

Functional Stability of a Ferromagnetic Polycrystalline Ni₂MnGa High Temperature Shape Memory Alloy

M.A. Azeem^{a,b,c,*}, N.G. Jones^{a,d}, S.L. Raghunathan^{a,e}, V.A. Vorontsov^a, D. Dye^{a,*}

^aDept. Materials, Royal School of Mines, Imperial College, Prince Consort Road, South Kensington, London SW7 2BP, UK

^bSchool of Materials, The University of Manchester, Manchester M13 9PL, UK

^cResearch Complex at Harwell, Harwell Science & Innovation Campus, Didcot OX11 0FA, UK

^dDepartment of Materials Science and Metallurgy, University of Cambridge, 27 Charles Babbage Road, Cambridge, CB3 0FS, UK

^eCryogenic Ltd, Unit 6, Acton Park Estate, The Vale, Acton, London W3 7QE

Abstract

Electrocaloric Ni₂MnGa is of interest for solid state refrigeration applications, as well as a high temperature thermal shape memory alloy. Here, polycrystalline Ni₅₄Mn₂₅Ga₂₁ is examined using in situ synchrotron X-ray diffraction. The initial martensite (M_f) and austenite (A_f) finish temperatures were found to be 232 °C and 298 °C respectively. M_f was observed to decline by 8 °C/cycle and A_f increased by 1 °C/cycle. Both below and surprisingly, above the Curie temperature, the application of an e.m.f. was found to affect the lattice parameters measured. A change in the thermal expansion of the two phases was found around the Curie temperature.

Keywords: Ferromagnetic shape memory alloy; synchrotron radiation; thermal expansion; intermetallic compounds

1 Ferromagnetic shape memory alloys (fSMAs) based on the 29
2 Ni-Mn-Ga system [1] are of great interest both as actuators and 30
3 as magnetocaloric materials. Examples range from fast actua- 31
4 tion (using an applied magnetic field) in aerodynamic control 32
5 surfaces [2, 3] to magnetically operated solid-state heat pumps 33
6 offering coefficients of performance (COP) > 4 [4]. There 34
7 has also been interest in using the mechanically induced phase 35
8 transformation in conventional SMAs (e.g. NiTi) in a simi- 36
9 lar way [5, 6]. Large magnetic field induced strains (MFIS) 37
10 are widely reported in these alloys [7–11] and the underlying 38
11 mechanism is thought to be the magnetically induced mobility 39
12 of twin interfaces [8, 12]. 40

13 Shape memory functionality was first observed in the 41
14 Hoesler compound Ni₂MnGa [13] and the ferromagnetic phase 42
15 transformation was discovered almost a decade later in off- 43
16 stoichiometric Ni-Mn-Ga alloys [14]. The martensite transfor- 44
17 mation temperatures (T_m) are sensitive to deviations in stoi- 45
18 chiometry; Ni substitution for Ga result in changes of ~50 K 46
19 per at.% [1, 15, 16], but this results in only a modest change 47
20 in the Curie temperature (T_c) [1, 15]. It is observed that the T_m 48
21 increases linearly with the alloy electron concentration, i.e., in- 49
22 creases in the electron per atom ratio (e/a) [16–18]. Therefore, 50
23 in developing high temperature Ni-Mn-Ga alloys, replacement 51
24 of Ga by Ni are generally preferred as Ni has the greatest num- 52
25 ber of valence electrons and since the Ni fraction is greatest. For 53
26 example, the M_f in a Ni₅₄Mn_{24.7}Ga_{21.3} alloy ($e/a = 7.768$) [15] 54
27 is ~215 K higher than in near-stoichiometric Ni₅₀Mn₂₅Ga₂₅ 55
28 ($e/a = 7.5$) [7, 18–20]. 56

Irrespective of composition, the austenite in Ni-Mn-Ga al-
loys is reported to be Fm $\bar{3}$ m but the martensite crystal struc-
ture varies with small changes in composition [15, 21]. Both
modulated and non-modulated martensites have been reported;
5M and 7M martensites are observed in alloys with Curie tem-
peratures (T_c) higher than T_m whilst non-modulated marten-
sites are observed in alloys where $T_c < T_m$ [15]. In Mn-rich
alloys with ~20 at.% Ga, orthorhombic and tetragonal marten-
sites are observed, and for certain compositions both forms can
co-exist [21]. In the tetragonal martensite increases in e/a in-
creases the tetragonality, c/a_t [19].

Single crystals have been used in most investigations of Ni-
Mn-Ga alloys, but single crystal fabrication will limit the range
of product forms that can be economically produced, so poly-
crystalline Ni-Mn-Ga alloys are highly desirable. In this con-
text, purely thermal high temperature Ni₂MnGa actuators are
also of interest, as conventional NiTi-based SMAs are limited to
austenite finish temperatures of ~90°C. Normally, SMAs show
significant functional fatigue, with the transformation tempera-
tures dropping with each cycle, and so practical actuators with
service lives of thousands of cycles remain difficult to produce.

Here, we focus on the characterisation of the thermal trans-
formation in polycrystalline off-stoichiometry Ni₅₄Mn₂₅Ga₂₁
by synchrotron X-ray diffraction, with an emphasis on the evo-
lution of the unit cell parameters of the phases. This particu-
lar composition was chosen to satisfy the requirement that the
transformation be single step (cubic↔tetragonal) [22, 23] and
to obtain a high temperature shape memory candidate with T_m
in excess of 200°C [22].

The alloy was prepared by vacuum arc melting pure elemen-
tal Ni, Mn and Ga in pure argon. The ingot was remelted four

*Corresponding author

Email address: mohammed.azeem@manchester.ac.uk,
david.dye@imperial.ac.uk (D. Dye)

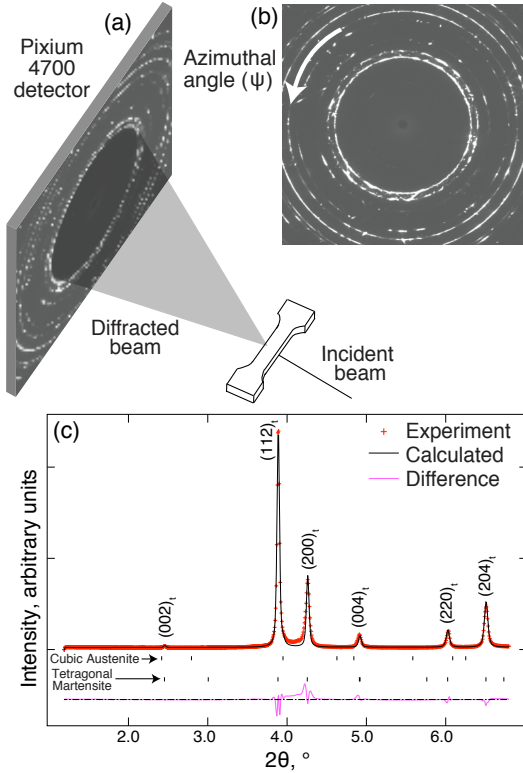


Figure 1: (a) Schematic of the synchrotron X-ray diffraction setup, (b) Complete Debye-Scherrer rings from tetragonal martensite and (c) Diffraction spectra after 360° intensity integration across diffraction rings in (b)

times and inter-melt flipped to minimise elemental segregation. The ≈ 60 g ingot was encapsulated in a quartz tube and then homogenised at 900°C for 24 h. The ingot was then packed in Ti powder in a mild steel can and 80% hot rolled at 900°C . $0.5 \times 0.5 \times 30$ mm gauge length specimens were then fabricated by grinding and electro-discharge machining. Differential scanning calorimetry using a Mettler Toledo DSC 822^e instrument was used to obtain initial estimates of the transformation temperatures. The composition was also checked at this stage by SEM-EDX and found to be within the measurement error (~ 0.5 at.%) of the target composition. Light element pickup such as O was therefore not determined. The specimen was then thermally cycled at a rate of 2°C s^{-1} at zero load in an Instron5 kN resistance heating electro-thermal mechanical testing machine (ETMT) at the ID15B powder diffraction beamline at the ESRF, Grenoble France, Figure 1. The sample temperature was measured by using R-type thermocouples mounted on the sample surface, within 0.1 mm of the scanned volume. A wavelength of 0.1427 \AA was used and the incident beam was approx. 0.5×0.5 mm in size, giving a sampled volume of $0.5 \times 0.5 \times 0.5 = 0.125 \text{ mm}^3$. A Thales Pixium 4700 detector [24] was used to collect complete Debye-Scherrer diffraction rings (Figure 1b) at a sample-detector distance of 1.228 m. A LaB_6 calibration powder was used to determine the instrument parameters. Phase fractions and lattice parameters were evaluated using the generalized structural analysis package (GSAS) [26, 27] and the area detector analysis pack-

age Fit2D [28]. Whole-ring integrations were used; the analysis was not found to be sensitive to orientation around the ring; *i.e.* there were no effects induced by, *e.g.* the rolling texture. This issue has been examined in our previous work [29, 30].

The high temperature (austenite) phase in $\text{Ni}_{54}\text{Mn}_{25}\text{Ga}_{21}$ has an $Fm\bar{3}m$ crystal structure with Ni at $(1/4, 1/4, 1/4)$, Mn at $(1/2, 1/2, 1/2)$ and Ga at $(0, 0, 0)$ and $a_c = 5.8468 \pm 0.0003$ at 320°C . For thermal cycling in the absence of stress and magnetic field, the martensite is tetragonal. The room temperature lattice parameters of the martensite were found to be $a_t = 3.8391 \pm 0.0003$, $c_t = 6.650 \pm 0.001$ and $c_t/a_t = 1.732$. Comparing these with the lattice parameters for $\text{Ni}_{53}\text{Mn}_{25}\text{Ga}_{22}$ reported by Cong *et al.* [23] of $a_t = 3.865$, $c_t = 6.596$ and $c_t/a_t = 1.707$, the increase in e/a achieved by increasing Ni (at the expense of Ga) in the current alloy appears to change the c/a ratio considerably. The main orientation relationship in $\text{Ni}_{53}\text{Mn}_{25}\text{Ga}_{22}$ is reported to be $(111)_c \parallel (101)_t$, $[110]_c \parallel [111]_t$ [1, 31].

The transformation sequence in the first thermal cycle is shown in Figure 2. A large dilation in the $d_{(200)}$ interplanar spacing in the tetragonal martensite was clearly visible on heating and cooling. A new peak emerged in vicinity of the tetragonal (200) peak during the initial stages of cooling, but apart from this no additional peaks were observed elsewhere in the spectra. Therefore it is difficult to confirm the existence of an intermediate martensite phase. A subtle yet visible shift in the (200) tetragonal peak is observed (arrowed) on final cooling and before reheating began.

The Rietveld-refined lattice parameters and phase fractions are shown in Figure 3. The transformation from tetragonal to cubic begins on heating at around 213°C with 80% of the cubic phase appearing over a temperature range of $227 - 258^\circ\text{C}$, after which the transformation rate decreases with temperature before final completion at 286°C . The measurement points occur approximately every 5°C , allowing the austenite and martensite start and finish temperatures (at 90% / 10% phase fraction) to be interpolated with an accuracy of around 3°C , Table 1. The fitted lattice parameters, Figure 3, show almost no shift between successive thermal cycles, which is a good indicator of thermal stability. For example, the change in a_t at 125°C after four thermal cycles was only 0.04% (Table 1). On initial heating, a_t increased at a rate of $\approx 77 \times 10^{-6}^\circ\text{C}^{-1}$ and c_t decreases with a thermal expansion coefficient of $-69 \times 10^{-6}^\circ\text{C}^{-1}$. At $\approx 160^\circ\text{C}$ (near the Curie temperature estimated elsewhere [15, 23]), a visible change in linearity is observed and the thermal expansion / contraction appears to accelerate; *i.e.*, a_t then dilates and c_t constricts at a more rapidly than before. These events are schematically highlighted in the final subfigure in Figure 3. The reverse behaviour is observed on cooling and again the inflection point is $\approx 160^\circ\text{C}$. One can speculate that this change in thermal expansion is associated with a change in the magnetic ordering in the material in the vicinity of the reported Curie temperature; rationalisation of this result would benefit from *in situ* magnetic neutron diffraction measurements.

In the cubic high temperature phase, the corresponding thermal expansion coefficient is measured to be $13 \times 10^{-6}^\circ\text{C}^{-1}$. During transformation, contraction of c_t initially ceases (at

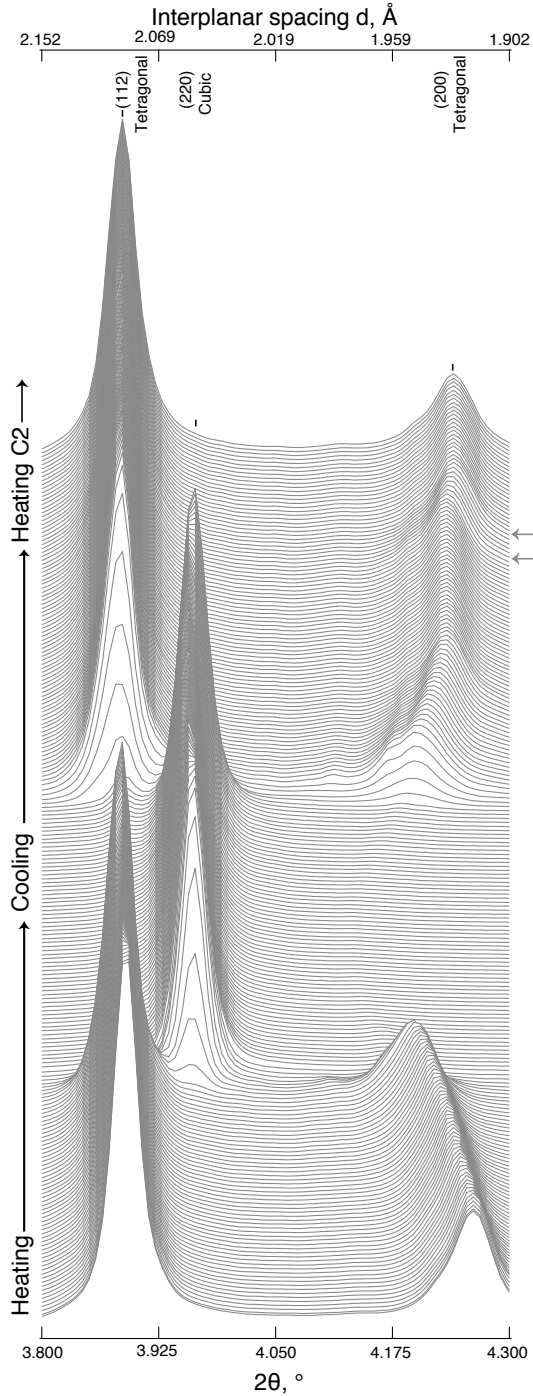


Figure 2: Evolution of (sections of the) diffraction pattern during the first thermal cycle in $\text{Ni}_{54}\text{Mn}_{25}\text{Ga}_{21}$. The e.m.f reversal regime during which a shift in martensite lattice parameters was observed is indicated by arrows on the right.

144 $\approx 235^\circ\text{C}$), but during the final stages of transformation it re-
 145 sumes; a similar behaviour in reverse is observed for a_t . During
 146 cooling, at approximately 256°C , indicated by arrow in Fig-
 147 ure 3, there appears to be a rather large instantaneous jump in
 148 both c_t and the austenite lattice parameter. Remarkably, this
 149 phenomenon is only observed during cooling and in case of
 150 tetragonal martensite, it is only observed in the c_t lattice pa-
 151 rameter.

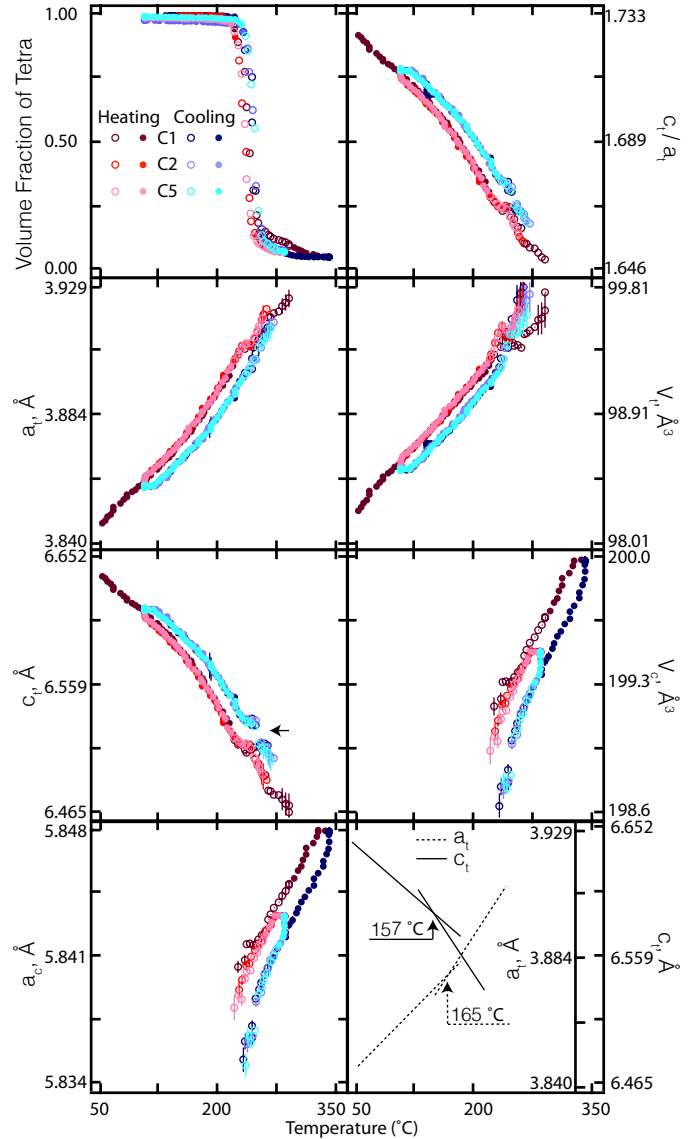


Figure 3: Tetragonal martensite and cubic austenite unit cell evolution evaluated as a function of sample temperature. Solid symbols are in the single phase region and open symbols during transformation. 'V' is the unit cell volume and subscripts t and c signify the tetragonal and cubic phases, respectively. Cycles 1 (C1), 2 (C2) and 5 (C5) are shown.

152 Interestingly, a hysteresis is observed in the lattice param-
 153 eters in both the cubic and tetragonal phases, not associated with
 154 the transformation itself. In the first thermal cycle (brown / dark blue)
 155 which reversed at 345°C the lattice parameters converge
 156 at that point, whereas in subsequent cycles they converge at the
 157 end point of 300°C . This is a consequence of the resistance
 158 heating method used, where the direct current applied produces
 159 an e.m.f. which, in this ferromagnetic material, results in a
 160 change in the lattice parameters. However, the Curie tempera-
 161 ture in these alloys is widely reported to be $\approx 110^\circ\text{C}$ [15]. This
 162 observation therefore poses a new question such as to whether
 163 these alloys shed their ferromagnetic properties completely, or
 164 not.

The lattice parameters are shown for the first, second and

Table 1: Lattice parameters (\AA) of tetragonal martensite on cooling and heating (at $125\text{ }^\circ\text{C}$ for martensite and at $270\text{ }^\circ\text{C}$ for austenite), where, c and h superscripts are for cooling and heating, A_s , A_f , M_s and M_f are austenite start, finish, martensite start and finish respectively.

	Cycle 1	Cycle 5
a_t^h	3.8653 ± 0.0003	3.8667 ± 0.0004
a_t^c	3.8606 ± 0.0004	3.8604 ± 0.0003
c_t^h	6.603 ± 0.001	6.600 ± 0.001
c_t^c	6.611 ± 0.001	6.612 ± 0.001
a_c^h	5.8433 ± 0.0003	5.8425 ± 0.0002
a_c^c	5.8403 ± 0.0003	5.8403 ± 0.0002
Austenite start, A_s	226	227
Austenite finish, A_f	298	258
Martensite start, M_s	267	268
Martensite finish, M_f	232	236

fifth thermal cycle applied. It is notable that, whilst the lattice parameters at a given temperature are invariant with cycling, the transformation temperatures do evolve, by $\sim 8\text{ K/cycle}$ for the austenite finish temperature A_f .

The evolution of the ETMT actuator position, along with the ex-situ DSC data, are shown in Figure 4. For comparison, DSC

data for binary NiTi and two $\text{Ni}_{50}\text{Ti}_{35}\text{X}_{15}$ ($\text{X}=\text{Hf}$ or Zr) high temperature shape memory alloys are also shown. The DSC results reproduce the observed transformation temperatures from the diffraction measurements and show a hysteresis (A_f-M_f) of around of 22 K in cycle 5. The austenite finish temperature gradually declines on cycling, but the martensite transformation temperatures remain constant, which is unusual compared to NiTi SMAs. In addition, the ETMT actuator position, which is in effect a dilatometry measurement over the entire sample length (some of which remained cold and did not transform), is shown. The transformation on heating results in a contraction of the sample which is reversed on cooling, but there is an evolution in the end-point actuator positions with cycling that indicates expansion of the sample from cycle 2 through 5. Since this occurs for the *fcc* austenite as well as the martensite, this cannot be due to the selection of preferred crystallographic variants and therefore must be due to the accumulation of plasticity, *i.e.* interface dislocations. By reference to conventional NiTi SMAs [32, 33], this would also be the reason for the gradual decline in the austenite transformation temperatures with cycling.

In summary, polycrystalline $\text{Ni}_{54}\text{Mn}_{25}\text{Ga}_{21}$ has been produced by a conventional ingot metallurgy and hot working approach. It is found to be a viable high temperature shape memory alloy material with $M_f = 236\text{ }^\circ\text{C}$. The lattice parameters measured in both phases were found to be affected by the magnetic fields induced in the specimen by the applied ohmic heating method used, even above the expected Curie temperature of the alloy. Some functional fatigue was observed in the austenite transformation temperatures on heating. The thermal expansion / contraction of the tetragonal martensite a_t and c_t cell parameters were observed to accelerate above the expected Curie temperature.

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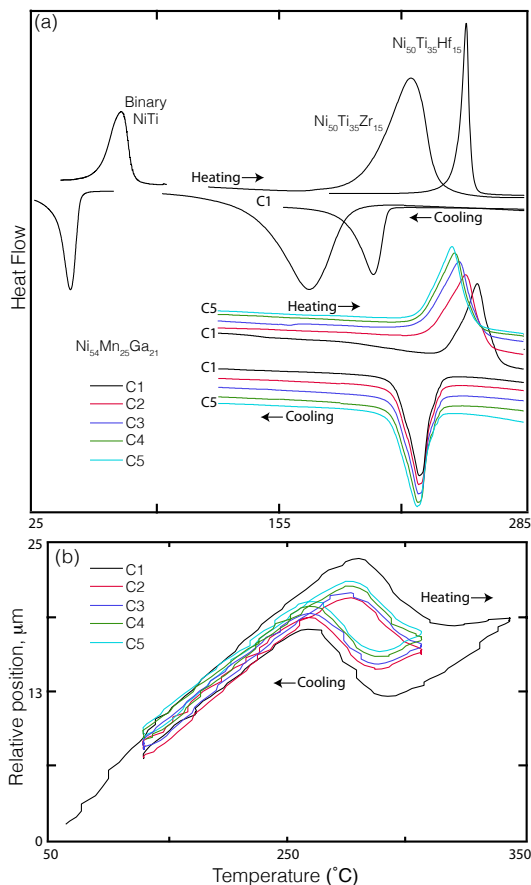


Figure 4: (a) Differential Scanning calorimeter (DSC) results obtained for five heating and five cooling cycles of as-rolled $\text{Ni}_{54}\text{Mn}_{25}\text{Ga}_{21}$ are shown, in comparison with the first cycle results for binary NiTi and as-rolled high temperature SMAs $\text{Ni}_{50}\text{Ti}_{35}\text{Hf}_{15}$ and $\text{Ni}_{50}\text{Ti}_{35}\text{Zr}_{15}$ alloys and (b) Dilation in macroscopic $\text{Ni}_{54}\text{Mn}_{25}\text{Ga}_{21}$ along the rolling direction during thermal cycling.

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