignment is achieved by performing a periodic electronic structure calculation at a given level of theory. The level of theory will affect the reliability of the results and a hybrid exchange-correlational functional (or beyond) is recommended for quantitative insights. The electrostatic (Hartree) density at the centre of the largest pore is then evaluated; this can be achieved using the MacroDensity (https://github.com/WMD-group/MacroDensity) package, which we have developed for this purpose. One must ensure that the density has reached a reliable convergence, which is achieved by sampling a large enough area (a sphere of radius > 2 Å) and ensuing that the variance within the sampling volume is small $(< 0.01$ V). The Hartree potential thus obtained is now referred to as the vacuum potential (V_{vac}) . The ionization potential is simply determined by subtracting the highest occupied band eigenvalue (E_{VB}): IP = V_{vac} . E_{VB} . Similarly, the electron affinity is determined using the lowest unoccupied band eigenvalue (E_{CB}) . The resultant ionization potentials and electron affinities can be compared across all materials.

Access to absolute band energies allows for the intelligent design of a range of physical applications for MOFs. Recently the group of Cheng found that using a newly designed ligand they could obtain a material, which electronic structure calculations predict to have band energies suitable for hydrolysis and degradation of organic pollutants³⁵. Moreover, they experimentally demonstrate the strong photocatalytic activity of the new material. The alignment of energy levels has also been applied to explain the differences in photocatalytic activity in d⁰ MOFs¹⁴. Grau-Crespo and co-workers demonstrated control of frontier orbital positions by metal substitution in porphyrin-based MOFs, allowing crystal engineering for solar fuel production³⁶. In the context of device design, access to absolute energy levels as well as lattice parameters from sources such as the CoRE-MOF database will allow for the application of design strategies such as the recently developed ELS (electronic lattice site) metric³⁷ allows us to fulfill the conditions of Kroemer's lemma and opens the field for the design of heterojunctions featuring MOFs.

Figure 3. Four avenues for engineering the electronic structure and optoelectronic activity of metal-organic frameworks by modifying the inorganic (metal substitution) and organic (ligand engineering) units, incorporating extrinsic chemical moieties (host-guest inclusion) and mechanical effects from interfaces or embedding (lattice strain). The schematic is inspired by Figure 48 of the seminal nanochemistry perspective by G. A. Ozin³⁸.

Principles for electroactivation

Following our brief outline of chemical bonding in MOFs as well as the factors influencing the band energies, we now collate design principles for tailoring the physical properties. The critical criterion for realizing semiconducting MOFs is that the frontier electronic bands are delocalized (low carrier effective masses), which requires effective communication between the organic and inorganic building blocks. However, long-range transport is not always essential, and short-range electron transfer (e.g. to an electrolyte at the surface or penetrating the pore) can be sufficient to support a wide range of redox processes.

A set of four chemical principles for electroactivation of MOFs are outlined in Figure 3 and below. These range from modifying the organic and inorganic building blocks that form the framework to post-synthetic modification and the application of lattice and chemical strain.

1. Metal substitution

Depending on the chemical identity and charge state, the orbitals of the metal may form the upper valence or lower conduction band of the MOF. Therefore, metal substitution has the potential to influence both oxidation (hole injection) and reduction (electron injection) processes. For example in HKUST-1 the Cu 3d orbitals are found at the valence band maximum³⁹, while in MIL-125 the Ti 3d orbitals are found at the conduction band minimum. Beyond complete metal substitution, the formation of mixed-metal systems offers a promising route to tune redox activity as demonstrated by Brozek and Dinca for MOF-5 with Ti, V, Cr, Mn and Fe incorporation²⁵.

2. Ligand engineering

The choice of ligand can be used to tune the electronic structure *directly* by modifying the orbital composition or *indirectly* by changing the framework topology. Many MOFs have at least one frontier band centered on an organic conjugated region. The electron energies of the organic regions are therefore tunable using conventional "push-pull" principles, where the band energies are influenced by the electron donating or withdrawing capability of additional substituents. For example, the introduction of an electron-donating primary amine to a benzene ring will result in an enhancement of electron density and a lowering of the ionization potential. Amination has been effective at introducing visible light photoactivity in a range of simple metal-organic frameworks by raising the valence band energy, e.g. modification of benzenedicarboxylayte in MIL-125²⁶ and $UiO-66¹⁴$. As an alternative to ligand modification, different types of ligands are being explored, e.g. linear conjugated carboxylates have been shown to exhibit extended helical orbitals 40 with the potential for long-range magnetic coupling⁴¹.

3. Host-guest inclusion

The introduction of redox active molecular guests into a framework is one way to change the electronic properties, either through spontaneous or lightactivated charge transfer. A related process is the ability to "rewire" an insulating MOF though an auxiliary electroactive linker. The champion system in this regard is HKUST-1 modified with the molecule TCNQ (7,7,8,8 tetracyanoquinododimethane) that provided a route to tunable conductivity⁴². In the recent work of Allen and Cohen, cross-linking in a series of isorecticular MOFs was demonstrated²¹, which could be extended to install conductive pathways for a wider range of MOF topologies.

4. Lattice strain

A feature of porous frameworks is that they are mechanically soft and flexible in comparison to close-packed materials, which provides another route to engineer their properties. Stress and strain can be introduced in several ways, e.g. epitaxial growth on a rigid substrate, solid-state embedding of MOF particles in a host, the application of pressure either mechanically or chemically (e.g. defects such as missing ligands are sensitive to the growth conditions and reagents)⁴³. The electronic structure response is described by the deformation potential: the change in ionization potential, electronic affinity or band gap with respect to a volume (pressure) change. It has been shown that the volume deformation potential are comparable to inorganic semiconductors⁴⁴; however, critically due to smaller bulk moduli the associate pressure coefficients are much larger. In the compound NCOF-1 the band gap pressure coefficient of -2.2 eV/GPa is 20 times larger than for bulk $Si⁴⁵$. For these materials, a small external stimulus can result in a large change in the electronic structure.

Figure 4. The increase in the number of publications on conductive MOFs (Web of Science, June 2016) with a timeline for future progress for MOF devices building on chemical and physical functionality.

Summary and Outlook

Our discussion has focused on the successes of modern simulation techniques in describing the chemical bonding and physical properties of metalorganic frameworks. While recent progress has been substantial, the *a priori* design of functional hybrid solids remains a daunting challenge. We have outlined fertile avenues for the pursuit of materials where the injection, extraction, excitation and transport of charge carriers is controllable, which build on the hybrid compounds reported in this special issue. Not only can the composition of the framework itself be engineered, but the porosity can be exploited for absorbate interactions that activate optical or electronic activity.

The approaches we have outlined for engineering physical properties can be exploited for the use of MOFs in a range of devices, where the hybrid advantage offers the prospect of disruptive technologies. With the number of reports on conductive MOFs rapidly increasing, a roadmap for future progress is illustrated in Figure 4, drawing from the earlier perspective of Allendorf et al^{46} . In Table I we list several possibilities that highlight the potential role for materials modelling. Referring back to the Nobel lecture of Herbert Kroemer: "The

principal applications of any sufficiently new and innovative technology have always been – and will continue to be – applications *created* by that technology".

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Tables

Table I: Potential applications for electroactive MOFs, with key physical properties and potential for input from materials modelling.

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