



**Effect of extreme ocean precipitation on sea surface elevation and storm surges**

Journal:	<i>QJRMS</i>
Manuscript ID	QJ-15-0243.R1
Wiley - Manuscript type:	Research Article
Date Submitted by the Author:	09-Feb-2016
Complete List of Authors:	Wong, Benjamin; Imperial College London, Physics Toumi, Ralf; Imperial College London, Physics
Keywords:	ocean modelling, Cyclone Monica, rain mass, rain stress

SCHOLARONE™  
Manuscripts

Review

1  
2  
3 **Effect of extreme ocean precipitation on sea surface elevation and storm surges**  
4

5  
6  
7 Running head: Effect of extreme rainfall on sea surface elevation and storm surges  
8  
9

10  
11  
12  
13 Benjamin Wong (corresponding author): benjamin.wong11@imperial.ac.uk  
14  
15

16  
17 Ralf Toumi: r.toumi@imperial.ac.uk  
18  
19

20  
21  
22  
23 **Space and Atmospheric Physics Group,**  
24

25  
26 **Blackett Laboratory,**  
27

28  
29 **Imperial College London,**  
30

31  
32 **London SW7 2AZ,**  
33

34  
35 **UK**  
36  
37  
38  
39  
40  
41

42 **Keywords:** ocean modelling, Cyclone Monica, rain mass, rain stress  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

## Abstract

Ocean models that neglect mass and momentum contributions from precipitation can have a systematic bias in sea surface height (SSH). Here, a new rainfall scheme is introduced into the Regional Ocean Modelling System (ROMS) to incorporate the effects of precipitation mass. When precipitation is added to the sea surface, it spreads out via surface gravity waves that increase in propagation speed with increasing water depth. Over several days, the SSH increase due to the precipitation mass added created a geostrophic adjustment, generating anti-cyclonic geostrophic currents around the SSH increase. The transfer of momentum from precipitation to the sea surface, or rain stress, can also be important. In the case study of a real tropical cyclone, Monica passing North Australia, the effect of incorporating precipitation mass is compared with other processes affecting the storm surge: surface wind, inverse barometer effect and rain stress. The maximum SSH response is 170.6 cm for the wind effect, 61.5 cm for the inverse barometer effect, 7.5 cm for the effect of rain stress and 6.4 cm for the effect of rain mass. Each process has been shown to have different spatial influences. The effect of rain mass has a strong remote influence compared to the inverse barometer effect and the effect of rain stress. This is particularly seen in semi-enclosed bays.

## 1 Introduction

The contribution of oceanic precipitation to storm surges has so far been neglected in storm surge models, typically only incorporating the contributions of the inverse barometer effect (Roden et al., 1999) and wind driven effect towards the overall surge height. There is very little work addressing the contribution of oceanic precipitation to surge levels. This is expected, considering that many ocean models do not incorporate the mass and momentum effect of precipitation for it to be studied in detail. Sheng et al. (2012) described a test

1  
2  
3 platform used to review several established storm surge models such as CH3D-SSMS (Sheng  
4 et al., 2010), ADCIRC (Westerink et al., 1992), FVCOM (Chen et al., 2008), CMEPS (Peng  
5 et al., 2004) and SLOSH (Jarvinen et al., 1985). Thirty scenarios were used to test their  
6 sensitivity towards parameters such as bathymetry, storm forcing, wind drag coefficient,  
7 bottom friction, Coriolis and 2D-3D formulation, but precipitation was not investigated.  
8 Ocean models such as ROMS (Shchepetkin et al., 2005) and FVCOM only incorporate the  
9 effect of temperature, salinity and momentum from rain, but not the effect of rain mass.  
10 Storm surge models such as ADCIRC and SLOSH ignore the effects of ocean rain entirely.  
11 Ponte (2006) incorporated realistic freshwater fluxes in a global barotropic model and found  
12 that they can cause annual standard deviations in sea surface level as large as 1 cm, especially  
13 in shallow and semi-enclosed regions. In highly precipitating storm events with extreme rain  
14 rates, the sea surface height increase may be more substantial. For instance, the highest  
15 rainfall recorded by Typhoon Morakot in 2009 was up to  $74 \text{ cm day}^{-1}$  (Wu, 2013). How this  
16 contributes to storm surges and to what extent rain mass can be neglected in this case is not  
17 known.

18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37 Here, a new rainfall scheme has been introduced into ROMS to account for the effect of rain  
38 mass. This scheme complements the existing rainfall routine in ROMS where only  
39 temperature and salinity changes due to rainfall are accounted for. An idealized set-up of  
40 ROMS will be used to isolate the effects of adding rain mass to the sea surface height. The  
41 Weather Research and Forecasting model (WRF, Skamarock et al., 2008) is used to generate  
42 the rain forcing to evaluate the sea surface height response in ROMS.

43  
44  
45  
46  
47  
48  
49  
50  
51  
52 The transfer of momentum from raindrops to the water surface, or rain stress, was first  
53 documented by Van Dorn (1953), who concluded that the stress contributions from rainfall  
54 can considerably intensify surface stresses. Caldwell et al. (1970) have since built on his  
55  
56  
57  
58  
59  
60

1  
2  
3 work, parameterizing rain stress as a function of rain rate and wind speed. This  
4  
5 parameterization has been widely incorporated in many model studies such as the one-  
6  
7 dimensional mixed layer model developed by Clayson et al. (1999) and the bulk  
8  
9 parameterization outlined by Fairall et al. (1996) relating the near-surface atmospheric and  
10  
11 the oceanographic bulk variables. This air-sea bulk parameterization has also been adopted in  
12  
13 ROMS. While the parameterization of rain stress has been widely incorporated, there has  
14  
15 been no study so far that shows its isolated effect in the models it has been implemented in.  
16  
17 For the case of tropical cyclones, where extreme wind speeds and precipitation rates are  
18  
19 found, the isolated effect of rain stress has not been shown and its contribution to storm  
20  
21 surges is not known. This section shall employ the use of an idealised set-up in ROMS to  
22  
23 show the isolated effects of rain stress in the context of storm surges.  
24  
25  
26  
27

28  
29 Finally, a real case scenario is studied to compare the various processes affecting storm  
30  
31 surges. Tropical Cyclone Monica (Durden, 2010) has been chosen for her high wind speed,  
32  
33 low system pressure and high rainfall, which are the key storm surge effects investigated  
34  
35 here. A validated atmospheric state is generated using WRF and is prescribed onto ROMS to  
36  
37 investigate how various processes affect the storm surge. The significance of rain mass and  
38  
39 rain stress is compared with two other established storm surge processes: the inverse  
40  
41 barometer effect and the effect of surface wind. Each of these physical effects is investigated  
42  
43 based on their isolated contribution to the total SSH and storm surge patterns.  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

## 2 Model set-up and validation

### 2.1 Set-up of idealized case

Rain mass is introduced to ROMS through a new rainfall scheme that incorporates rain rate as point sources. This approach is similar to how river run-off is introduced to the model, but in this case, vertical fluxes are added instead of horizontal fluxes. The rainfall rate from the atmospheric rain forcing is converted into point sources without further prescribing its temperature and salinity. The temperature and salinity of rain is accounted for within the existing bulk flux routine where air-sea momentum and energy exchanges are calculated (Fairall et al., 1996). An idealized set-up of ROMS is used to create a controlled environment, isolating the specific effects investigated here. The ROMS grid is set up with closed boundaries, a horizontal resolution of 15 km, 21 vertical levels and a spatial extent 5°N to 35°N and 120°E to 180°E. The initial ocean state has a uniform temperature of 26°C and salinity of 35 psu. The land mask is specified at the western boundary up to 140°E to simulate the coastline. A third-order upwind scheme is used for horizontal momentum advection, and a Smagorinsky-like viscosity is applied (Griffies et al., 2000). The turbulence closure scheme used to calculate vertical mixing is based on the Generic Length Scale (GLS) parameterization as described by Umlauf et al. (2003).

An idealized tropical cyclone is simulated using WRF and prescribed onto ROMS. The set-up of the initial cyclone profile and atmospheric state follows Wang et al. (2015). For the WRF simulation, one single nest is set up with 31 vertical levels and a horizontal resolution of 15 km in the domain spanning 5°N to 35°N and 120°E to 180°E. The microphysical parameterization used is based on Lin et al. (1983). Cumulus parameterization is not included, which reduces the computational time required for each simulation. Without

1  
2  
3 cumulus parameterization, the amount of precipitation may be underestimated at a 15 km  
4  
5 horizontal resolution. However, the aim of this part of the study is to understand the specific  
6  
7 effect of adding precipitation mass to the ocean model in a simple, idealized set-up, and not  
8  
9 to reproduce the atmospheric domain with high accuracy. The surface layer scheme is based  
10  
11 on Monin-Obukhov, with Zilitinkevich thermal roughness length and standard similarity  
12  
13 functions (Paulson, 1970). Planetary boundary layer processes are parameterized using the  
14  
15 Mellor-Yamada-Janjic scheme (Janjic, 1994). The longwave and shortwave radiation  
16  
17 schemes are implemented from Mlawer et al. (1997) and J. Dudhia (1989) respectively.  
18  
19  
20  
21  
22

## 23 2.2 Rain stress parameterization

24  
25  
26 In ROMS, rain stress is implemented in the bulk fluxes routine from Fairall et al. (1996) and  
27  
28 based on the parameterization by Caldwell and Elliott (1970) as follows:  
29  
30  
31

$$32 \tau_{rx} = 0.85R\rho_r |\vec{v}_{10}| * SIGN(U_{10}) \quad (1)$$

$$33 \tau_{ry} = 0.85R\rho_r |\vec{v}_{10}| * SIGN(V_{10}) \quad (2)$$

34  
35  
36 where R is the rain rate,  $\rho_r$  is the density of rainwater,  $\tau_{rx}$  is the rain stress in the x direction,  
37  
38  $SIGN(U_{10})$  denotes the sign of  $U_{10}$  (U-component wind speed at 10 m),  $\tau_{ry}$  is the rain stress  
39  
40 in the y direction and  $SIGN(V_{10})$  denotes the sign of  $V_{10}$  (V-component wind speed at 10 m).  
41  
42  
43  
44  
45  
46  
47  
48

49 The rain stress parameterization of Caldwell and Elliott (1970) is derived from a one-  
50  
51 dimensional model where  $\vec{v}_{10}$  refers to the wind velocity in the direction of reference. When  
52  
53 considering wind velocity resolved into its orthogonal components, it would be incorrect to  
54  
55 utilize the total velocity in calculating the rain stress. Here, the magnitude of the u-component  
56  
57  
58  
59  
60

winds ( $U$ ) and the v-component winds ( $V$ ) used in the calculation of rain stress will always be  $\sqrt{U^2 + V^2}$ . This results in a systematic overestimation in the magnitude of rain stress by a factor of  $\sqrt{2}$ . The error in the direction of rain stress is more inconsistent. Since  $U$  and  $V$  are always set to be equal in magnitude (equal to the net wind speed), the directions of the rain stress can only be in any of the four directions:  $45^\circ$ ,  $135^\circ$ ,  $225^\circ$  and  $315^\circ$ , unless there is zero rainfall or zero wind speed. Furthermore, the existing ROMS rain stress code always assigns a positive sign to zero values. For example, if  $U$  is zero, it will still be assigned the magnitude of  $V$  (now equal to the magnitude of the net wind speed), with the direction of positive  $U$ . The error in the direction of the rain stress term can range from  $0^\circ$  (when  $U=V$ ) to  $45^\circ$  (when either  $U=0$  or  $V=0$ ). Overall, this results in a systematic overestimation of rain stress and an incorrect direction calculated for the overall wind stress.

The corrected implementation of the rain stress parameterization is as follows:

$$\tau_{rx} = 0.85R\rho_r U_{10} \quad (3)$$

$$\tau_{ry} = 0.85R\rho_r V_{10} \quad (4)$$

An idealized set-up of ROMS is used to verify the corrected implementation of rain stress. Here, the u-component wind is set to a constant  $-1 \text{ m s}^{-1}$  and the v-component wind is set to zero. One centimetre of rain is introduced to a single cell (15 km by 15 km) in the centre of the domain ( $20^\circ\text{N } 150^\circ\text{E}$ ) at the first hour. Two simulation runs are conducted to compare the existing implementation of rain stress (using equations (1) and (2)) with the corrected implementation of rain stress (using equations (3) and (4)). Both simulation runs are subtracted from a control run without any rain stress to eliminate any noise such as the SSH response to the background u-component wind.



### 2.3 Set-up of real case

The ROMS set-up uses a single grid at a horizontal resolution of 6 km, with the open boundaries using the Flather (1976) conditions for the barotropic velocities and the Chapman (1985) conditions for sea surface elevation. A minimum depth of 10 m and maximum depth of 5500 m has been specified. Twenty layers are chosen in the vertically stretched terrain-following sigma-coordinate, with higher resolution placed at the surface and lower resolution placed at the bottom. Geometry of the bathymetry is obtained from Global Topography v14.1 (Smith, 1997). The tidal forcing is applied at the open boundaries and imposed on the elevation and the barotropic velocities. It is derived from 11 tidal harmonics that are extracted from the 1/12° resolution Pacific Ocean Atlas solution provided by the Oregon State University (OSU) Tidal Data Inversion (Egbert et al., 1994). A nudging relaxation zone with eight grids from the boundary is set up to relax the baroclinic structure to the forcing fields at the boundary (Marchesiello et al., 2001). The 1/12° resolution HYbrid Coordinate Ocean Model (HYCOM) (Bleck et al., 1981) global daily analysis data is used as the initial, boundary and nudging conditions for the model.

The pressure-driven effect in ROMS is obtained by applying the atmospheric pressure field from WRF and the sea level is adjusted based on the inverse barometer approximation (Gill and Niiler, 1973). The inverse barometer effect is generally a good approximation for sea level responses to local atmospheric pressure. For non-isostatic pressure effects, Hirose et al. (2001) found that deviations from the inverse barometer effect contributed only up to 2 cm in large-scale basins such as the Pacific and Atlantic Ocean. For smaller-scale basins such as Cape of Carpentaria, the non-isostatic deviation was under 0.5 cm and is much smaller compared to the precipitation effects in this case (Figure 10b).

1  
2  
3 WRF is used to simulate a high resolution atmospheric state of tropical cyclone Monica to be  
4 prescribed onto the ocean model for the case study. Two nests are employed with 28 vertical  
5 levels. The final grid has a horizontal resolution of 6 km. NCEP Climate Forecast System  
6 Reanalysis (CFSR) (Saha et al., 2010) has been used for initial and lateral boundary  
7 conditions. The boundary forcing is prescribed every 6 hours and the model interior is  
8 allowed to evolve freely. The ETA scheme (Rogers et al., 2001) is chosen for microphysical  
9 parameterization. The Kain-Fritsch cumulus parameterization scheme (Kain, 2004) has been  
10 used for both nests. The surface layer scheme is based on Monin-Obukhov with Carls-  
11 Boland viscous sub-layer and standard similarity functions. The selected planetary boundary  
12 layer is from Yonsei University (Hong et al., 2006). The Mesoscale Model 5 (MM5) 5-layer  
13 thermal soil temperature model (Dudhia, 1996) is used to calculate the heat and moisture  
14 fluxes over the land. The longwave and shortwave radiation schemes are implemented from  
15 Mlawer et al. (1997) and Chou et al. (1994) respectively.

16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33 The isolated effects of including precipitation mass source, rain stress, inverse barometer  
34 effect and surface wind forcing are individually investigated by comparing with four separate  
35 configurations that are simulated with precipitation mass source, without rain stress, without  
36 the inverse barometer correction and without surface wind forcing respectively. The five  
37 experiments conducted are tabulated in Table 2.1. Spatial plots of SSH are differenced from  
38 the control simulation (Case A) to illustrate the isolated effects of each process.  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

<i>Cases</i>	<i>WRF wind conditions applied to ROMS</i>	<i>Inverse barometer effect in ROMS</i>	<i>Rain stress effect in ROMS</i>	<i>New rain mass scheme</i>
<i>Case A</i>	Wind forcing applied	Turned ON	Turned ON	NOT used
<i>Case B</i>	NO wind forcing applied	Turned ON	Turned ON	NOT used
<i>Case C</i>	Wind forcing applied	Turned OFF	Turned ON	NOT used
<i>Case D</i>	Wind forcing applied	Turned ON	Turned OFF	NOT used
<i>Case E</i>	Wind forcing applied	Turned ON	Turned ON	In use

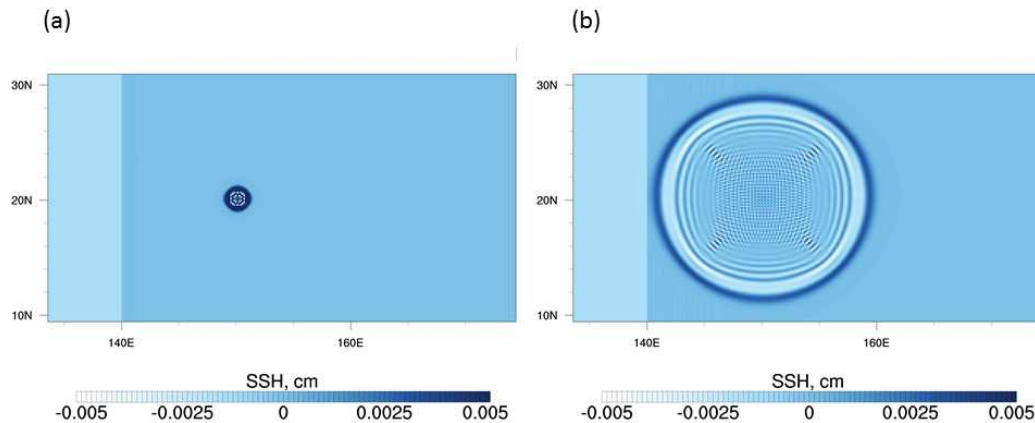
Table 2.1: Cases set up with different ROMS model options and wind conditions.

### 3 Results

#### 3.1 Idealized case: response to rain mass

An idealized set-up of ROMS is used to study the effect of rainfall mass on sea surface elevation. The effect of rainfall is isolated by setting the wind field to zero and all other atmospheric forcings are set to a constant nominal value. Rain is introduced to a single grid cell (15 km by 15 km) at 20°N 150°E for 1 hour at the rate of 1 cm/h, to understand the effect of a single-point perturbation on the free surface. Figure 1 shows the ocean surface response after 2 hours and 9 hours. The illustrations are zoomed into the area 9°N to 31°N and 134°E to 176°E. A surface gravity wave is generated, propagating radially outwards from the point of perturbation. At time = 1 h, the maximum SSH is 0.0181 cm and the minimum is -0.0076 cm. At time = 9 h, the maximum SSH is 0.0043 cm, the minimum is -0.0045 cm and the

1  
2  
3 wave propagates to about  $10^\circ$  from the initial point of perturbation. Over 9 hours, the  
4  
5 amplitude of the gravity wave drops by about 76%.  
6  
7  
8  
9



26 Figure 1: SSH response to rain mass addition to a single surface grid cell after (a) 1 hour and  
27 (b) 9 hours. Land is west of  $140^\circ\text{E}$ , ocean is east of  $140^\circ\text{E}$ . Wind forcing is set to zero.  
28  
29  
30

31  
32 Next, the SSH response to rainfall at a longer time scale is studied by forcing ROMS with  
33 rainfall from an idealized tropical cyclone simulated using WRF over 8 days. The maximum  
34 rain rate is not at the centre of the cyclone, but along the rain bands. As the cyclone treks  
35 from the east to the west of the domain, it gradually grows in size and intensity. The rain  
36 bands of the tropical cyclone also increase in size with time. Figure 2 shows the cumulative  
37 rainfall throughout the event, with a maximum of 319 mm at any one grid. To avoid any  
38 interference from the winds, the wind field is set to zero so the sea surface is unaffected by  
39 wind stress and responds only to the cyclone rainfall.  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

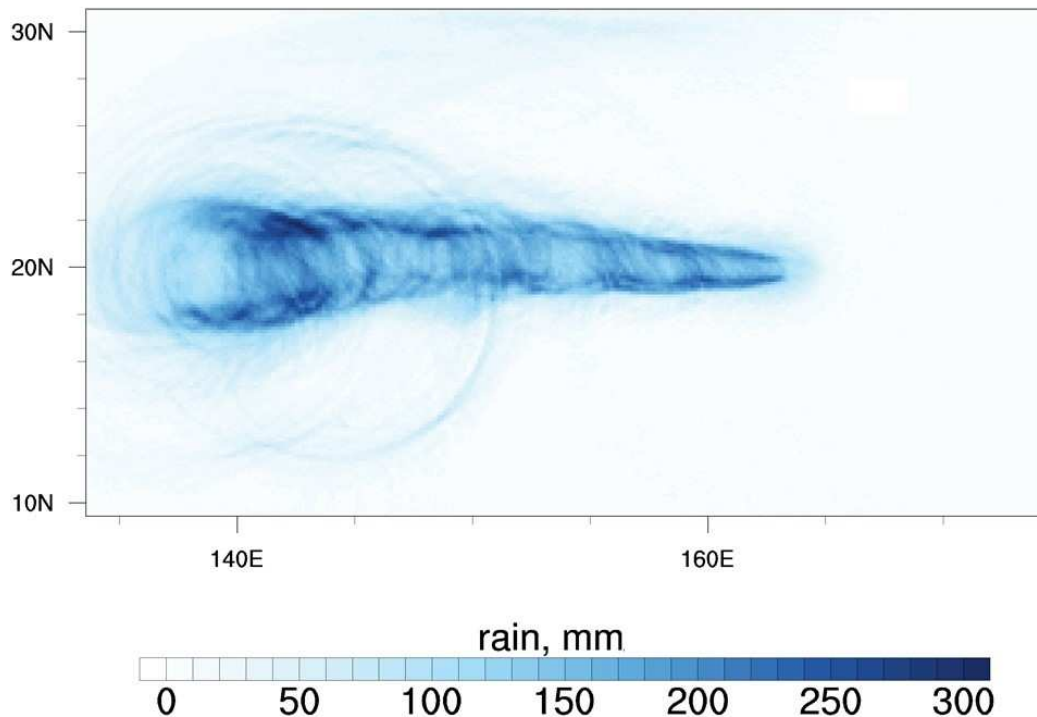


Figure 2: Cumulative rainfall of the idealized tropical cyclone over 8 days.

The SSH response to rainfall from the idealized tropical cyclone is simulated for two ocean depths at 100 m and 1000 m deep. Figure 3 shows the SSH at the 170<sup>th</sup> hour of the simulation when the tropical cyclone makes landfall at the coastline along 140°E. For both depths, the SSH increases along the rainfall path and gradually decreases towards the domain boundaries. For the case with a depth of 100 m, the maximum SSH is 5.8 cm and the minimum SSH is 1.3 cm at the 170<sup>th</sup> hour. A clockwise gyre is formed, with the sea surface currents at a magnitude of  $10^{-2} \text{ m s}^{-1}$ . For the case with a depth of 1000 m, the maximum SSH is 3.2 cm while the minimum SSH is 2.3 cm. A weaker clockwise gyre is formed with the sea surface currents at a magnitude of  $10^{-3} \text{ m s}^{-1}$ .

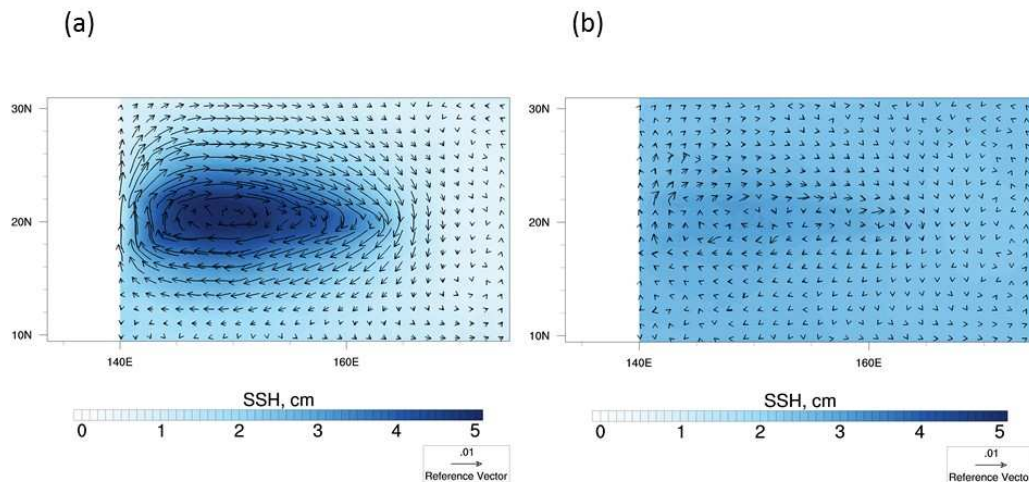


Figure 3: SSH and sea surface current response to the rain mass added to the ocean with (a) 100 m depth and (b) 1000 m depth at the 170<sup>th</sup> hour.

### 3.2 Idealized case: response to rain stress

Figure 4 shows the rain stress perturbation setting off a surface gravity wave that propagates outwards from the point of perturbation towards the domain boundaries. Two implementations are shown for the existing and the corrected model. In the existing implementation of rain stress (Figure 4a), the resulting surface gravity wave is at an angle of 45° to the wind direction. The front of the wave is led by a crest in the north-westward direction with an amplitude of up to 0.01 cm. In the south-eastward direction, the wave is led by a trough of up to 0.01 cm in amplitude. Figure 4b shows the SSH response to the corrected implementation of rain stress in ROMS. The surface gravity wave generated from the rain stress perturbation is now symmetrical to the wind direction. In the westward direction, the front of the wave is led by a crest of up to 0.005 cm in amplitude. In the eastward direction, the wave is led by a trough of up to 0.005 cm in amplitude.

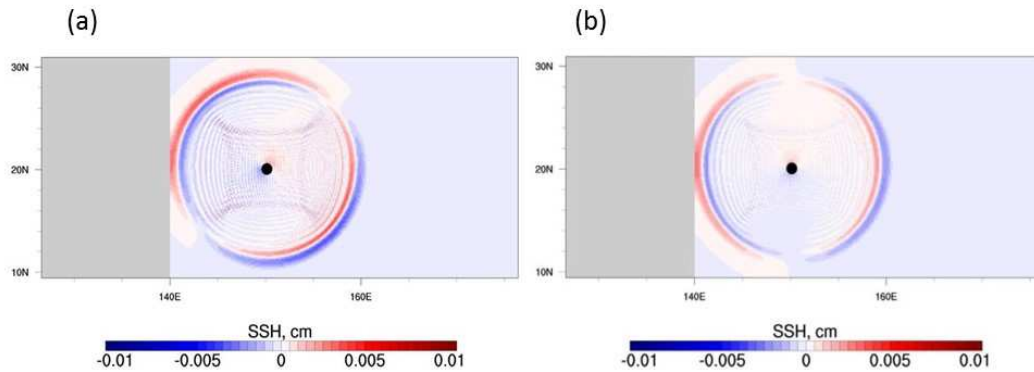


Figure 4: Sea surface height at 10 hours after introducing a rain stress perturbation for (a) the existing implementation of rain stress in ROMS and (b) the corrected implementation of rain stress in ROMS. The black dot in the centre of the domain ( $20^{\circ}\text{N}$   $150^{\circ}\text{E}$ ) denotes the initial position of perturbation. Wind forcing is set to zero.

Next, the contribution of rain stress to storm surges is investigated by forcing ROMS with an idealized tropical cyclone simulated using WRF. Here, the wind and rainfall fields are retained, but the other forcing fields are set to a nominal and constant value in order not to introduce interfering signals into the system. To show the isolated effect of rain stress, the response of SSH and sea surface currents from the rain stress simulation was subtracted from a control simulation (Figure 5). Here, an anticlockwise circulation centering  $20^{\circ}\text{N}$   $141^{\circ}\text{E}$  is generated with a decrease in SSH at the circulation centre. The positive surge is up to  $+0.44$  cm and the negative surge is up to  $-0.63$  cm.



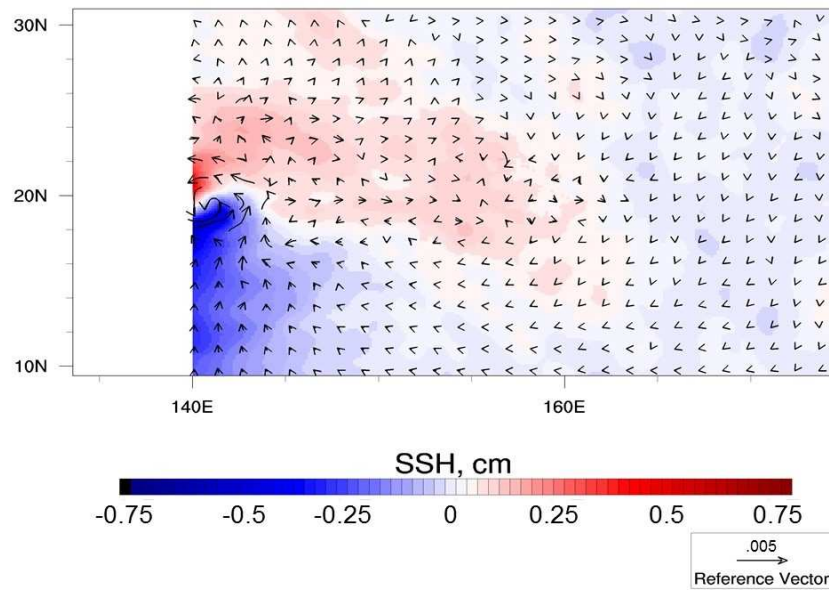


Figure 5: Response of SSH and sea surface currents to rain stress at  $t = 163$  h.

### 3.3 Case study: validation of Cyclone Monica

The atmospheric state is simulated for the period between 20 April 2006 and 24 April 2006. Minimum sea level pressure from WRF is used to generate the cyclone path for comparison as shown in Figure 6. This cyclone path is validated using the storm track provided by the Australian Bureau of Meteorology (BOM, <http://www.bom.gov.au/cyclone>). On 20 April 2006, both the modelled cyclone and the observed cyclone are located at the western coast of Cape York Peninsula, but the modelled cyclone is situated more northwards compared to the observed cyclone. The modelled path is largely similar to the observed path, but has a northward offset of up to  $1^\circ$ . The modelled cyclone also travels faster than the observed cyclone by up to a day.



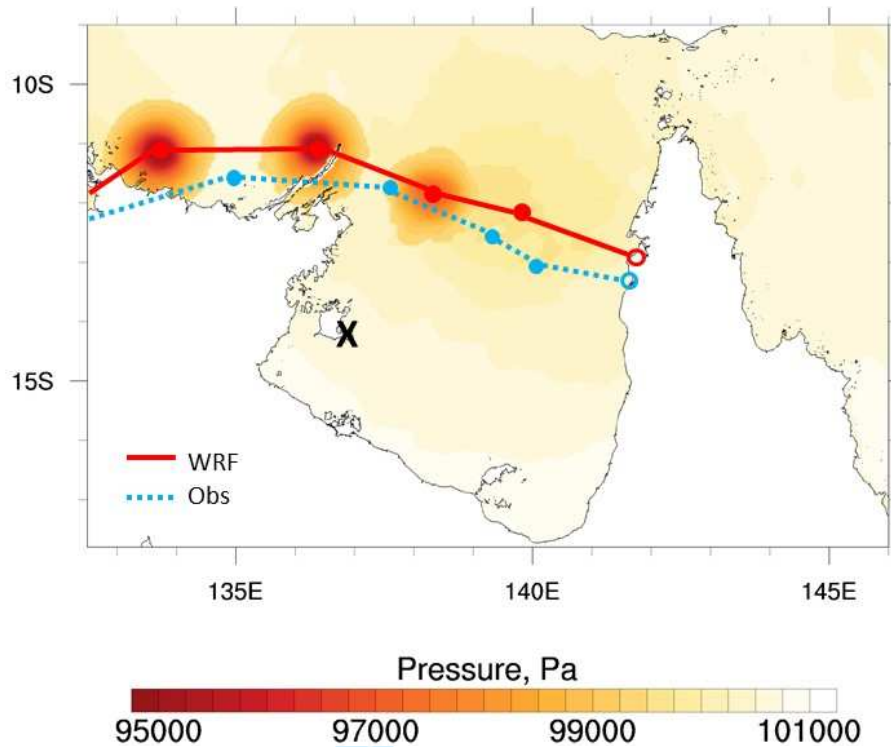


Figure 6: Cyclone path of WRF (solid line) and observations from BOM cyclone report (dashed line). Filled circles are positions every 24 hours from 0000 UTC 20 April 2006 (open circle). Location of the Groote Elyandt weather station is indicated with an 'X'.

Based on the observed weather conditions provided by BOM, the maximum wind speed observed at Cape Wessel (11°S, 136.8°E) and Maningrida (12°S, 134.2°E) were 130 km h<sup>-1</sup> and 148 km h<sup>-1</sup> respectively. In WRF, the maximum modelled wind speed at Cape Wessel and Maningrida were 141 km h<sup>-1</sup> and 146 km h<sup>-1</sup> respectively, which compared well with observations (Figure 7).

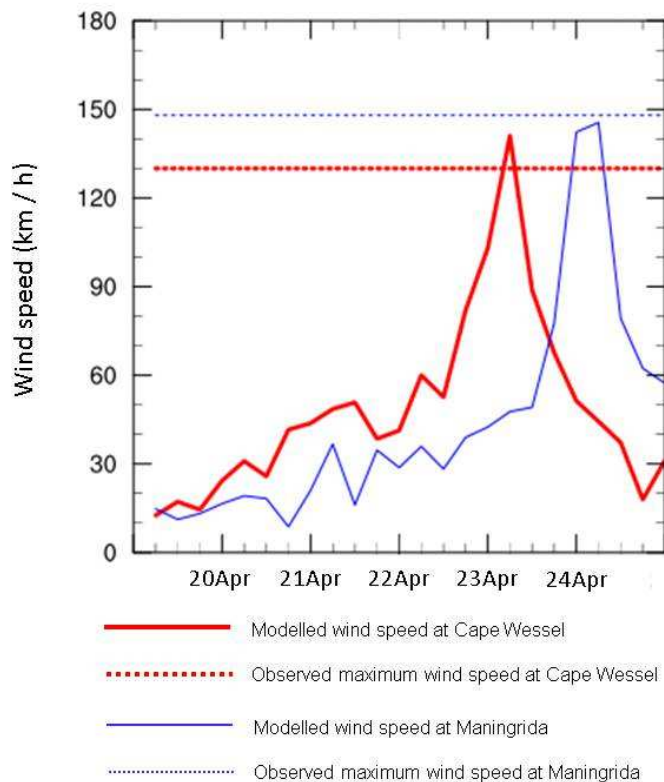
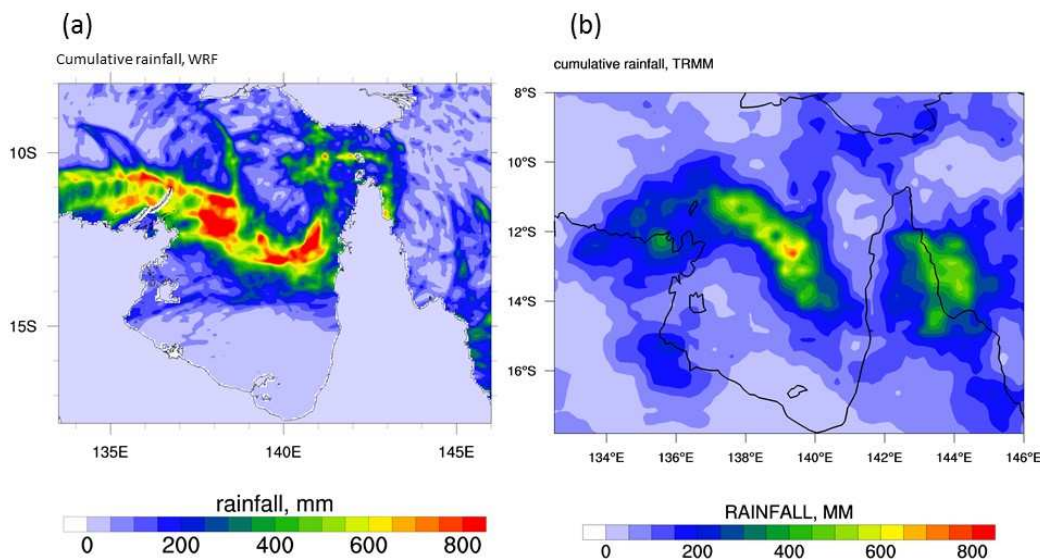


Figure 7: Comparison of the modelled wind speeds at Cape Wessel (11°S, 136.8°E) and Maningrida (12°S, 134.2°E) with the maximum wind speeds observed.

The simulated precipitation is compared with the multi-satellite rainfall estimates from the Tropical Rainfall Measuring Mission (TRMM) 3B42 version 6 (Huffman et al., 2007) (Figure 8). The modelled rainfall path is largely in line with the observed rainfall path but the rainfall is more intense along the cyclone track and less intense away from the track. The observed cumulative rainfall (TRMM) is more homogeneous, having more regions with 100 – 300 mm of rainfall. Nevertheless, the domain cumulative rainfall in WRF is 6% more than that observed by TRMM, although this percentage is expected to be higher since the WRF rainfall is summed with the land mask applied, while TRMM rainfall is summed over every cell in the domain. The total domain cumulative rainfall simulated using WRF is  $4.7 \times 10^6$  mm and the observed cumulative rainfall in TRMM is  $4.4 \times 10^6$  mm over 8 days. However, this is not

1  
2  
3 a large concern in this case study since the main objectives are to simulate and compare  
4 various processes in a realistic environment and not to closely replicate the exact storm event.  
5  
6  
7 The precipitation rates modelled here are not uncommon among tropical cyclones in  
8  
9  
10 Australia and similar storm systems in other parts of the world. From the BOM reports on  
11 historical tropical cyclones, the highest observed precipitation rates of Monica was 340 mm  
12 day<sup>-1</sup>, Laurence (December 2009) at 402 mm day<sup>-1</sup>, Paul (March 2010) at 443 mm day<sup>-1</sup> and  
13  
14  
15 Yasi (January 2011) at 471 mm day<sup>-1</sup>. Outside of Australia, the estimated maximum rainfall  
16  
17 of Hurricane Katrina (August 2005) was 300 mm day<sup>-1</sup> (Dodla et al., 2011) and Typhoon  
18  
19  
20 Morakot (August 2009) was 741 mm day<sup>-1</sup> (Wu, 2013). Hence, the rainfall produced by  
21  
22  
23 Monica is not an extreme case compared to other tropical cyclones and the overestimation in  
24  
25  
26 the rainfall simulated here does not lead to an unrealistic scenario.  
27  
28



49 Figure 8: 8-Day cumulative rainfall output from 19 April 2006 for (a) WRF and (b) TRMM.  
50  
51

52  
53 The configuration of ROMS used for validating the ocean state incorporated the effects of  
54 rain mass, rain stress, the inverse barometer effect and all atmospheric forcings including  
55 wind. The storm surge (Figure 9) is validated against the tidal gauge measurements taken  
56  
57  
58  
59  
60

1  
2  
3 from Milner Bay in Groote Elyandt. The tidal gauge data is obtained from the Australian  
4  
5 Baseline Sea Level Monitoring Project (ABSLMP) (Watson, 2011). To compare the storm  
6  
7 surge, the tidal components are detided using the tidal package by Codiga (2011). The tidal  
8  
9 range is very large and there remains a residual tidal-like signal after de-tiding eleven tidal  
10  
11 components that were introduced into ROMS in the tidal forcing set-up. At this location, the  
12  
13 surge simulated in ROMS (23 April 2006) is earlier than the observed data (24 April 2006)  
14  
15 by a day. However, the magnitude of the surge is very well-captured. The difference in  
16  
17 modelled and observed time of peak surge is also consistent with the different translation  
18  
19 speed and track of the WRF cyclone. The modelled cyclone is located north of the Groote  
20  
21 Elyandt station around 1200 UTC 22 April 2006, while the observed cyclone reaches north of  
22  
23 the Groote Elyandt station around 0000 UTC 24 April 2006. However, the modelled cyclone  
24  
25 path is situated more north compared to the observed cyclone path and a longer time is  
26  
27 expected for the storm surge to be picked up at Groote Elyandt. Hence, an overall difference  
28  
29 of a day can be expected between the modelled and observed storm surge.  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

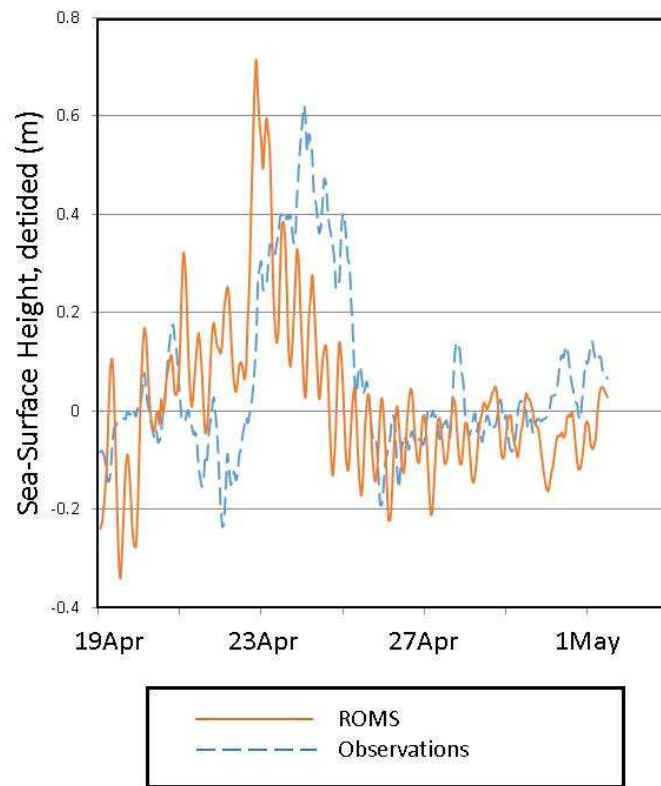


Figure 9: Comparison between the storm surge generated in ROMS with tidal gauge measurements (detided) at Groote Elyandt weather station obtained from ABSLMP, with the location indicated with an 'X' on Figure 6.

### 3.4 Case study: contribution of rain to SSH

We examine the individual contributions to the SSH response to Cyclone Monica. The SSH response to the effect of wind (Figure 10a) is the largest compared to the other effects. On 21 April 2006, the negative surge is up to -58 cm at the south-eastern boundary of the Gulf of Carpentaria and the positive surge is up to 61 cm on the western boundary. On 22 April 2006, the negative surge at the south-eastern boundary of the gulf and the positive surge at the western boundary still persist. There is a region of negative SSH located at the centre of the

1  
2  
3 gulf. The minimum SSH is -57 cm and maximum SSH is 60 cm. On 23 April 2006, the  
4  
5 cyclone reaches the Wessel Islands and the SSH rises up to 171 cm in Buckingham Bay  
6  
7 (12°S 136°E). The positive surge at the western boundary of the gulf is still present at this  
8  
9 time, but the negative surge at the south-eastern boundary has reverted to its nominal levels.  
10  
11 At the northern coast of the Arnhem Land, the negative surge is up to -111 cm.  
12  
13

14  
15 Figure 10b shows the SSH change due to the inverse barometer effect. The SSH increase is  
16  
17 localized to the position of the cyclone and increases as the cyclone intensifies through the  
18  
19 event. The distribution of the SSH scalar field is similar to atmospheric pressure, with the  
20  
21 highest SSH at the centre of the cyclone, gradually decreasing with distance from the cyclone  
22  
23 centre. The SSH increases from 13 cm on 21 April 2006 to 61.5 cm on 23 April 2006 as the  
24  
25 cyclone intensifies.  
26  
27

28  
29 The additional SSH response due to the effect of rain stress is shown in Figure 10c. On 21  
30  
31 April 2006, the additional negative surge is up to -2.4 cm on west coast of Cape York  
32  
33 Peninsula. There is an additional positive SSH of up to 0.9 cm near the position of the  
34  
35 cyclone centre. On 22 April 2006, when Monica is at the centre of the Gulf of Carpentaria,  
36  
37 there is an additional positive SSH of up to 0.7 cm ahead of the cyclone and a negative SSH  
38  
39 of up to -1.1 cm behind the cyclone. On 23 April 2006, there is an additional positive surge of  
40  
41 up to 7.5 cm to the east of the Wessel islands and an additional negative surge of up to -2.9  
42  
43 cm to the west of the islands.  
44  
45

46  
47  
48 Figure 10d shows the additional SSH response due to rain mass, increasing as the cyclone  
49  
50 tracks across the Gulf of Carpentaria and decreasing when the cyclone passes. On 21 April  
51  
52 2006, Monica is situated at the western coast of Cape York Peninsula and the SSH rises up to  
53  
54 3.7 cm at the coast. On 22 April 2006, Monica is centred in the gulf and the SSH increases up  
55  
56 to 4.6 cm, spreading out to the rest of the gulf. The sea level at the southern boundary of the  
57  
58  
59  
60



gulf increases up to 3.5 cm. On 23 April 2006, the cyclone reaches the Wessel Islands and the SSH rises up to an additional 6.4 cm. The sea level change at the southern boundary of the gulf decreases to 2 cm.

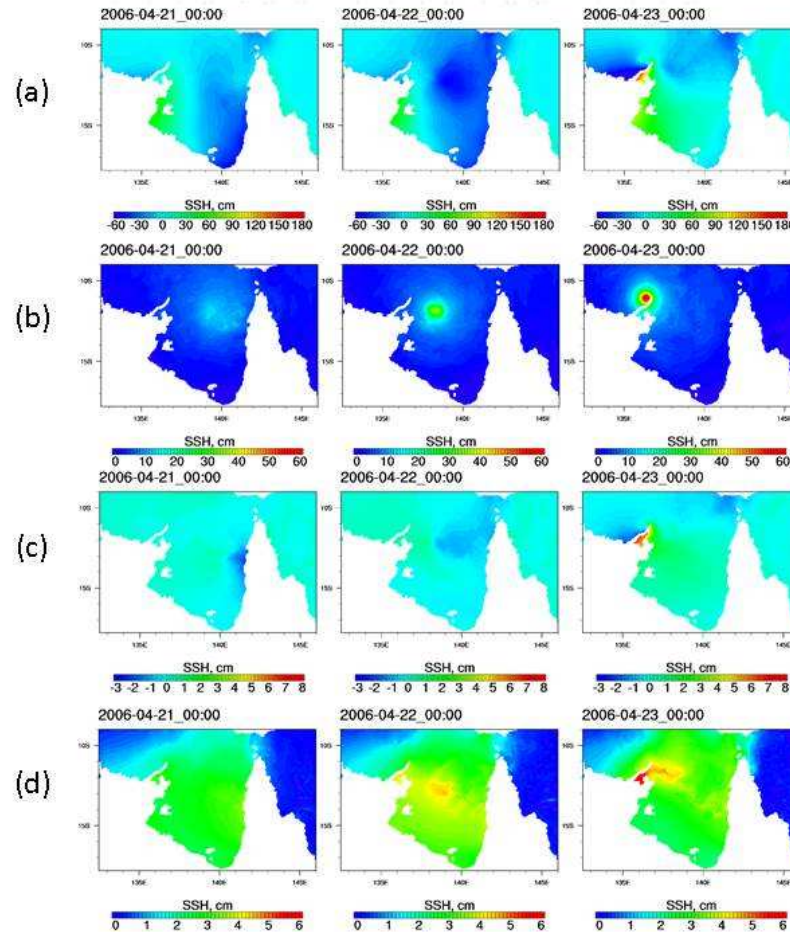


Figure 10: Additional SSH response to the effect of (a) wind, (b) atmospheric pressure, (c) rain stress and (d) rain mass from 21 April 2006 to 23 April 2006.

#### 4 Discussion

The rain mass perturbation on the water surface sets off a surface gravity wave propagating outwards towards the domain boundaries over a time scale of several hours (Figure 1).

1  
2  
3 Characteristics of surface waves are inherently non-linear. But considering shallow water  
4 approximations (Stewart, 2009), where the water depth is much smaller than the wavelength,  
5 the phase speed ( $c$ ) can be expressed as:  
6  
7  
8

$$c = \sqrt{gd} \quad (5)$$

9  
10  
11  
12  
13  
14  
15  
16 where  $g$  is the gravitational acceleration and  $d$  is the depth of water.

17  
18  
19  
20 From the initial point of perturbation, the surface wave propagated from 20°N 150°E to 20°N  
21 160°E in 9 hours translating at a speed of 30.9 m s<sup>-1</sup>. The shallow water approximation is  
22 valid for water depths much lesser than the wavelength. Here, the water depth is 100 m while  
23 the wavelength is about 100 km. The water depth in this case is only 0.1% that of the  
24 wavelength, making the shallow water approximation a valid assumption. From equation (5),  
25 the surface gravity wave in this case should advance at a phase speed of 31.3 m s<sup>-1</sup>, which is  
26 consistent with the ROMS output.  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37

38 Next, the SSH response at a longer time scale is examined using an idealized tropical cyclone  
39 to force the ocean model over 8 days. In Figure 3, the added mass spreads outwards from the  
40 mass sources towards the domain boundaries via gravity waves. In the deeper ocean case, the  
41 maximum SSH is smaller (up to 3.2 cm) compared to the shallow ocean case (up to 5.3 cm).  
42 From equation (5), gravity waves in the deeper case travel 3 times faster than in the shallow  
43 water case, which accounts for the faster decrease in the maximum SSH for the deeper ocean  
44 case. With the rain mass spreading out faster in the deeper ocean, the minimum SSH reaches  
45 a higher level for the deep ocean case (2.3 cm) compared to the shallow ocean case (1.3 cm).  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56 At longer time scales (after 7 days), a pressure gradient is generated from the sea surface  
57  
58  
59  
60



1  
2  
3 height increase. This creates a geostrophic adjustment with an opposing Coriolis force.

4  
5 Geostrophic currents of the magnitude  $0.01 \text{ m s}^{-1}$  are generated, rotating in a clockwise  
6  
7 direction.  
8  
9

10  
11 With the implementation of rain stress now corrected in ROMS, the SSH response to rain  
12 stress is shown by forcing ROMS with an idealized tropical cyclone (Figure 5). Rain stress is  
13 parameterized as a function of rain rate and wind speed and follows the direction of the  
14 winds. Hence, rain stress is only exerted in areas where non-zero values of rain rate and wind  
15 speed coincide. In the case of a tropical cyclone, rainfall is highest along the rain bands  
16 towards the centre of the cyclone. Rain stress contributes to storm surges in a manner similar  
17 to wind stress. The anti-clockwise rotating rain stress creates an anti-clockwise gyre at the sea  
18 surface. Coriolis force deflects the currents to the right and generates an outward diverging  
19 flow that decreases the SSH at the centre of the gyre.  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31

32  
33 Finally, a real case study is investigated to compare the SSH response to different storm surge  
34 processes. The effect of wind (Figure 10a) is the largest compared to the effect of  
35 precipitation and the inverse barometer effect. It drives positive and negative SSH anomalies  
36 based on the wind direction and the geometry of the coastline. On 21 April 2006, the strong  
37 easterly wind to the south of the domain generates a positive surge at the western boundary of  
38 the Gulf of Carpentaria due to the onshore wind. A negative surge is generated at the south-  
39 eastern boundary of the gulf due to the offshore wind. On 22 April 2006, the easterly wind to  
40 the south of the domain continues to generate a positive and negative surge at the western and  
41 south-eastern boundaries of the gulf respectively. There is a SSH depression in the middle of  
42 the gulf due to Ekman pumping. This SSH depression is present in all the 6-hourly output  
43 from 1800 UTC 21 April 2006 to 1800 UTC 22 April 2006 when Monica is trekking across  
44 the Gulf of Carpentaria. The SSH depression is absent on 21 April 2006 and 23 April 2006  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

1  
2  
3 when the cyclone is near Cape York Peninsula and the Wessel Islands respectively. On 23  
4 April 2006, the cyclone is relatively far from the eastern boundary of the gulf and the  
5 offshore winds here weakens from the range of  $15 \text{ m s}^{-1}$  (21 April 2006) to the range of  $10 \text{ m}$   
6  $\text{s}^{-1}$ . The negative surge at the eastern boundary of the gulf weakens correspondingly. On the  
7 other hand, the positive surge at the western boundary of the gulf strengthens with the strong  
8 onshore winds. At this time, the strong winds blowing into Buckingham Bay create a  
9 funnelling effect, channelling water into the bay and generating a high positive surge of up to  
10 171 cm.  
11  
12  
13  
14  
15  
16  
17  
18  
19

20  
21 The SSH response to the inverse barometer effect is spatially constrained by the location of  
22 the cyclone's low pressure centre and varies in magnitude based on the storm intensity  
23 (Figure 10b). SSH is expected to increase by 1 cm for every 1 mbar drop in atmospheric  
24 pressure (Roden & Rossby, 1999). The inverse barometer effect contributes substantially to  
25 the storm surge when the low pressure centre of the cyclone coincides with the coastline.  
26  
27 When the cyclone is away from the coastline and over the open sea, the inverse barometer  
28 effect raises the SSH but does not contribute significantly to the storm surge. On 21 April  
29 2006, Monica crosses the west coast of Cape York Peninsula but due to its low intensity at  
30 this time, the contribution to the storm surge is only up to 13 cm. On 22 April 2006, the storm  
31 intensifies as it crosses the Gulf of Carpentaria but due to its distance from the coast, the  
32 contribution to the storm surge is only around 10 cm, despite raising the SSH up to 35.2 cm  
33 in the centre of the gulf. On 23 April 2006, Monica intensifies further, crossing the Wessel  
34 Islands and increases the storm surge by up to 61.5 cm.  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49

50  
51 A region is subjected to rain stress when rainfall and winds coincide. On 21 April 2006, the  
52 location of intense rain is situated just off the west coast of Cape York Peninsula (Figure  
53 10c). The cyclone is traversing westward with an overall westward wind direction. Rain  
54  
55  
56  
57  
58  
59  
60

1  
2  
3 stress perturbation on the sea surface creates a positive front ahead of the perturbation and a  
4 negative front behind it as shown in Figure 4. The offshore winds and high rainfall at the  
5 eastern boundary of the gulf generates a SSH response to rain stress, enhancing the negative  
6 surge here. While the wind stress generated positive and negative surges at the western and  
7 south-eastern boundaries of the Gulf of Carpentaria, these surges are absent in the rain stress  
8 effect shown here due to the lower rainfall at these locations. On 22 April 2006, the cyclone  
9 is in the middle of the gulf trekking westwards. This creates a similar effect with the positive  
10 SSH ahead of the cyclone and a negative SSH behind it. The positive and negative surges at  
11 the western and south-eastern boundaries of the gulf are similarly absent here, due to the light  
12 rainfall in these areas. On 23 April 2006 when the cyclone crosses the Wessel Islands, the  
13 effect of rain stress is very similar to that of wind stress, with the intense rain coinciding with  
14 the strong winds here. The additional positive surge in Buckingham Bay is generated due to  
15 rain stress channelling water into the bay, which is similar to the effect of wind stress. At the  
16 northern coast of the Arnhem Land just west of the Wessel Islands, the offshore winds  
17 coincide with the high rainfall from the storm to generate an offshore rain stress and a  
18 negative surge of up to -2.9 cm.  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38

39 Figure 10d shows the SSH response to precipitation mass addition. From 21 April 2006 to 22  
40 April 2006, the areas of SSH change do not confine solely to the precipitation footprint (as  
41 traced out in Figure 8a), but spreads out, advecting to the south of the Gulf of Carpentaria. On  
42 23 April 2006, the cyclone crosses the Wessel Islands and the rain mass is trapped in  
43 Buckingham Bay, generating an additional local SSH increase of 6.4 cm. The effect of  
44 precipitation is essentially an in-situ addition of mass, and inlets and semi-enclosed basins  
45 will serve to enhance this effect by trapping the rain. Buckingham Bay is fairly well-resolved  
46 in this 6 km resolution model, spanning 13 x 14 grid cells, so the size of the effect is unlikely  
47 to be just numerical. The geometry of the gulf itself is an important contributing factor  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

1  
2  
3 towards the effect of precipitation seen here. From Figure 10d, it can be seen that the SSH  
4  
5 response is generally higher to the southern half of the gulf than the north. If the gulf was  
6  
7 entirely enclosed, the additional mass will be expected to distribute evenly by the clockwise  
8  
9 gyre in both the northern and southern regions. The constrained southern boundary of the gulf  
10  
11 enhances the build-up of SSH within the gulf, since the north-western section of the gulf  
12  
13 opens up to the Arafura Sea. Typical remote enhancements are of the order of 2 cm.  
14  
15

16  
17 Evaporation, often considered in tandem with precipitation, can have notable effects on sea  
18  
19 surface elevation. Roden et al. (1999) have reported oceanic evaporation to be between 0.08  
20  
21 cm day<sup>-1</sup> to 0.55 cm day<sup>-1</sup>, depending on wind conditions. Hurricane evaporation rates have  
22  
23 been found to be in the order of 1 – 2 cm day<sup>-1</sup> (Riehl and Malkus (1958), Machta (1968),  
24  
25 etc). At this order of magnitude, the evaporation rate is much smaller compared to the effect  
26  
27 of rainfall in this case, with local effects of up to 14 cm (mass and momentum transfer).  
28  
29 Nevertheless, it will be interesting to include the effect of evaporation rate on sea surface  
30  
31 height in future work for a comprehensive comparison.  
32  
33  
34  
35

36  
37 The Gulf of Carpentaria is not an isolated case of a semi-enclosed basin that is prone to  
38  
39 severe storms and many other cases have been studied. Zhong et al. (2010) used ROMS to  
40  
41 study the storm surge predictions during the passage of Hurricane Isabel in Chesapeake Bay  
42  
43 and found that the hurricane translation speed and the resolution of the horizontal wind field  
44  
45 are important factors affecting storm surges. Rego et al. (2010) used FVCOM to simulate  
46  
47 Hurricane Ike at Galveston Bay and found that the Bolivar Peninsula provided a significant  
48  
49 surge barrier in protecting the bay. Westerink et al. (2008) developed a model of south  
50  
51 Louisiana (which included many inlets, bays and channels) using ADCIRC to simulate the  
52  
53 effect of storm surges during the events Hurricane Betsy and Andrew. Weisberg et al. (2006)  
54  
55 used FVCOM to simulate the storm surge produced from several idealized hurricanes  
56  
57  
58  
59  
60

1  
2  
3 crossing Tampa Bay and found that the hurricanes with slow transitional speeds produced  
4  
5 larger surges within the bay. However, none of these studies or models include the effect of  
6  
7 rain mass. The semi-enclosed bay comprising Buckingham Bay, extending to the Wessel  
8  
9 Islands and Gove Peninsula ( $12^{\circ}2'S$   $136^{\circ}5'E$ ) consists of about 250 cells, which is equivalent  
10  
11 to an area of  $9000\text{ km}^2$ . The areas investigated in literature, such as Tampa Bay ( $5700\text{ km}^2$ ),  
12  
13 Galveston Bay ( $1600\text{ km}^2$ ) and Chesapeake Bay ( $11600\text{ km}^2$ ) are all of similar sizes or  
14  
15 smaller. Not including the effect of rain mass in the models used in these studies will imply a  
16  
17 systematic error in the storm surge levels modelled. To further compare how each of the  
18  
19 physical effects contribute to coastal impacts in a worst case scenario, the SSH at the  
20  
21 coastline starting from the Wessel Islands and ending at the northern tip of Cape York  
22  
23 Peninsula is investigated. The maximum storm surge (for the duration of the storm) due to  
24  
25 each physical effect is shown against the coastal locations (Figure 11) to illustrate the worst  
26  
27 case scenario that can arise from each effect during this event. Considering the maximum  
28  
29 storm surge response to the different processes at each coastal point, the additional effect of  
30  
31 rain mass ranges from 1.3 cm to 6.4 cm, the effect of rain stress is from 0.4 cm to 7.5 cm, the  
32  
33 inverse barometer effect is from 0.2 cm to 45.4 cm and the effect of wind stress is from 20.5  
34  
35 cm to 170.6 cm. While the inverse barometer effect is generally larger than the effect of rain  
36  
37 mass, it falls below the effect of rain mass at the southern boundary of the gulf (coastal points  
38  
39 180 – 330) when the cyclone is located away from the coastline. This shows that the effect of  
40  
41 rain mass can have a stronger remote presence than the inverse barometer effect. Although  
42  
43 the maximum surge level due to the effect of rain stress is higher than the effect of rain mass,  
44  
45 the effect of rain mass is larger than the effect of rain stress in most of the coastal points.  
46  
47 ROMS (and many ocean models) include the effect of rain stress but not the rain mass even  
48  
49 when the effect of rain mass can be larger as shown here. Ocean models that consider the  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

effect of rain stress to be sufficiently substantial to include its effect should not neglect the effect of rain mass.

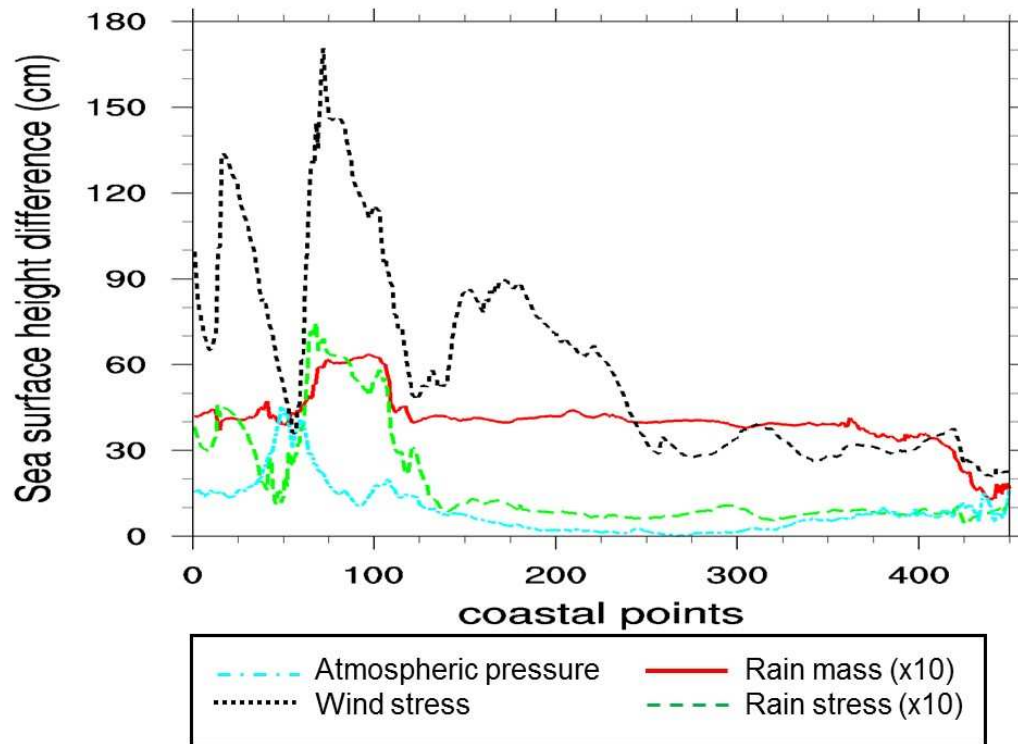


Figure 11: Comparison of the SSH change due to various physical effects, along the coastline at Gulf of Carpentaria. The coastal points can be attributed to different locations along the coastline as follows: the Wessel Islands (points 1 – 120), the western boundary of the Gulf of Carpentaria (points 121 – 180), the southern boundary of the gulf (points 181 – 280) and the eastern boundary of the gulf (points 281 – 450). SSH responses to rain mass and rain stress are scaled up by a factor of 10 for illustration.

## 5 Conclusion

Having successfully incorporated a new rain scheme into ROMS that considers the mass contribution of rainfall, its effect is shown in several cases. This new rain scheme does not interfere with the existing temperature and salinity treatment of rainfall in ROMS. The added rain mass generated surface waves that increase in propagation speed with increasing depth. An error in the existing implementation of rain stress in ROMS produced a systematic overestimation of 41% in magnitude and up to 45° in direction. This error has been corrected in this study and its isolated effect on SSH in a tropical cyclone case is shown in an idealized set-up. Finally, the effect of rain mass and rain stress in a real storm surge case (Cyclone Monica) is compared with the effect of winds and atmospheric pressure. In this case, the effect of wind and pressure generated a surge of up to 170.6 cm and 61.5 cm respectively. The effect of rain mass increases the SSH up to 6.4 cm in Buckingham Bay. The effect of rain mass has a stronger remote presence compared to the effect of pressure and rain stress. Especially in semi-enclosed areas, the effect of rain mass can substantially contribute to the storm surge even when the storm is relatively far from the coastline, and both the effect of pressure and rain stress are no longer significant. An area is subjected to high rain stress when intense winds coincide with heavy rain. Here, the effect of rain stress increases the SSH up to 7.5 cm, but the effect is generally smaller than the effect of rain mass along the coast. Many ocean models ignore the mass effect of rainfall and only incorporate the effect of wind, atmospheric pressure and rain stress. However, this study shows that the effect of rain mass can be larger than the effect of pressure and rain stress in certain scenarios, and should not be neglected. Each process shows different spatial influence and cannot be accounted for by simply scaling up a certain effect, for example, recalibrating the surge model by scaling the wind field.

## References

- 1  
2  
3  
4  
5  
6  
7 Bleck, R., & Boudra, D. B. (1981). Initial Testing of a Numerical Ocean Circulation Model  
8  
9 Using a Hybrid (Quasi-Isopycnic) Vertical Coordinate. *Journal of Physical*  
10  
11 *Oceanography*, 11(6), 755-770. doi: Doi 10.1175/1520-  
12  
13 0485(1981)011<0755:Itoano>2.0.Co;2  
14  
15  
16 Caldwell, D. R., & Elliott, W. P. (1970). Surface stresses produced by rainfall. *J. Phys.*  
17  
18 *Oceanography*, 1, 145-148.  
19  
20  
21 Chapman, D. C. (1985). Numerical treatment of cross-shelf open boundaries in a barotropic  
22  
23 coastal ocean model. *Journal of Physical Oceanography*, 15(8), 1060-1075. doi:  
24  
25 10.1175/1520-0485(1985)015<1060:ntocso>2.0.co;2  
26  
27  
28 Chen, C. S., Qi, J. H., Li, C. Y., Beardsley, R. C., Lin, H. C., Walker, R., & Gates, K. (2008).  
29  
30 Complexity of the flooding/drying process in an estuarine tidal-creek salt-marsh  
31  
32 system: An application of FVCOM. *Journal of Geophysical Research-Oceans*,  
33  
34 113(C7). doi: Artn C07052 Doi 10.1029/2007jc004328  
35  
36  
37 Chou, M., & Suarez, and M. J. (1994). An efficient thermal infrared radiation  
38  
39 parameterization for use in general circulation models. *NASA Tech. Memo.*,  
40  
41 104606(3), 85pp.  
42  
43  
44 Clayson, C. A., & Kantha, L. H. (1999). Turbulent kinetic energy and its dissipation rate in  
45  
46 the equatorial mixed layer. *Journal of Physical Oceanography*, 29(9), 2146-2166. doi:  
47  
48 Doi 10.1175/1520-0485(1999)029<2146:Tkeaid>2.0.Co;2  
49  
50  
51 Codiga, D.L. (2011). Unified Tidal Analysis and Prediction Using the UTide Matlab  
52  
53 Functions. *Technical Report 2011-01 Graduate School of Oceanography, University*  
54  
55 *of Rhode Island, Narragansett, RI. 59pp*  
56  
57  
58  
59  
60



- 1  
2  
3 Dodla, V. B., Desamsetti, S., & Yerramilli, A. (2011). A Comparison of HWRF, ARW and  
4  
5 NMM Models in Hurricane Katrina (2005) Simulation. *International Journal of*  
6  
7 *Environmental Research and Public Health*, 8(6), 2447-2469. doi: DOI  
8  
9 10.3390/ijerph8062447  
10  
11  
12 Dudhia, J. (1989). Numerical Study of Convection Observed during the Winter Monsoon  
13  
14 Experiment Using a Mesoscale Two-Dimensional Model. *Journal of the Atmospheric*  
15  
16 *Sciences*, 46(20), 3077-3107. doi: Doi 10.1175/1520-  
17  
18 0469(1989)046<3077:Nsocod>2.0.Co;2  
19  
20  
21 Dudhia, Jimy. (1996). A multi-layer soil temperature model for MM5. *Sixth PSU/NCAR*  
22  
23 *Mesoscale Model Users' Workshop*.  
24  
25  
26 Durden, S. L. (2010). Remote Sensing and Modeling of Cyclone Monica near Peak Intensity.  
27  
28 *Atmosphere*, 1(1), 15-33. doi: Doi 10.3390/Atmos1010015  
29  
30  
31 Egbert, G. D., Bennett, A. F., & Foreman, M. G. G. (1994). TOPEX/POSEIDON tides  
32  
33 estimated using a global inverse model. *Journal of Geophysical Research-Oceans*,  
34  
35 99(C12), 24821-24852. doi: 10.1029/94jc01894  
36  
37  
38 Fairall, C. W., Bradley, E. F., Rogers, D. P., Edson, J. B., & Young, G. S. (1996). Bulk  
39  
40 parameterization of air-sea fluxes for Tropical Ocean Global Atmosphere Coupled  
41  
42 Ocean Atmosphere Response Experiment. *Journal of Geophysical Research-Oceans*,  
43  
44 101(C2), 3747-3764. doi: Doi 10.1029/95jc03205  
45  
46  
47  
48 Flather, R. A. (1976). A tidal model of the northwest European continental shelf *Mémoires de*  
49  
50 *la Société Royale de Sciences de Liege* (Vol. 6, pp. 141-164).  
51  
52  
53  
54 Gill, A. E., and P. P. Niiler (1973). The theory of the seasonal variability in the ocean, *Deep*  
55  
56 *Sea Res.*, 20, 141-177  
57  
58  
59  
60

- 1  
2  
3 Griffies, S. M., & Hallberg, R. W. (2000). Biharmonic friction with a Smagorinsky-like  
4  
5 viscosity for use in large-scale eddy-permitting ocean models. *Monthly Weather*  
6  
7 *Review*, 128(8), 2935-2946.  
8  
9
- 10 Hirose, N., I. Fukumori, and R. M. Ponte (2001). A non-isostatic global sea level response to  
11  
12 barometric pressure near 5 days, *Geophys. Res. Lett.*, 28, 2441–2444,  
13  
14 doi:10.1029/2001GL012907.  
15  
16
- 17 Hong, S. Y., Noh, Y., & Dudhia, J. (2006). A new vertical diffusion package with an explicit  
18  
19 treatment of entrainment processes. *Monthly Weather Review*, 134(9), 2318-2341.  
20  
21 doi: Doi 10.1175/Mwr3199.1  
22  
23
- 24 Huffman, G. J., Adler, R. F., Bolvin, D. T., Gu, G. J., Nelkin, E. J., Bowman, K. P., . . .  
25  
26 Wolff, D. B. (2007). The TRMM multisatellite precipitation analysis (TMPA): Quasi-  
27  
28 global, multiyear, combined-sensor precipitation estimates at fine scales. *Journal of*  
29  
30 *Hydrometeorology*, 8(1), 38-55. doi: Doi 10.1175/Jhm560.1  
31  
32
- 33 Janjic, Z. I. (1994). The Step-Mountain Eta Coordinate Model - Further Developments of the  
34  
35 Convection, Viscous Sublayer, and Turbulence Closure Schemes. *Monthly Weather*  
36  
37 *Review*, 122(5), 927-945. doi: Doi 10.1175/1520-  
38  
39 0493(1994)122<0927:Tsmecm>2.0.Co;2  
40  
41  
42
- 43 Jarvinen, B. R., & Lawrence, M. B. (1985). An Evaluation of the Slosh Storm-Surge Model.  
44  
45 *Bulletin of the American Meteorological Society*, 66(11), 1408-1411.  
46  
47
- 48 Kain, J. S. (2004). The Kain-Fritsch convective parameterization: An update. *Journal of*  
49  
50 *Applied Meteorology*, 43(1), 170-181. doi: Doi 10.1175/1520-  
51  
52 0450(2004)043<0170:Tkcpan>2.0.Co;2  
53  
54  
55  
56  
57  
58  
59  
60

- 1  
2  
3 Lin, Y. L., Farley, R. D., & Orville, H. D. (1983). Bulk Parameterization of the Snow Field in  
4  
5 a Cloud Model. *Journal of Climate and Applied Meteorology*, 22(6), 1065-1092. doi:  
6  
7 Doi 10.1175/1520-0450(1983)022<1065:Bpotsf>2.0.Co;2  
8  
9
- 10 Machta, L. (1969). Evaporation rates based on tritium measurements for hurricane Betsy.  
11  
12 *Tellus*, 21(3), 404-408.  
13  
14
- 15 Marchesiello, P., McWilliams, J. C. , & Shchepetkin, A. F. . (2001). Open boundary  
16  
17 conditions for long-term integration of regional ocean models. *Ocean Modelling*, 3, 1-  
18  
19 20.  
20  
21
- 22 Mlawer, E. J., Taubman, S. J., Brown, P. D., Iacono, M. J., & Clough, S. A. (1997). Radiative  
23  
24 transfer for inhomogeneous atmospheres: RRTM, a validated correlated-k model for  
25  
26 the longwave. *Journal of Geophysical Research-Atmospheres*, 102(D14), 16663-  
27  
28 16682. doi: Doi 10.1029/97jd00237  
29  
30
- 31 Paulson, C. A. (1970). The mathematical representation of wind speed and temperature  
32  
33 profiles in the unstable atmospheric surface layer. *J. Appl. Meteor.*, 9, 857-861.  
34  
35
- 36 Peng, M. C., Xie, L., & Pietrafesa, L. J. (2004). A numerical study of storm surge and  
37  
38 inundation in the Croatan-Albemarle-Pamlico Estuary System. *Estuarine Coastal and*  
39  
40 *Shelf Science*, 59(1), 121-137. doi: DOI 10.1016/j.ecss.2003.07.010  
41  
42
- 43 Pickard, G. L., & Emery, W. J. (1990). *Descriptive physical oceanography: an introduction*.  
44  
45 Elsevier.  
46  
47
- 48 Ponte, R. M. (2006). Oceanic response to surface loading effects neglected in volume-  
49  
50 conserving models. *Journal of Physical Oceanography*, 36(3), 426-434. doi:  
51  
52 10.1175/Jpo2843.1  
53  
54  
55  
56  
57  
58  
59  
60

- 1  
2  
3 Rego, J. L., & Li, C. Y. (2010). Storm surge propagation in Galveston Bay during Hurricane  
4  
5 Ike. *Journal of Marine Systems*, 82(4), 265-279. doi: DOI  
6  
7 10.1016/j.jmarsys.2010.06.001  
8  
9  
10 Riehl H, Malkus J. (2011). Some aspects of hurricane Daisy, 1958. *Tellus A* 13(2)  
11  
12  
13 Roden, G. I., & Rossby, H. T. (1999). Early Swedish contribution to oceanography: Nils  
14  
15 Gissler (1715-71) and the inverted barometer effect. *Bulletin of the American*  
16  
17 *Meteorological Society*, 80(4), 675-682. doi: Doi 10.1175/1520-  
18  
19 0477(1999)080<0675:Escton>2.0.Co;2  
20  
21  
22 Rogers, Eric, , Thomas Black, , Brad Ferrier, , Ying Lin, , David Parrish, & , and Geoffrey  
23  
24 DiMego. (2001). National Oceanic and Atmospheric Administration Changes to the  
25  
26 NCEP Meso Eta Analysis and Forecast System: Increase in resolution, new cloud  
27  
28 microphysics, modified precipitation assimilation, modified 3DVAR analysis. *NOAA*.  
29  
30  
31 Saha, S., Moorthi, S., Pan, H. L., Wu, X. R., Wang, J. D., Nadiga, S., Goldberg, M. (2010).  
32  
33 The Ncep Climate Forecast System Reanalysis. *Bulletin of the American*  
34  
35 *Meteorological Society*, 91(8), 1015-1057. doi: Doi 10.1175/2010bams3001.1  
36  
37  
38 Shchepetkin, A. F., & McWilliams, J. C. (2005). The regional oceanic modeling system  
39  
40 (ROMS): a split-explicit, free-surface, topography-following-coordinate oceanic  
41  
42 model. *Ocean Modelling*, 9(4), 347-404. doi: DOI 10.1016/j.ocemod.2004.08.002  
43  
44  
45  
46 Sheng, Y. P., Alymov, V., & Paramygin, V. A. (2010). Simulation of storm surge, wave,  
47  
48 currents, and inundation in the Outer Banks and Chesapeake Bay during Hurricane  
49  
50 Isabel in 2003: The importance of waves. *Journal of Geophysical Research-Oceans*,  
51  
52 115. doi: Artn C04008 Doi 10.1029/2009jc005402  
53  
54  
55  
56  
57  
58  
59  
60

- 1  
2  
3 Sheng, Y. P., Davis, J. R., Figueiredo, R., Liu, B., Liu, H., Luettich, R., Zheng, L. (2012). A  
4  
5 regional testbed for storm surge and coastal inundation models - an overview.  
6  
7 *Estuarine and Coastal Modeling, American Society of Civil Engineers*, 476-495.  
8  
9  
10 Skamarock, W. C, Klemp, J. B, Dudhia, J, Gill, D. O, Baker, D. M, Duda, M. G, JG., Power.  
11  
12 (2008). A Description of the Advanced Research WRF Version 3. . *NCAR/TN-*  
13  
14 *475+STR*.  
15  
16  
17 Smith, W. H. F., and Sandwell D. T. . (1997). Global seafloor topography from satellite  
18  
19 altimetry and ship depth soundings. *Science*, 277, 1957-1962.  
20  
21  
22 Stewart, Robert H. (2009). Introduction to Physical Oceanography. *University Press of*  
23  
24 *Florida*.  
25  
26  
27 Umlauf, L., & Burchard, H. (2003). A generic length-scale equation for geophysical  
28  
29 turbulence models. *Journal of Marine Research*, 61(2), 235-265. doi:  
30  
31 10.1357/002224003322005087  
32  
33  
34 Van Dorn, W. G. (1953). Wind stress on an artificial pond. *Journal of marine research*, 12,  
35  
36 249-276.  
37  
38  
39 Wang, Shuai, Toumi, Ralf, Czaja, Arnaud, & Van Kan, Adrian. (2015). An Analytic Model  
40  
41 of Tropical Cyclone Wind Profiles. *Quarterly Journal of the Royal Meteorological*  
42  
43 *Society*, n/a-n/a. doi: 10.1002/qj.2586  
44  
45  
46 Watson, P. J. (2011). Is There Evidence Yet of Acceleration in Mean Sea Level Rise around  
47  
48 Mainland Australia? *Journal of Coastal Research*, 27(2), 368-377. doi: Doi  
49  
50 10.2112/Jcoastres-D-10-00141.1  
51  
52  
53 Weisberg, R. H., & Zheng, L. Y. (2006). Hurricane storm surge simulations for Tampa Bay.  
54  
55 *Estuaries and Coasts*, 29(6A), 899-913.  
56  
57  
58  
59  
60

- 1  
2  
3 Westerink, J. J., Luettich, R. A., Baptista, A. M., Scheffner, N. W., & Farrar, P. (1992). Tide  
4  
5 and Storm-Surge Predictions Using Finite-Element Model. *Journal of Hydraulic*  
6  
7 *Engineering-Asce*, 118(10), 1373-1390. doi: Doi 10.1061/(Asce)0733-  
8  
9 9429(1992)118:10(1373)  
10  
11  
12 Westerink, J. J., Luettich, R. A., Feyen, J. C., Atkinson, J. H., Dawson, C., Roberts, H. J.,  
13  
14 Pourtaheri, H. (2008). A basin- to channel-scale unstructured grid hurricane storm  
15  
16 surge model applied to southern Louisiana. *Monthly Weather Review*, 136(3), 833-  
17  
18 864. doi: Doi 10.1175/2007mwr1946.1  
19  
20  
21  
22 Wu, C. C. (2013). TYPHOON MORAKOT Key Findings from the Journal TAO for  
23  
24 Improving Prediction of Extreme Rains at Landfall. *Bulletin of the American*  
25  
26 *Meteorological Society*, 94(2), 155-160. doi: Doi 10.1175/Bams-D-11-00155.1  
27  
28  
29 Zhong, L. J., Li, M., & Zhang, D. L. (2010). How do uncertainties in hurricane model  
30  
31 forecasts affect storm surge predictions in a semi-enclosed bay? *Estuarine Coastal*  
32  
33 *and Shelf Science*, 90(2), 61-72. doi: DOI 10.1016/j.ecss.2010.07.001  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

1  
2  
3 **Effect of extreme ocean precipitation on sea surface elevation and storm surges**  
4

5  
6  
7 Running head: Effect of extreme rainfall on sea surface elevation and storm surges  
8  
9

10  
11  
12  
13 Benjamin Wong (corresponding author): benjamin.wong11@imperial.ac.uk  
14  
15

16  
17 Ralf Toumi: r.toumi@imperial.ac.uk  
18  
19

20  
21  
22  
23 **Space and Atmospheric Physics Group,**  
24

25  
26 **Blackett Laboratory,**  
27

28  
29 **Imperial College London,**  
30

31  
32 **London SW7 2AZ,**  
33

34  
35  
36 **UK**  
37  
38

39  
40  
41  
42 **Keywords:** ocean modelling, Cyclone Monica, rain mass, rain stress  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

## Abstract

Ocean models that neglect mass and momentum contributions from precipitation can have a systematic bias in sea surface height (SSH). Here, a new rainfall scheme is introduced into the Regional Ocean Modelling System (ROMS) to incorporate the effects of precipitation mass. When precipitation is added to the sea surface, it spreads out via surface gravity waves that increase in propagation speed with increasing water depth. Over several days, the SSH increase due to the precipitation mass added created a geostrophic adjustment, generating anti-cyclonic geostrophic currents around the SSH increase. The transfer of momentum from precipitation to the sea surface, or rain stress, can also be important. In the case study of a real tropical cyclone, Monica passing North Australia, the effect of incorporating precipitation mass is compared with other processes affecting the storm surge: surface wind, inverse barometer effect and rain stress. The maximum SSH response is 170.6 cm for the wind effect, 61.5 cm for the inverse barometer effect, 7.5 cm for the effect of rain stress and 6.4 cm for the effect of rain mass. Each process has been shown to have different spatial influences. The effect of rain mass has a strong remote influence compared to the inverse barometer effect and the effect of rain stress. This is particularly seen in semi-enclosed bays.

## 1 Introduction

The contribution of oceanic precipitation to storm surges has so far been neglected in storm surge models, typically only incorporating the contributions of the inverse barometer effect (Roden et al., 1999) and wind driven effect towards the overall surge height. There is very little work addressing the contribution of oceanic precipitation to surge levels. This is expected, considering that many ocean models do not incorporate the mass and momentum effect of precipitation for it to be studied in detail. Sheng et al. (2012) described a test



1  
2  
3 platform used to review several established storm surge models such as CH3D-SSMS (Sheng  
4 et al., 2010), ADCIRC (Westerink et al., 1992), FVCOM (Chen et al., 2008), CMEPS (Peng  
5 et al., 2004) and SLOSH (Jarvinen et al., 1985). Thirty scenarios were used to test their  
6 sensitivity towards parameters such as bathymetry, storm forcing, wind drag coefficient,  
7 bottom friction, Coriolis and 2D-3D formulation, but precipitation was not investigated.  
8  
9 Ocean models such as ROMS (Shchepetkin et al., 2005) and FVCOM only incorporate the  
10 effect of temperature, salinity and momentum from rain, but not the effect of rain mass.  
11 Storm surge models such as ADCIRC and SLOSH ignore the effects of ocean rain entirely.  
12 Ponte (2006) incorporated realistic freshwater fluxes in a global barotropic model and found  
13 that they can cause annual standard deviations in sea surface level as large as 1 cm, especially  
14 in shallow and semi-enclosed regions. In highly precipitating storm events with extreme rain  
15 rates, the sea surface height increase may be more substantial. For instance, the highest  
16 rainfall recorded by Typhoon Morakot in 2009 was up to  $74 \text{ cm day}^{-1}$  (Wu, 2013). How this  
17 contributes to storm surges and to what extent rain mass can be neglected in this case is not  
18 known.  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36

37 Here, a new rainfall scheme has been introduced into ROMS to account for the effect of rain  
38 mass. This scheme complements the existing rainfall routine in ROMS where only  
39 temperature and salinity changes due to rainfall are accounted for. An idealized set-up of  
40 ROMS will be used to isolate the effects of adding rain mass to the sea surface height. The  
41 Weather Research and Forecasting model (WRF, Skamarock et al., 2008) is used to generate  
42 the rain forcing to evaluate the sea surface height response in ROMS.  
43  
44  
45  
46  
47  
48  
49  
50

51 The transfer of momentum from raindrops to the water surface, or rain stress, was first  
52 documented by Van Dorn (1953), who concluded that the stress contributions from rainfall  
53 can considerably intensify surface stresses. Caldwell et al. (1970) have since built on his  
54  
55  
56  
57  
58  
59  
60

1  
2  
3 work, parameterizing rain stress as a function of rain rate and wind speed. This  
4  
5 parameterization has been widely incorporated in many model studies such as the one-  
6  
7 dimensional mixed layer model developed by Clayson et al. (1999) and the bulk  
8  
9 parameterization outlined by Fairall et al. (1996) relating the near-surface atmospheric and  
10  
11 the oceanographic bulk variables. This air-sea bulk parameterization has also been adopted in  
12  
13 ROMS. While the parameterization of rain stress has been widely incorporated, there has  
14  
15 been no study so far that shows its isolated effect in the models it has been implemented in.  
16  
17 For the case of tropical cyclones, where extreme wind speeds and precipitation rates are  
18  
19 found, the isolated effect of rain stress has not been shown and its contribution to storm  
20  
21 surges is not known. This section shall employ the use of an idealised set-up in ROMS to  
22  
23 show the isolated effects of rain stress in the context of storm surges.  
24  
25  
26  
27

28  
29 Finally, a real case scenario is studied to compare the various processes affecting storm  
30  
31 surges. Tropical Cyclone Monica (Durden, 2010) has been chosen for her high wind speed,  
32  
33 low system pressure and high rainfall, which are the key storm surge effects investigated  
34  
35 here. A validated atmospheric state is generated using WRF and is prescribed onto ROMS to  
36  
37 investigate how various processes affect the storm surge. The significance of rain mass and  
38  
39 rain stress is compared with two other established storm surge processes: the inverse  
40  
41 barometer effect and the effect of surface wind. Each of these physical effects is investigated  
42  
43 based on their isolated contribution to the total SSH and storm surge patterns.  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

## 2 Model set-up and validation

### 2.1 Set-up of idealized case

Rain mass is introduced to ROMS through a new rainfall scheme that incorporates rain rate as point sources. This approach is similar to how river run-off is introduced to the model, but in this case, vertical fluxes are added instead of horizontal fluxes. The rainfall rate from the atmospheric rain forcing is converted into point sources without further prescribing its temperature and salinity. The temperature and salinity of rain is accounted for within the existing bulk flux routine where air-sea momentum and energy exchanges are calculated (Fairall et al., 1996). An idealized set-up of ROMS is used to create a controlled environment, isolating the specific effects investigated here. The ROMS grid is set up with closed boundaries, a horizontal resolution of 15 km, 21 vertical levels and a spatial extent 5°N to 35°N and 120°E to 180°E. The initial ocean state has a uniform temperature of 26°C and salinity of 35 psu. The land mask is specified at the western boundary up to 140°E to simulate the coastline. A third-order upwind scheme is used for horizontal momentum advection, and a Smagorinsky-like viscosity is applied (Griffies et al., 2000). The turbulence closure scheme used to calculate vertical mixing is based on the Generic Length Scale (GLS) parameterization as described by Umlauf et al. (2003).

An idealized tropical cyclone is simulated using WRF and prescribed onto ROMS. The set-up of the initial cyclone profile and atmospheric state follows Wang et al. (2015). For the WRF simulation, one single nest is set up with 31 vertical levels and a horizontal resolution of 15 km in the domain spanning 5°N to 35°N and 120°E to 180°E. The microphysical parameterization used is based on Lin et al. (1983). Cumulus parameterization is not included, which reduces the computational time required for each simulation. Without

1  
2  
3 cumulus parameterization, the amount of precipitation may be underestimated at a 15 km  
4  
5 horizontal resolution. However, the aim of this part of the study is to understand the specific  
6  
7 effect of adding precipitation mass to the ocean model in a simple, idealized set-up, and not  
8  
9 to reproduce the atmospheric domain with high accuracy. The surface layer scheme is based  
10  
11 on Monin-Obukhov, with Zilitinkevich thermal roughness length and standard similarity  
12  
13 functions (Paulson, 1970). Planetary boundary layer processes are parameterized using the  
14  
15 Mellor-Yamada-Janjic scheme (Janjic, 1994). The longwave and shortwave radiation  
16  
17 schemes are implemented from Mlawer et al. (1997) and J. Dudhia (1989) respectively.  
18  
19  
20  
21  
22

## 23 2.2 Rain stress parameterization

24  
25  
26 In ROMS, rain stress is implemented in the bulk fluxes routine from Fairall et al. (1996) and  
27  
28 based on the parameterization by Caldwell and Elliott (1970) as follows:  
29  
30  
31

$$32 \tau_{rx} = 0.85R\rho_r |\vec{v}_{10}| * SIGN(U_{10}) \quad (1)$$

$$33 \tau_{ry} = 0.85R\rho_r |\vec{v}_{10}| * SIGN(V_{10}) \quad (2)$$

34  
35  
36 where R is the rain rate,  $\rho_r$  is the density of rainwater,  $\tau_{rx}$  is the rain stress in the x direction,  
37  
38  $SIGN(U_{10})$  denotes the sign of  $U_{10}$  (U-component wind speed at 10 m),  $\tau_{ry}$  is the rain stress  
39  
40 in the y direction and  $SIGN(V_{10})$  denotes the sign of  $V_{10}$  (V-component wind speed at 10 m).  
41  
42  
43  
44  
45  
46  
47  
48

49  
50 The rain stress parameterization of Caldwell and Elliott (1970) is derived from a one-  
51  
52 dimensional model where  $\vec{v}_{10}$  refers to the wind velocity in the direction of reference. When  
53  
54 considering wind velocity resolved into its orthogonal components, it would be incorrect to  
55  
56 utilize the total velocity in calculating the rain stress. Here, the magnitude of the u-component  
57  
58  
59  
60

winds ( $U$ ) and the v-component winds ( $V$ ) used in the calculation of rain stress will always be  $\sqrt{U^2 + V^2}$ . This results in a systematic overestimation in the magnitude of rain stress by a factor of  $\sqrt{2}$ . The error in the direction of rain stress is more inconsistent. Since  $U$  and  $V$  are always set to be equal in magnitude (equal to the net wind speed), the directions of the rain stress can only be in any of the four directions:  $45^\circ$ ,  $135^\circ$ ,  $225^\circ$  and  $315^\circ$ , unless there is zero rainfall or zero wind speed. Furthermore, the existing ROMS rain stress code always assigns a positive sign to zero values. For example, if  $U$  is zero, it will still be assigned the magnitude of  $V$  (now equal to the magnitude of the net wind speed), with the direction of positive  $U$ . The error in the direction of the rain stress term can range from  $0^\circ$  (when  $U=V$ ) to  $45^\circ$  (when either  $U=0$  or  $V=0$ ). Overall, this results in a systematic overestimation of rain stress and an incorrect direction calculated for the overall wind stress.

The corrected implementation of the rain stress parameterization is as follows:

$$\tau_{rx} = 0.85R\rho_r U_{10} \quad (3)$$

$$\tau_{ry} = 0.85R\rho_r V_{10} \quad (4)$$

An idealized set-up of ROMS is used to verify the corrected implementation of rain stress. Here, the u-component wind is set to a constant  $-1 \text{ m s}^{-1}$  and the v-component wind is set to zero. One centimetre of rain is introduced to a single cell (15 km by 15 km) in the centre of the domain ( $20^\circ\text{N } 150^\circ\text{E}$ ) at the first hour. Two simulation runs are conducted to compare the existing implementation of rain stress (using equations (1) and (2)) with the corrected implementation of rain stress (using equations (3) and (4)). Both simulation runs are subtracted from a control run without any rain stress to eliminate any noise such as the SSH response to the background u-component wind.

### 2.3 Set-up of real case

The ROMS set-up uses a single grid at a horizontal resolution of 6 km, with the open boundaries using the Flather (1976) conditions for the barotropic velocities and the Chapman (1985) conditions for sea surface elevation. A minimum depth of 10 m and maximum depth of 5500 m has been specified. Twenty layers are chosen in the vertically stretched terrain-following sigma-coordinate, with higher resolution placed at the surface and lower resolution placed at the bottom. Geometry of the bathymetry is obtained from Global Topography v14.1 (Smith, 1997). The tidal forcing is applied at the open boundaries and imposed on the elevation and the barotropic velocities. It is derived from 11 tidal harmonics that are extracted from the 1/12° resolution Pacific Ocean Atlas solution provided by the Oregon State University (OSU) Tidal Data Inversion (Egbert et al., 1994). A nudging relaxation zone with eight grids from the boundary is set up to relax the baroclinic structure to the forcing fields at the boundary (Marchesiello et al., 2001). The 1/12° resolution HYbrid Coordinate Ocean Model (HYCOM) (Bleck et al., 1981) global daily analysis data is used as the initial, boundary and nudging conditions for the model.

The pressure-driven effect in ROMS is obtained by applying the atmospheric pressure field from WRF and the sea level is adjusted based on the inverse barometer approximation (Gill and Niiler, 1973). The inverse barometer effect is generally a good approximation for sea level responses to local atmospheric pressure. For non-isostatic pressure effects, Hirose et al. (2001) found that deviations from the inverse barometer effect contributed only up to 2 cm in large-scale basins such as the Pacific and Atlantic Ocean. For smaller-scale basins such as Cape of Carpentaria, the non-isostatic deviation was under 0.5 cm and is much smaller compared to the precipitation effects in this case (Figure 10b).

1  
2  
3 WRF is used to simulate a high resolution atmospheric state of tropical cyclone Monica to be  
4 prescribed onto the ocean model for the case study. Two nests are employed with 28 vertical  
5 levels. The final grid has a horizontal resolution of 6 km. NCEP Climate Forecast System  
6 Reanalysis (CFRS) (Saha et al., 2010) has been used for initial and lateral boundary  
7 conditions. The boundary forcing is prescribed every 6 hours and the model interior is  
8 allowed to evolve freely. The ETA scheme (Rogers et al., 2001) is chosen for microphysical  
9 parameterization. The Kain-Fritsch cumulus parameterization scheme (Kain, 2004) has been  
10 used for both nests. The surface layer scheme is based on Monin-Obukhov with Carlson-  
11 Boland viscous sub-layer and standard similarity functions. The selected planetary boundary  
12 layer is from Yonsei University (Hong et al., 2006). The Mesoscale Model 5 (MM5) 5-layer  
13 thermal soil temperature model (Dudhia, 1996) is used to calculate the heat and moisture  
14 fluxes over the land. The longwave and shortwave radiation schemes are implemented from  
15 Mlawer et al. (1997) and Chou et al. (1994) respectively.

16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33 The isolated effects of including precipitation mass source, rain stress, inverse barometer  
34 effect and surface wind forcing are individually investigated by comparing with four separate  
35 configurations that are simulated with precipitation mass source, without rain stress, without  
36 the inverse barometer correction and without surface wind forcing respectively. The five  
37 experiments conducted are tabulated in Table 2.1. Spatial plots of SSH are differenced from  
38 the control simulation (Case A) to illustrate the isolated effects of each process.  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

<i>Cases</i>	<i>WRF wind conditions applied to ROMS</i>	<i>Inverse barometer effect in ROMS</i>	<i>Rain stress effect in ROMS</i>	<i>New rain mass scheme</i>
<i>Case A</i>	Wind forcing applied	Turned ON	Turned ON	NOT used
<i>Case B</i>	NO wind forcing applied	Turned ON	Turned ON	NOT used
<i>Case C</i>	Wind forcing applied	Turned OFF	Turned ON	NOT used
<i>Case D</i>	Wind forcing applied	Turned ON	Turned OFF	NOT used
<i>Case E</i>	Wind forcing applied	Turned ON	Turned ON	In use

Table 2.1: Cases set up with different ROMS model options and wind conditions.

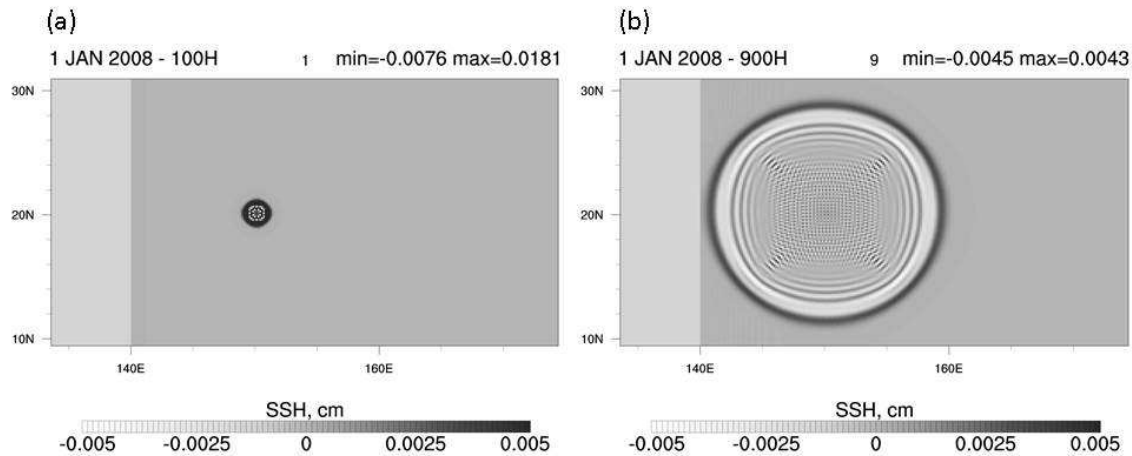
### 3 Results

#### 3.1 Idealized case: response to rain mass

An idealized set-up of ROMS is used to study the effect of rainfall mass on sea surface elevation. The effect of rainfall is isolated by setting the wind field to zero and all other atmospheric forcings are set to a constant nominal value. Rain is introduced to a single grid cell (15 km by 15 km) at 20°N 150°E for 1 hour at the rate of 1 cm/h, to understand the effect of a single-point perturbation on the free surface. Figure 1 shows the ocean surface response after 2 hours and 9 hours. The illustrations are zoomed into the area 9°N to 31°N and 134°E to 176°E. A surface gravity wave is generated, propagating radially outwards from the point of perturbation. At time = 1 h, the maximum SSH is 0.0181 cm and the minimum is -0.0076 cm. At time = 9 h, the maximum SSH is 0.0043 cm, the minimum is -0.0045 cm and the



1  
2  
3 wave propagates to about  $10^\circ$  from the initial point of perturbation. Over 9 hours, the  
4  
5 amplitude of the gravity wave drops by about 76%.  
6  
7  
8  
9



26  
27 Figure 1: SSH response to rain mass addition to a single surface grid cell after (a) 1 hour and  
28  
29 (b) 9 hours. Land is west of  $140^\circ\text{E}$ , ocean is east of  $140^\circ\text{E}$ . Wind forcing is set to zero.  
30  
31

32  
33 Next, the SSH response to rainfall at a longer time scale is studied by forcing ROMS with  
34 rainfall from an idealized tropical cyclone simulated using WRF over 8 days. The maximum  
35 rain rate is not at the centre of the cyclone, but along the rain bands. As the cyclone treks  
36 from the east to the west of the domain, it gradually grows in size and intensity. The rain  
37 bands of the tropical cyclone also increase in size with time. Figure 2 shows the cumulative  
38 rainfall throughout the event, with a maximum of 319 mm at any one grid. To avoid any  
39 interference from the winds, the wind field is set to zero so the sea surface is unaffected by  
40 wind stress and responds only to the cyclone rainfall.  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

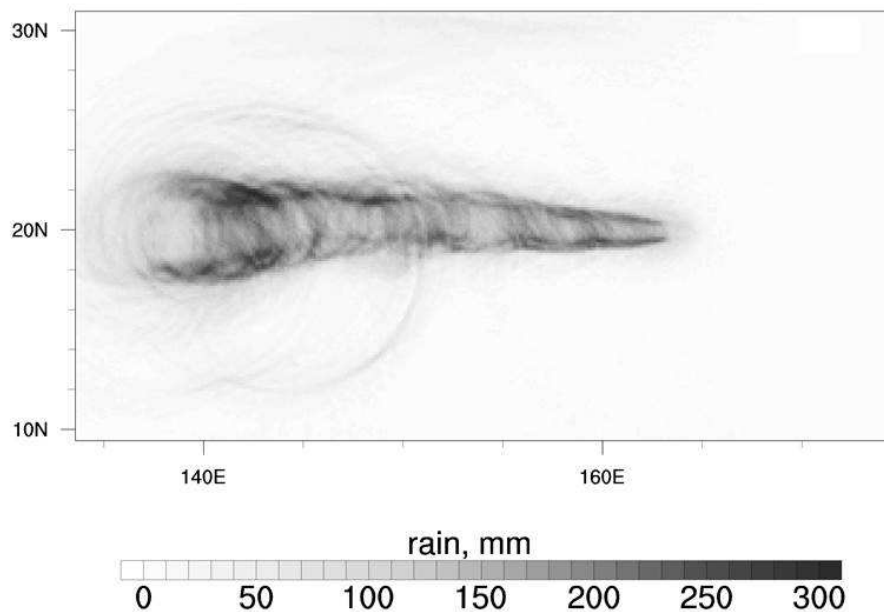


Figure 2: Cumulative rainfall of the idealized tropical cyclone over 8 days.

The SSH response to rainfall from the idealized tropical cyclone is simulated for two ocean depths at 100 m and 1000 m deep. Figure 3 shows the SSH at the 170<sup>th</sup> hour of the simulation when the tropical cyclone makes landfall at the coastline along 140°E. For both depths, the SSH increases along the rainfall path and gradually decreases towards the domain boundaries. For the case with a depth of 100 m, the maximum SSH is 5.8 cm and the minimum SSH is 1.3 cm at the 170<sup>th</sup> hour. A clockwise gyre is formed, with the sea surface currents at a magnitude of  $10^{-2} \text{ m s}^{-1}$ . For the case with a depth of 1000 m, the maximum SSH is 3.2 cm while the minimum SSH is 2.3 cm. A weaker clockwise gyre is formed with the sea surface currents at a magnitude of  $10^{-3} \text{ m s}^{-1}$ .

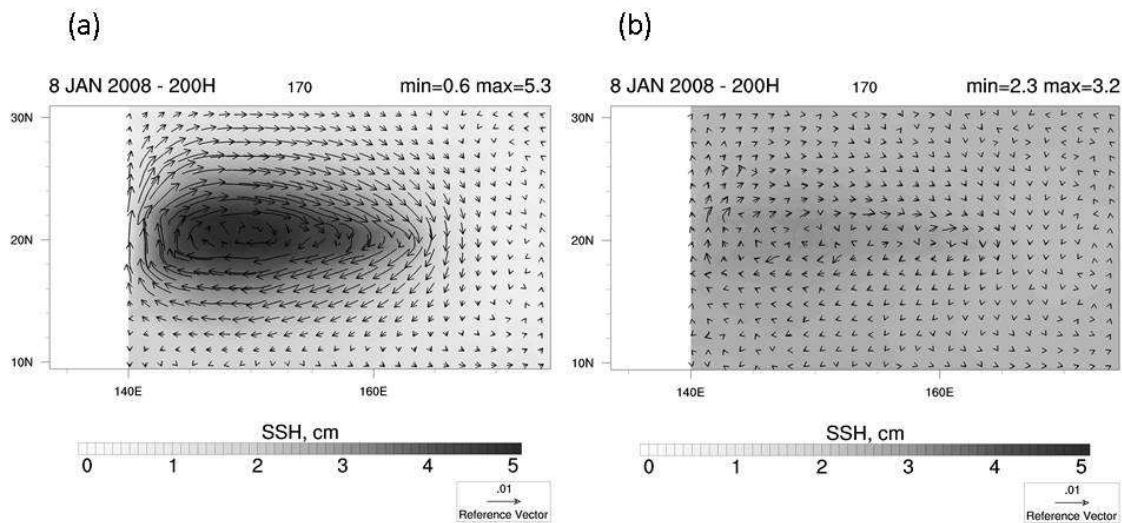


Figure 3: SSH and sea surface current response to the rain mass added to the ocean with (a) 100 m depth and (b) 1000 m depth at the 170<sup>th</sup> hour.

### 3.2 Idealized case: response to rain stress

Figure 4 shows the rain stress perturbation setting off a surface gravity wave that propagates outwards from the point of perturbation towards the domain boundaries. Two implementations are shown for the existing and the corrected model. In the existing implementation of rain stress (Figure 4a), the resulting surface gravity wave is at an angle of 45° to the wind direction. The front of the wave is led by a crest in the north-westward direction with an amplitude of up to 0.01 cm. In the south-eastward direction, the wave is led by a trough of up to 0.01 cm in amplitude. Figure 4b shows the SSH response to the corrected implementation of rain stress in ROMS. The surface gravity wave generated from the rain stress perturbation is now symmetrical to the wind direction. In the westward direction, the front of the wave is led by a crest of up to 0.005 cm in amplitude. In the eastward direction, the wave is led by a trough of up to 0.005 cm in amplitude.

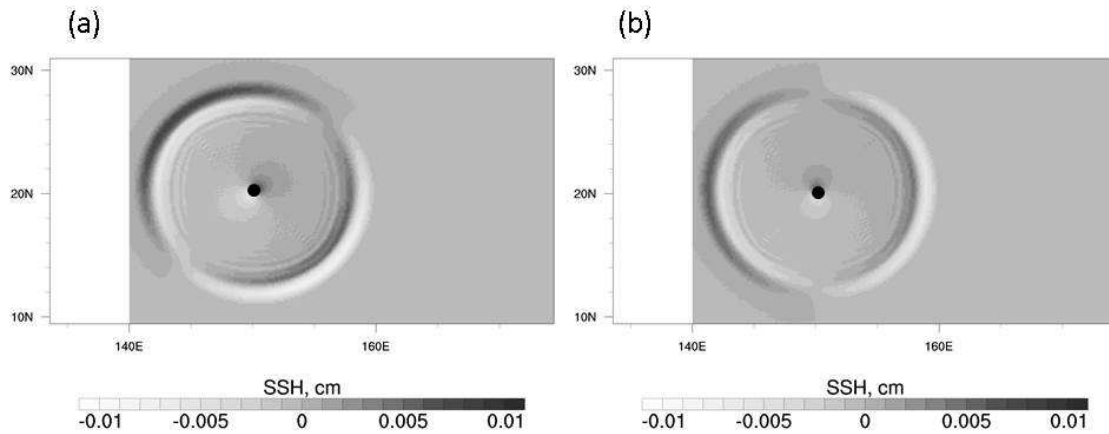


Figure 4: Sea surface height at 10 hours after introducing a rain stress perturbation for (a) the existing implementation of rain stress in ROMS and (b) the corrected implementation of rain stress in ROMS. The black dot in the centre of the domain ( $20^{\circ}\text{N } 150^{\circ}\text{E}$ ) denotes the initial position of perturbation. Wind forcing is set to zero.

Next, the contribution of rain stress to storm surges is investigated by forcing ROMS with an idealized tropical cyclone simulated using WRF. Here, the wind and rainfall fields are retained, but the other forcing fields are set to a nominal and constant value in order not to introduce interfering signals into the system. To show the isolated effect of rain stress, the response of SSH and sea surface currents from the rain stress simulation was subtracted from a control simulation (Figure 5). Here, an anticlockwise circulation centering  $20^{\circ}\text{N } 141^{\circ}\text{E}$  is generated with a decrease in SSH at the circulation centre. The positive surge is up to +0.44 cm and the negative surge is up to -0.63 cm.

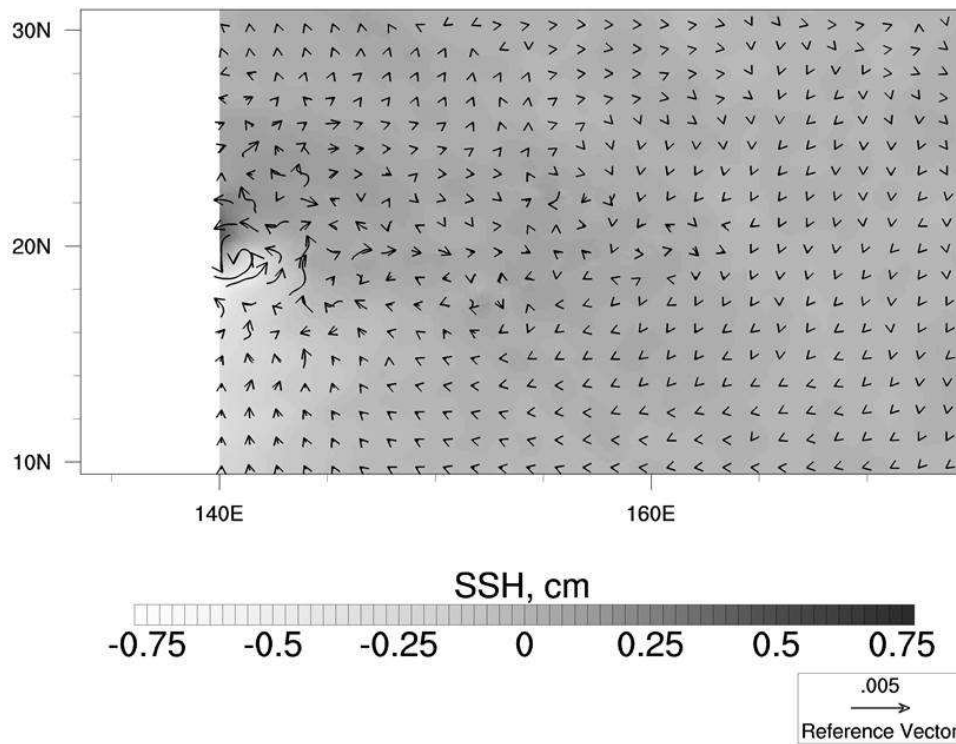


Figure 5: Response of SSH and sea surface currents to rain stress at  $t = 163$  h.

### 3.3 Case study: validation of Cyclone Monica

The atmospheric state is simulated for the period between 20 April 2006 and 24 April 2006. Minimum sea level pressure from WRF is used to generate the cyclone path for comparison as shown in Figure 6. This cyclone path is validated using the storm track provided by the Australian Bureau of Meteorology (BOM, <http://www.bom.gov.au/cyclone>). On 20 April 2006, both the modelled cyclone and the observed cyclone are located at the western coast of Cape York Peninsula, but the modelled cyclone is situated more northwards compared to the observed cyclone. The modelled path is largely similar to the observed path, but has a

northward offset of up to  $1^\circ$ . The modelled cyclone also travels faster than the observed cyclone by up to a day.

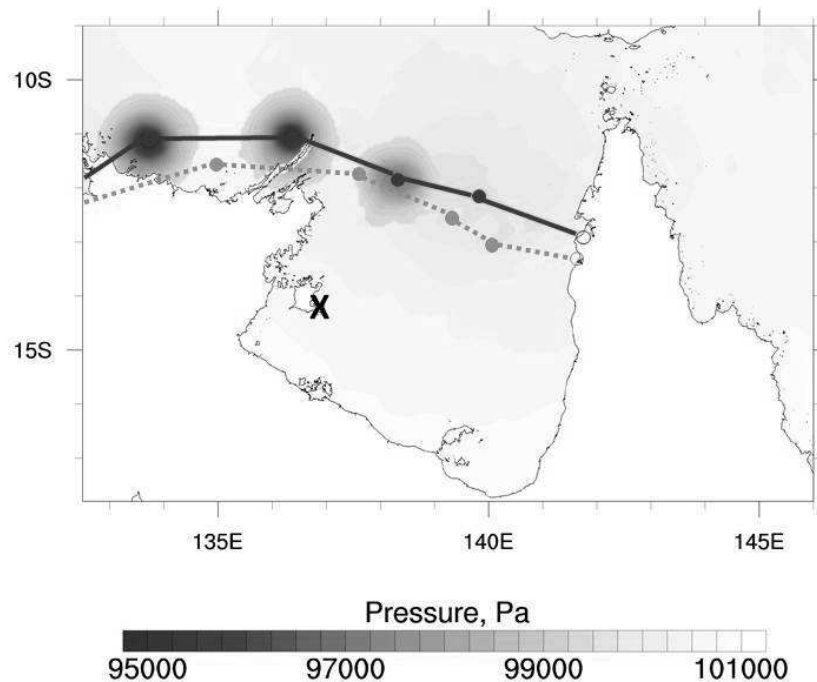


Figure 6: Cyclone path of WRF (solid line) and observations from BOM cyclone report (dashed line). Filled circles are positions every 24 hours from 0000 UTC 20 April 2006 (open circle). Location of the Groote Elyandt weather station is indicated with an 'X'.

Based on the observed weather conditions provided by BOM, the maximum wind speed observed at Cape Wessel ( $11^\circ\text{S}$ ,  $136.8^\circ\text{E}$ ) and Maningrida ( $12^\circ\text{S}$ ,  $134.2^\circ\text{E}$ ) were  $130\text{ km h}^{-1}$  and  $148\text{ km h}^{-1}$  respectively. In WRF, the maximum modelled wind speed at Cape Wessel and Maningrida were  $141\text{ km h}^{-1}$  and  $146\text{ km h}^{-1}$  respectively, which compared well with observations (Figure 7).

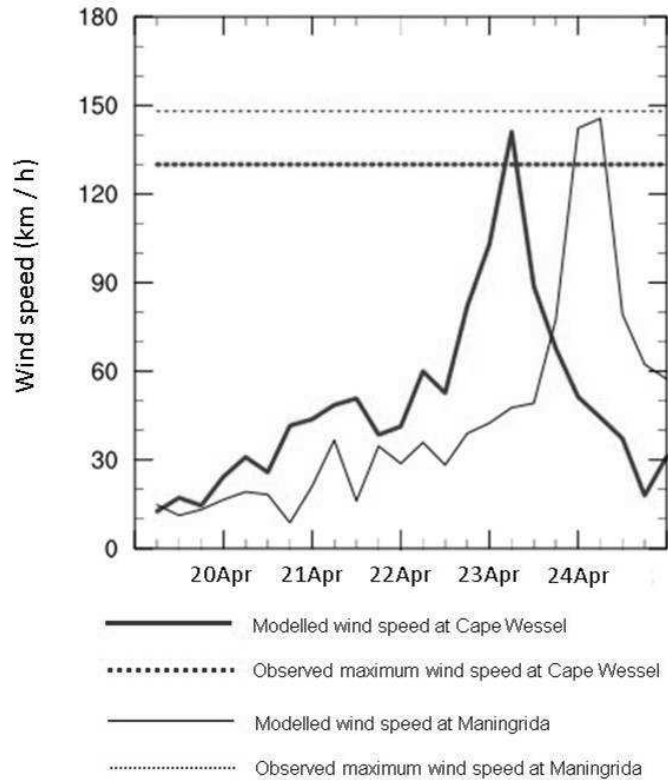
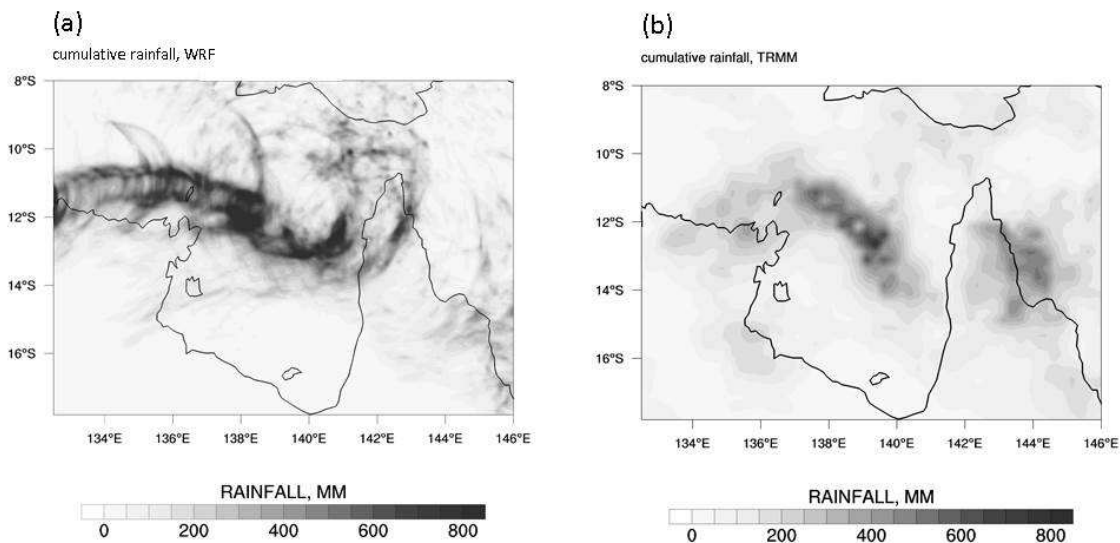


Figure 7: Comparison of the modelled wind speeds at Cape Wessel (11°S, 136.8°E) and Maningrida (12°S, 134.2°E) with the maximum wind speeds observed.

The simulated precipitation is compared with the multi-satellite rainfall estimates from the Tropical Rainfall Measuring Mission (TRMM) 3B42 version 6 (Huffman et al., 2007) (Figure 8). The modelled rainfall path is largely in line with the observed rainfall path but the rainfall is more intense along the cyclone track and less intense away from the track. The observed cumulative rainfall (TRMM) is more homogeneous, having more regions with 100 – 300 mm of rainfall. Nevertheless, the domain cumulative rainfall in WRF is 6% more than that observed by TRMM, although this percentage is expected to be higher since the WRF rainfall is summed with the land mask applied, while TRMM rainfall is summed over every cell in the domain. The total domain cumulative rainfall simulated using WRF is  $4.7 \times 10^6$  mm and



1  
2  
3 the observed cumulative rainfall in TRMM is  $4.4 \times 10^6$  mm over 8 days. However, this is not  
4  
5 a large concern in this case study since the main objectives are to simulate and compare  
6  
7 various processes in a realistic environment and not to closely replicate the exact storm event.  
8  
9 The precipitation rates modelled here are not uncommon among tropical cyclones in  
10  
11 Australia and similar storm systems in other parts of the world. From the BOM reports on  
12  
13 historical tropical cyclones, the highest observed precipitation rates of Monica was 340 mm  
14  
15  $\text{day}^{-1}$ , Laurence (December 2009) at 402  $\text{mm day}^{-1}$ , Paul (March 2010) at 443  $\text{mm day}^{-1}$  and  
16  
17 Yasi (January 2011) at 471  $\text{mm day}^{-1}$ . Outside of Australia, the estimated maximum rainfall  
18  
19 of Hurricane Katrina (August 2005) was 300  $\text{mm day}^{-1}$  (Dodla et al., 2011) and Typhoon  
20  
21 Morakot (August 2009) was 741  $\text{mm day}^{-1}$  (Wu, 2013). Hence, the rainfall produced by  
22  
23 Monica is not an extreme case compared to other tropical cyclones and the overestimation in  
24  
25 Monica is not an extreme case compared to other tropical cyclones and the overestimation in  
26  
27 the rainfall simulated here does not lead to an unrealistic scenario.  
28  
29



50 Figure 8: 8-Day cumulative rainfall output from 19 April 2006 for (a) WRF and (b) TRMM.

51  
52  
53 The configuration of ROMS used for validating the ocean state incorporated the effects of  
54  
55 rain mass, rain stress, the inverse barometer effect and all atmospheric forcings including  
56  
57 wind. The storm surge (Figure 9) is validated against the tidal gauge measurements taken  
58  
59  
60



1  
2  
3 from Milner Bay in Groote Elyandt. The tidal gauge data is obtained from the Australian  
4 Baseline Sea Level Monitoring Project (ABSLMP) (Watson, 2011). To compare the storm  
5 surge, the tidal components are detided using the tidal package by Codiga (2011). The tidal  
6 range is very large and there remains a residual tidal-like signal after de-tiding eleven tidal  
7 components that were introduced into ROMS in the tidal forcing set-up. At this location, the  
8 surge simulated in ROMS (23 April 2006) is earlier than the observed data (24 April 2006)  
9 by a day. However, the magnitude of the surge is very well-captured. The difference in  
10 modelled and observed time of peak surge is also consistent with the different translation  
11 speed and track of the WRF cyclone. The modelled cyclone is located north of the Groote  
12 Elyandt station around 1200 UTC 22 April 2006, while the observed cyclone reaches north of  
13 the Groote Elyandt station around 0000 UTC 24 April 2006. However, the modelled cyclone  
14 path is situated more north compared to the observed cyclone path and a longer time is  
15 expected for the storm surge to be picked up at Groote Elyandt. Hence, an overall difference  
16 of a day can be expected between the modelled and observed storm surge.  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

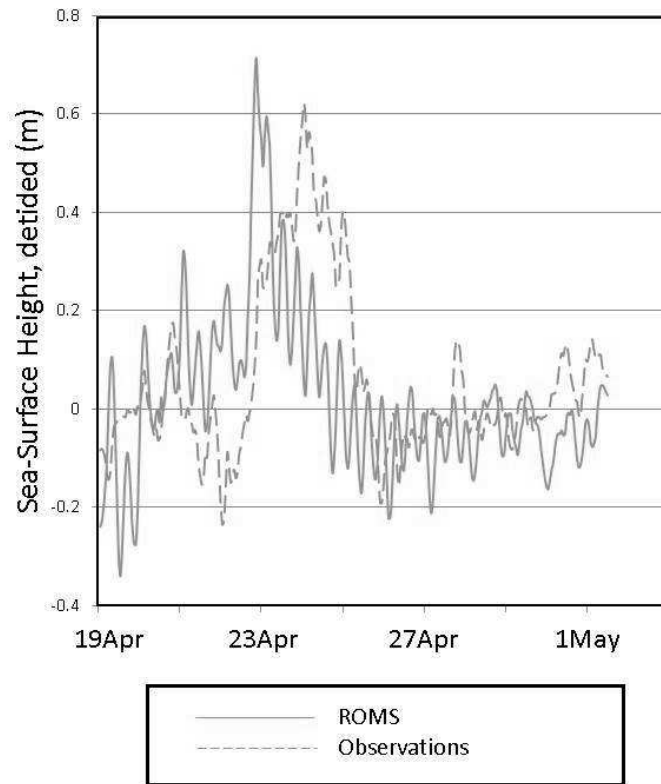


Figure 9: Comparison between the storm surge generated in ROMS with tidal gauge measurements (detided) at Groote Elyandt weather station obtained from ABSLMP, with the location indicated with an 'X' on Figure 6.

### 3.4 Case study: contribution of rain to SSH

We examine the individual contributions to the SSH response to Cyclone Monica. The SSH response to the effect of wind (Figure 10a) is the largest compared to the other effects. On 21 April 2006, the negative surge is up to -58 cm at the south-eastern boundary of the Gulf of Carpentaria and the positive surge is up to 61 cm on the western boundary. On 22 April 2006, the negative surge at the south-eastern boundary of the gulf and the positive surge at the western boundary still persist. There is a region of negative SSH located at the centre of the

1  
2  
3 gulf. The minimum SSH is -57 cm and maximum SSH is 60 cm. On 23 April 2006, the  
4  
5 cyclone reaches the Wessel Islands and the SSH rises up to 171 cm in Buckingham Bay  
6  
7 (12°S 136°E). The positive surge at the western boundary of the gulf is still present at this  
8  
9 time, but the negative surge at the south-eastern boundary has reverted to its nominal levels.  
10  
11 At the northern coast of the Arnhem Land, the negative surge is up to -111 cm.  
12  
13

14  
15 Figure 10b shows the SSH change due to the inverse barometer effect. The SSH increase is  
16  
17 localized to the position of the cyclone and increases as the cyclone intensifies through the  
18  
19 event. The distribution of the SSH scalar field is similar to atmospheric pressure, with the  
20  
21 highest SSH at the centre of the cyclone, gradually decreasing with distance from the cyclone  
22  
23 centre. The SSH increases from 13 cm on 21 April 2006 to 61.5 cm on 23 April 2006 as the  
24  
25 cyclone intensifies.  
26  
27

28  
29 The additional SSH response due to the effect of rain stress is shown in Figure 10c. On 21  
30  
31 April 2006, the additional negative surge is up to -2.4 cm on west coast of Cape York  
32  
33 Peninsula. There is an additional positive SSH of up to 0.9 cm near the position of the  
34  
35 cyclone centre. On 22 April 2006, when Monica is at the centre of the Gulf of Carpentaria,  
36  
37 there is an additional positive SSH of up to 0.7 cm ahead of the cyclone and a negative SSH  
38  
39 of up to -1.1 cm behind the cyclone. On 23 April 2006, there is an additional positive surge of  
40  
41 up to 7.5 cm to the east of the Wessel islands and an additional negative surge of up to -2.9  
42  
43 cm to the west of the islands.  
44  
45  
46

47  
48 Figure 10d shows the additional SSH response due to rain mass, increasing as the cyclone  
49  
50 tracks across the Gulf of Carpentaria and decreasing when the cyclone passes. On 21 April  
51  
52 2006, Monica is situated at the western coast of Cape York Peninsula and the SSH rises up to  
53  
54 3.7 cm at the coast. On 22 April 2006, Monica is centred in the gulf and the SSH increases up  
55  
56 to 4.6 cm, spreading out to the rest of the gulf. The sea level at the southern boundary of the  
57  
58  
59  
60

1  
2  
3  
4  
5  
6  
7  
8  
9  
10  
11  
12  
13  
14  
15  
16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

gulf increases up to 3.5 cm. On 23 April 2006, the cyclone reaches the Wessel Islands and the SSH rises up to an additional 6.4 cm. The sea level change at the southern boundary of the gulf decreases to 2 cm.

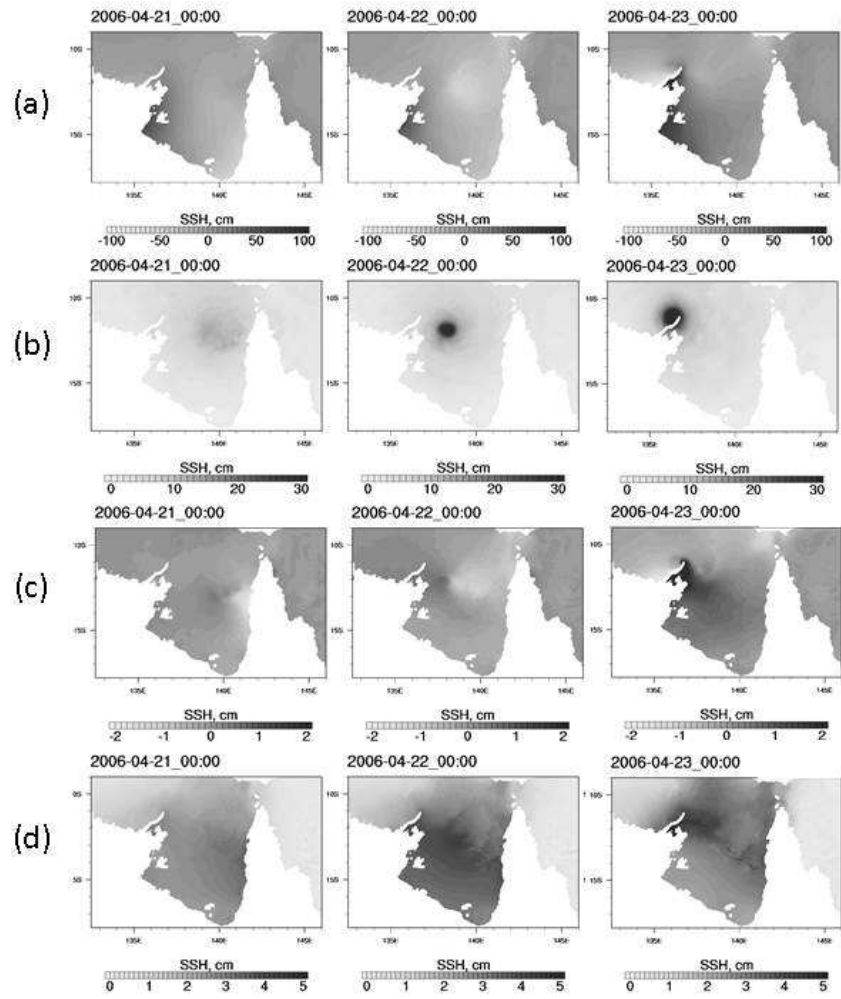


Figure 10: Additional SSH response to the effect of (a) wind, (b) atmospheric pressure, (c) rain stress and (d) rain mass from 21 April 2006 to 23 April 2006.

#### 4 Discussion

The rain mass perturbation on the water surface sets off a surface gravity wave propagating outwards towards the domain boundaries over a time scale of several hours (Figure 1).

Characteristics of surface waves are inherently non-linear. But considering shallow water approximations (Stewart, 2009), where the water depth is much smaller than the wavelength, the phase speed ( $c$ ) can be expressed as:

$$c = \sqrt{gd} \quad (5)$$

where  $g$  is the gravitational acceleration and  $d$  is the depth of water.

From the initial point of perturbation, the surface wave propagated from 20°N 150°E to 20°N 160°E in 9 hours translating at a speed of 30.9 m s<sup>-1</sup>. The shallow water approximation is valid for water depths much lesser than the wavelength. Here, the water depth is 100 m while the wavelength is about 100 km. The water depth in this case is only 0.1% that of the wavelength, making the shallow water approximation a valid assumption. From equation (5), the surface gravity wave in this case should advance at a phase speed of 31.3 m s<sup>-1</sup>, which is consistent with the ROMS output.

Next, the SSH response at a longer time scale is examined using an idealized tropical cyclone to force the ocean model over 8 days. In Figure 3, the added mass spreads outwards from the mass sources towards the domain boundaries via gravity waves. In the deeper ocean case, the maximum SSH is smaller (up to 3.2 cm) compared to the shallow ocean case (up to 5.3 cm). From equation (5), gravity waves in the deeper case travel 3 times faster than in the shallow water case, which accounts for the faster decrease in the maximum SSH for the deeper ocean

1  
2  
3 case. With the rain mass spreading out faster in the deeper ocean, the minimum SSH reaches  
4  
5 a higher level for the deep ocean case (2.3 cm) compared to the shallow ocean case (1.3 cm).  
6  
7 At longer time scales (after 7 days), a pressure gradient is generated from the sea surface  
8  
9 height increase. This creates a geostrophic adjustment with an opposing Coriolis force.  
10  
11 Geostrophic currents of the magnitude  $0.01 \text{ m s}^{-1}$  are generated, rotating in a clockwise  
12  
13 direction.  
14  
15

16  
17  
18 With the implementation of rain stress now corrected in ROMS, the SSH response to rain  
19  
20 stress is shown by forcing ROMS with an idealized tropical cyclone (Figure 5). Rain stress is  
21  
22 parameterized as a function of rain rate and wind speed and follows the direction of the  
23  
24 winds. Hence, rain stress is only exerted in areas where non-zero values of rain rate and wind  
25  
26 speed coincide. In the case of a tropical cyclone, rainfall is highest along the rain bands  
27  
28 towards the centre of the cyclone. Rain stress contributes to storm surges in a manner similar  
29  
30 to wind stress. The anti-clockwise rotating rain stress creates an anti-clockwise gyre at the sea  
31  
32 surface. Coriolis force deflects the currents to the right and generates an outward diverging  
33  
34 flow that decreases the SSH at the centre of the gyre.  
35  
36  
37  
38

39  
40 Finally, a real case study is investigated to compare the SSH response to different storm surge  
41  
42 processes. The effect of wind (Figure 10a) is the largest compared to the effect of  
43  
44 precipitation and the inverse barometer effect. It drives positive and negative SSH anomalies  
45  
46 based on the wind direction and the geometry of the coastline. On 21 April 2006, the strong  
47  
48 easterly wind to the south of the domain generates a positive surge at the western boundary of  
49  
50 the Gulf of Carpentaria due to the onshore wind. A negative surge is generated at the south-  
51  
52 eastern boundary of the gulf due to the offshore wind. On 22 April 2006, the easterly wind to  
53  
54 the south of the domain continues to generate a positive and negative surge at the western and  
55  
56 south-eastern boundaries of the gulf respectively. There is a SSH depression in the middle of  
57  
58  
59  
60

1  
2  
3 the gulf due to Ekman pumping. This SSH depression is present in all the 6-hourly output  
4  
5 from 1800 UTC 21 April 2006 to 1800 UTC 22 April 2006 when Monica is trekking across  
6  
7 the Gulf of Carpentaria. The SSH depression is absent on 21 April 2006 and 23 April 2006  
8  
9 when the cyclone is near Cape York Peninsula and the Wessel Islands respectively. On 23  
10  
11 April 2006, the cyclone is relatively far from the eastern boundary of the gulf and the  
12  
13 offshore winds here weakens from the range of  $15 \text{ m s}^{-1}$  (21 April 2006) to the range of  $10 \text{ m}$   
14  
15  $\text{s}^{-1}$ . The negative surge at the eastern boundary of the gulf weakens correspondingly. On the  
16  
17 other hand, the positive surge at the western boundary of the gulf strengthens with the strong  
18  
19 onshore winds. At this time, the strong winds blowing into Buckingham Bay create a  
20  
21 funnelling effect, channelling water into the bay and generating a high positive surge of up to  
22  
23 171 cm.  
24  
25  
26

27  
28 The SSH response to the inverse barometer effect is spatially constrained by the location of  
29  
30 the cyclone's low pressure centre and varies in magnitude based on the storm intensity  
31  
32 (Figure 10b). SSH is expected to increase by 1 cm for every 1 mbar drop in atmospheric  
33  
34 pressure (Roden & Rossby, 1999). The inverse barometer effect contributes substantially to  
35  
36 the storm surge when the low pressure centre of the cyclone coincides with the coastline.  
37  
38 When the cyclone is away from the coastline and over the open sea, the inverse barometer  
39  
40 effect raises the SSH but does not contribute significantly to the storm surge. On 21 April  
41  
42 2006, Monica crosses the west coast of Cape York Peninsula but due to its low intensity at  
43  
44 this time, the contribution to the storm surge is only up to 13 cm. On 22 April 2006, the storm  
45  
46 intensifies as it crosses the Gulf of Carpentaria but due to its distance from the coast, the  
47  
48 contribution to the storm surge is only around 10 cm, despite raising the SSH up to 35.2 cm  
49  
50 in the centre of the gulf. On 23 April 2006, Monica intensifies further, crossing the Wessel  
51  
52 Islands and increases the storm surge by up to 61.5 cm.  
53  
54  
55  
56  
57  
58  
59  
60

1  
2  
3 A region is subjected to rain stress when rainfall and winds coincide. On 21 April 2006, the  
4 location of intense rain is situated just off the west coast of Cape York Peninsula (Figure  
5 10c). The cyclone is traversing westward with an overall westward wind direction. Rain  
6 stress perturbation on the sea surface creates a positive front ahead of the perturbation and a  
7 negative front behind it as shown in Figure 4. The offshore winds and high rainfall at the  
8 eastern boundary of the gulf generates a SSH response to rain stress, enhancing the negative  
9 surge here. While the wind stress generated positive and negative surges at the western and  
10 south-eastern boundaries of the Gulf of Carpentaria, these surges are absent in the rain stress  
11 effect shown here due to the lower rainfall at these locations. On 22 April 2006, the cyclone  
12 is in the middle of the gulf trekking westwards. This creates a similar effect with the positive  
13 SSH ahead of the cyclone and a negative SSH behind it. The positive and negative surges at  
14 the western and south-eastern boundaries of the gulf are similarly absent here, due to the light  
15 rainfall in these areas. On 23 April 2006 when the cyclone crosses the Wessel Islands, the  
16 effect of rain stress is very similar to that of wind stress, with the intense rain coinciding with  
17 the strong winds here. The additional positive surge in Buckingham Bay is generated due to  
18 rain stress channelling water into the bay, which is similar to the effect of wind stress. At the  
19 northern coast of the Arnhem Land just west of the Wessel Islands, the offshore winds  
20 coincide with the high rainfall from the storm to generate an offshore rain stress and a  
21 negative surge of up to -2.9 cm.  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44

45  
46 Figure 10d shows the SSH response to precipitation mass addition. From 21 April 2006 to 22  
47 April 2006, the areas of SSH change do not confine solely to the precipitation footprint (as  
48 traced out in Figure 8a), but spreads out, advecting to the south of the Gulf of Carpentaria. On  
49 23 April 2006, the cyclone crosses the Wessel Islands and the rain mass is trapped in  
50 Buckingham Bay, generating an additional local SSH increase of 6.4 cm. The effect of  
51 precipitation is essentially an in-situ addition of mass, and inlets and semi-enclosed basins  
52  
53  
54  
55  
56  
57  
58  
59  
60



1  
2  
3 will serve to enhance this effect by trapping the rain. Buckingham Bay is fairly well-resolved  
4  
5 in this 6 km resolution model, spanning 13 x 14 grid cells, so the size of the effect is unlikely  
6  
7 to be just numerical. The geometry of the gulf itself is an important contributing factor  
8  
9 towards the effect of precipitation seen here. From Figure 10d, it can be seen that the SSH  
10  
11 response is generally higher to the southern half of the gulf than the north. If the gulf was  
12  
13 entirely enclosed, the additional mass will be expected to distribute evenly by the clockwise  
14  
15 gyre in both the northern and southern regions. The constrained southern boundary of the gulf  
16  
17 enhances the build-up of SSH within the gulf, since the north-western section of the gulf  
18  
19 opens up to the Arafura Sea. Typical remote enhancements are of the order of 2 cm.  
20  
21  
22

23  
24 Evaporation, often considered in tandem with precipitation, can have notable effects on sea  
25  
26 surface elevation. Roden et al. (1999) have reported oceanic evaporation to be between 0.08  
27  
28  $\text{cm day}^{-1}$  to  $0.55 \text{ cm day}^{-1}$ , depending on wind conditions. Hurricane evaporation rates have  
29  
30 been found to be in the order of  $1 - 2 \text{ cm day}^{-1}$  (Riehl and Malkus (1958), Machta (1968),  
31  
32 etc). At this order of magnitude, the evaporation rate is much smaller compared to the effect  
33  
34 of rainfall in this case, with local effects of up to 14 cm (mass and momentum transfer).  
35  
36 Nevertheless, it will be interesting to include the effect of evaporation rate on sea surface  
37  
38 height in future work for a comprehensive comparison.  
39  
40  
41

42  
43 The Gulf of Carpentaria is not an isolated case of a semi-enclosed basin that is prone to  
44  
45 severe storms and many other cases have been studied. Zhong et al. (2010) used ROMS to  
46  
47 study the storm surge predictions during the passage of Hurricane Isabel in Chesapeake Bay  
48  
49 and found that the hurricane translation speed and the resolution of the horizontal wind field  
50  
51 are important factors affecting storm surges. Rego et al. (2010) used FVCOM to simulate  
52  
53 Hurricane Ike at Galveston Bay and found that the Bolivar Peninsula provided a significant  
54  
55 surge barrier in protecting the bay. Westerink et al. (2008) developed a model of south  
56  
57  
58  
59  
60

1  
2  
3 Louisiana (which included many inlets, bays and channels) using ADCIRC to simulate the  
4  
5 effect of storm surges during the events Hurricane Betsy and Andrew. Weisberg et al. (2006)  
6  
7 used FVCOM to simulate the storm surge produced from several idealized hurricanes  
8  
9 crossing Tampa Bay and found that the hurricanes with slow transitional speeds produced  
10  
11 larger surges within the bay. However, none of these studies or models include the effect of  
12  
13 rain mass. The semi-enclosed bay comprising Buckingham Bay, extending to the Wessel  
14  
15 Islands and Gove Peninsula (12°2'S 136°5'E) consists of about 250 cells, which is equivalent  
16  
17 to an area of 9000 km<sup>2</sup>. The areas investigated in literature, such as Tampa Bay (5700 km<sup>2</sup>),  
18  
19 Galveston Bay (1600 km<sup>2</sup>) and Chesapeake Bay (11600 km<sup>2</sup>) are all of similar sizes or  
20  
21 smaller. Not including the effect of rain mass in the models used in these studies will imply a  
22  
23 systematic error in the storm surge levels modelled. To further compare how each of the  
24  
25 physical effects contribute to coastal impacts in a worst case scenario, the SSH at the  
26  
27 coastline starting from the Wessel Islands and ending at the northern tip of Cape York  
28  
29 Peninsula is investigated. The maximum storm surge (for the duration of the storm) due to  
30  
31 each physical effect is shown against the coastal locations (Figure 11) to illustrate the worst  
32  
33 case scenario that can arise from each effect during this event. Considering the maximum  
34  
35 storm surge response to the different processes at each coastal point, the additional effect of  
36  
37 rain mass ranges from 1.3 cm to 6.4 cm, the effect of rain stress is from 0.4 cm to 7.5 cm, the  
38  
39 inverse barometer effect is from 0.2 cm to 45.4 cm and the effect of wind stress is from 20.5  
40  
41 cm to 170.6 cm. While the inverse barometer effect is generally larger than the effect of rain  
42  
43 mass, it falls below the effect of rain mass at the southern boundary of the gulf (coastal points  
44  
45 180 – 330) when the cyclone is located away from the coastline. This shows that the effect of  
46  
47 rain mass can have a stronger remote presence than the inverse barometer effect. Although  
48  
49 the maximum surge level due to the effect of rain stress is higher than the effect of rain mass,  
50  
51 the effect of rain mass is larger than the effect of rain stress in most of the coastal points.  
52  
53  
54  
55  
56  
57  
58  
59  
60

ROMS (and many ocean models) include the effect of rain stress but not the rain mass even when the effect of rain mass can be larger as shown here. Ocean models that consider the effect of rain stress to be sufficiently substantial to include its effect should not neglect the effect of rain mass.

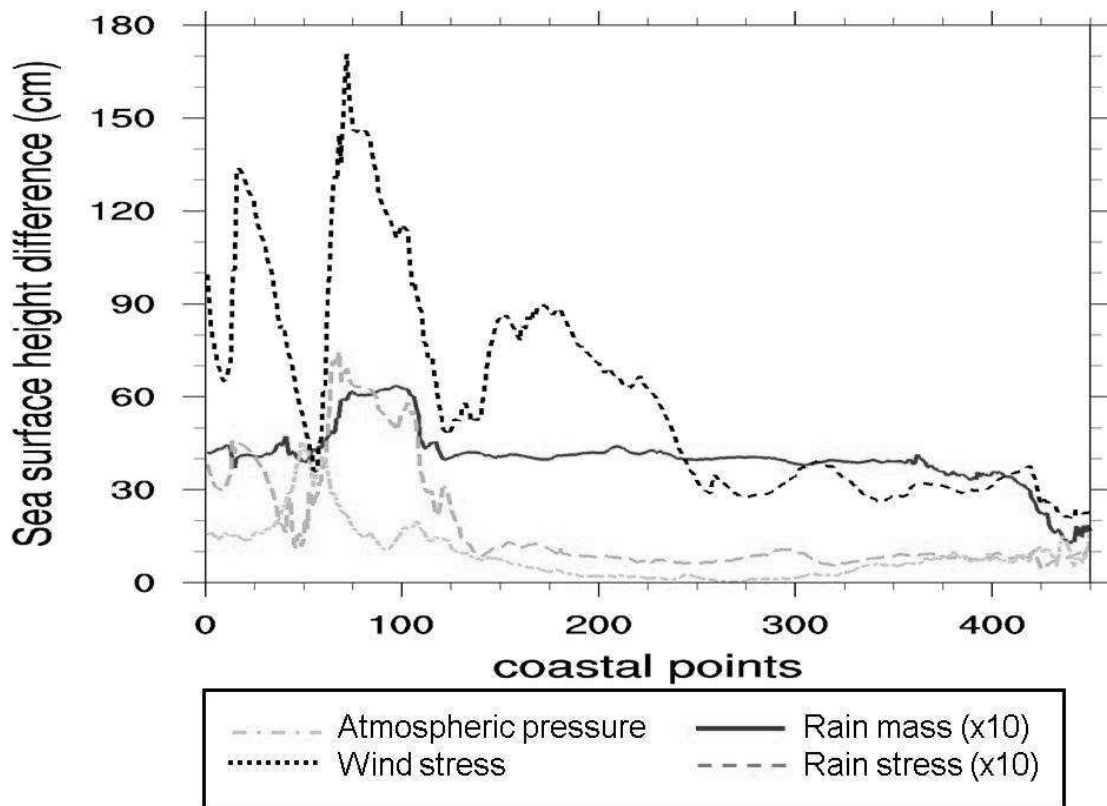


Figure 11: Comparison of the SSH change due to various physical effects, along the coastline at Gulf of Carpentaria. The coastal points can be attributed to different locations along the coastline as follows: the Wessel Islands (points 1 – 120), the western boundary of the Gulf of Carpentaria (points 121 – 180), the southern boundary of the gulf (points 181 – 280) and the eastern boundary of the gulf (points 281 – 450). SSH responses to rain mass and rain stress are scaled up by a factor of 10 for illustration.

## 5 Conclusion

Having successfully incorporated a new rain scheme into ROMS that considers the mass contribution of rainfall, its effect is shown in several cases. This new rain scheme does not interfere with the existing temperature and salinity treatment of rainfall in ROMS. The added rain mass generated surface waves that increase in propagation speed with increasing depth. An error in the existing implementation of rain stress in ROMS produced a systematic overestimation of 41% in magnitude and up to 45° in direction. This error has been corrected in this study and its isolated effect on SSH in a tropical cyclone case is shown in an idealized set-up. Finally, the effect of rain mass and rain stress in a real storm surge case (Cyclone Monica) is compared with the effect of winds and atmospheric pressure. In this case, the effect of wind and pressure generated a surge of up to 170.6 cm and 61.5 cm respectively. The effect of rain mass increases the SSH up to 6.4 cm in Buckingham Bay. The effect of rain mass has a stronger remote presence compared to the effect of pressure and rain stress. Especially in semi-enclosed areas, the effect of rain mass can substantially contribute to the storm surge even when the storm is relatively far from the coastline, and both the effect of pressure and rain stress are no longer significant. An area is subjected to high rain stress when intense winds coincide with heavy rain. Here, the effect of rain stress increases the SSH up to 7.5 cm, but the effect is generally smaller than the effect of rain mass along the coast. Many ocean models ignore the mass effect of rainfall and only incorporate the effect of wind, atmospheric pressure and rain stress. However, this study shows that the effect of rain mass can be larger than the effect of pressure and rain stress in certain scenarios, and should not be neglected. Each process shows different spatial influence and cannot be accounted for by simply scaling up a certain effect, for example, recalibrating the surge model by scaling the wind field.

**References**

- 1  
2  
3  
4  
5  
6  
7 Bleck, R., & Boudra, D. B. (1981). Initial Testing of a Numerical Ocean Circulation Model  
8  
9 Using a Hybrid (Quasi-Isopycnic) Vertical Coordinate. *Journal of Physical*  
10  
11 *Oceanography*, 11(6), 755-770. doi: Doi 10.1175/1520-  
12  
13 0485(1981)011<0755:Itoano>2.0.Co;2  
14  
15  
16 Caldwell, D. R., & Elliott, W. P. (1970). Surface stresses produced by rainfall. *J. Phys.*  
17  
18 *Oceanography*, 1, 145-148.  
19  
20  
21 Chapman, D. C. (1985). Numerical treatment of cross-shelf open boundaries in a barotropic  
22  
23 coastal ocean model. *Journal of Physical Oceanography*, 15(8), 1060-1075. doi:  
24  
25 10.1175/1520-0485(1985)015<1060:ntocso>2.0.co;2  
26  
27  
28 Chen, C. S., Qi, J. H., Li, C. Y., Beardsley, R. C., Lin, H. C., Walker, R., & Gates, K. (2008).  
29  
30 Complexity of the flooding/drying process in an estuarine tidal-creek salt-marsh  
31  
32 system: An application of FVCOM. *Journal of Geophysical Research-Oceans*,  
33  
34 113(C7). doi: Artn C07052 Doi 10.1029/2007jc004328  
35  
36  
37 Chou, M., & Suarez, and M. J. (1994). An efficient thermal infrared radiation  
38  
39 parameterization for use in general circulation models. *NASA Tech. Memo.*,  
40  
41 104606(3), 85pp.  
42  
43  
44 Clayson, C. A., & Kantha, L. H. (1999). Turbulent kinetic energy and its dissipation rate in  
45  
46 the equatorial mixed layer. *Journal of Physical Oceanography*, 29(9), 2146-2166. doi:  
47  
48 Doi 10.1175/1520-0485(1999)029<2146:Tkeaid>2.0.Co;2  
49  
50  
51 Codiga, D.L. (2011). Unified Tidal Analysis and Prediction Using the UTide Matlab  
52  
53 Functions. *Technical Report 2011-01 Graduate School of Oceanography, University*  
54  
55 *of Rhode Island, Narragansett, RI. 59pp*  
56  
57  
58  
59  
60

- 1  
2  
3 Dodla, V. B., Desamsetti, S., & Yerramilli, A. (2011). A Comparison of HWRF, ARW and  
4  
5 NMM Models in Hurricane Katrina (2005) Simulation. *International Journal of*  
6  
7 *Environmental Research and Public Health*, 8(6), 2447-2469. doi: DOI  
8  
9 10.3390/ijerph8062447  
10  
11  
12 Dudhia, J. (1989). Numerical Study of Convection Observed during the Winter Monsoon  
13  
14 Experiment Using a Mesoscale Two-Dimensional Model. *Journal of the Atmospheric*  
15  
16 *Sciences*, 46(20), 3077-3107. doi: Doi 10.1175/1520-  
17  
18 0469(1989)046<3077:Nsocod>2.0.Co;2  
19  
20  
21 Dudhia, Jimy. (1996). A multi-layer soil temperature model for MM5. *Sixth PSU/NCAR*  
22  
23 *Mesoscale Model Users' Workshop*.  
24  
25  
26 Durden, S. L. (2010). Remote Sensing and Modeling of Cyclone Monica near Peak Intensity.  
27  
28 *Atmosphere*, 1(1), 15-33. doi: Doi 10.3390/Atmos1010015  
29  
30  
31 Egbert, G. D., Bennett, A. F., & Foreman, M. G. G. (1994). TOPEX/POSEIDON tides  
32  
33 estimated using a global inverse model. *Journal of Geophysical Research-Oceans*,  
34  
35 99(C12), 24821-24852. doi: 10.1029/94jc01894  
36  
37  
38 Fairall, C. W., Bradley, E. F., Rogers, D. P., Edson, J. B., & Young, G. S. (1996). Bulk  
39  
40 parameterization of air-sea fluxes for Tropical Ocean Global Atmosphere Coupled  
41  
42 Ocean Atmosphere Response Experiment. *Journal of Geophysical Research-Oceans*,  
43  
44 101(C2), 3747-3764. doi: Doi 10.1029/95jc03205  
45  
46  
47  
48 Flather, R. A. (1976). A tidal model of the northwest European continental shelf *Mémoires de*  
49  
50 *la Société Royale de Sciences de Liege* (Vol. 6, pp. 141-164).  
51  
52  
53  
54 Gill, A. E., and P. P. Niiler (1973). The theory of the seasonal variability in the ocean, *Deep*  
55  
56 *Sea Res.*, 20, 141-177  
57  
58  
59  
60

- 1  
2  
3 Griffies, S. M., & Hallberg, R. W. (2000). Biharmonic friction with a Smagorinsky-like  
4  
5 viscosity for use in large-scale eddy-permitting ocean models. *Monthly Weather*  
6  
7 *Review*, 128(8), 2935-2946.  
8  
9
- 10 Hirose, N., I. Fukumori, and R. M. Ponte (2001). A non-isostatic global sea level response to  
11  
12 barometric pressure near 5 days, *Geophys. Res. Lett.*, 28, 2441–2444,  
13  
14 *doi:10.1029/2001GL012907*.  
15  
16
- 17 Hong, S. Y., Noh, Y., & Dudhia, J. (2006). A new vertical diffusion package with an explicit  
18  
19 treatment of entrainment processes. *Monthly Weather Review*, 134(9), 2318-2341.  
20  
21 *doi: Doi 10.1175/Mwr3199.1*  
22  
23
- 24 Huffman, G. J., Adler, R. F., Bolvin, D. T., Gu, G. J., Nelkin, E. J., Bowman, K. P., . . .  
25  
26 Wolff, D. B. (2007). The TRMM multisatellite precipitation analysis (TMPA): Quasi-  
27  
28 global, multiyear, combined-sensor precipitation estimates at fine scales. *Journal of*  
29  
30 *Hydrometeorology*, 8(1), 38-55. *doi: Doi 10.1175/Jhm560.1*  
31  
32
- 33 Janjic, Z. I. (1994). The Step-Mountain Eta Coordinate Model - Further Developments of the  
34  
35 Convection, Viscous Sublayer, and Turbulence Closure Schemes. *Monthly Weather*  
36  
37 *Review*, 122(5), 927-945. *doi: Doi 10.1175/1520-*  
38  
39 *0493(1994)122<0927:Tsmecm>2.0.Co;2*  
40  
41  
42
- 43 Jarvinen, B. R., & Lawrence, M. B. (1985). An Evaluation of the Slosh Storm-Surge Model.  
44  
45 *Bulletin of the American Meteorological Society*, 66(11), 1408-1411.  
46  
47
- 48 Kain, J. S. (2004). The Kain-Fritsch convective parameterization: An update. *Journal of*  
49  
50 *Applied Meteorology*, 43(1), 170-181. *doi: Doi 10.1175/1520-*  
51  
52 *0450(2004)043<0170:Tkcpan>2.0.Co;2*  
53  
54  
55  
56  
57  
58  
59  
60

- 1  
2  
3 Lin, Y. L., Farley, R. D., & Orville, H. D. (1983). Bulk Parameterization of the Snow Field in  
4  
5 a Cloud Model. *Journal of Climate and Applied Meteorology*, 22(6), 1065-1092. doi:  
6  
7 Doi 10.1175/1520-0450(1983)022<1065:Bpotsf>2.0.Co;2  
8  
9
- 10 Machta, L. (1969). Evaporation rates based on tritium measurements for hurricane Betsy.  
11  
12 *Tellus*, 21(3), 404-408.  
13  
14
- 15 Marchesiello, P., McWilliams, J. C. , & Shchepetkin, A. F. . (2001). Open boundary  
16  
17 conditions for long-term integration of regional ocean models. *Ocean Modelling*, 3, 1-  
18  
19 20.  
20  
21
- 22 Mlawer, E. J., Taubman, S. J., Brown, P. D., Iacono, M. J., & Clough, S. A. (1997). Radiative  
23  
24 transfer for inhomogeneous atmospheres: RRTM, a validated correlated-k model for  
25  
26 the longwave. *Journal of Geophysical Research-Atmospheres*, 102(D14), 16663-  
27  
28 16682. doi: Doi 10.1029/97jd00237  
29  
30
- 31 Paulson, C. A. (1970). The mathematical representation of wind speed and temperature  
32  
33 profiles in the unstable atmospheric surface layer. *J. Appl. Meteor.*, 9, 857-861.  
34  
35
- 36 Peng, M. C., Xie, L., & Pietrafesa, L. J. (2004). A numerical study of storm surge and  
37  
38 inundation in the Croatan-Albemarle-Pamlico Estuary System. *Estuarine Coastal and*  
39  
40 *Shelf Science*, 59(1), 121-137. doi: DOI 10.1016/j.ecss.2003.07.010  
41  
42
- 43 Pickard, G. L., & Emery, W. J. (1990). *Descriptive physical oceanography: an introduction*.  
44  
45 Elsevier.  
46  
47
- 48 Ponte, R. M. (2006). Oceanic response to surface loading effects neglected in volume-  
49  
50 conserving models. *Journal of Physical Oceanography*, 36(3), 426-434. doi:  
51  
52 10.1175/Jpo2843.1  
53  
54  
55  
56  
57  
58  
59  
60



- 1  
2  
3 Rego, J. L., & Li, C. Y. (2010). Storm surge propagation in Galveston Bay during Hurricane  
4  
5 Ike. *Journal of Marine Systems*, 82(4), 265-279. doi: DOI  
6  
7 10.1016/j.jmarsys.2010.06.001  
8  
9  
10 Riehl H, Malkus J. (2011). Some aspects of hurricane Daisy, 1958. *Tellus A* 13(2)  
11  
12  
13 Roden, G. I., & Rossby, H. T. (1999). Early Swedish contribution to oceanography: Nils  
14  
15 Gissler (1715-71) and the inverted barometer effect. *Bulletin of the American*  
16  
17 *Meteorological Society*, 80(4), 675-682. doi: Doi 10.1175/1520-  
18  
19 0477(1999)080<0675:Escton>2.0.Co;2  
20  
21  
22 Rogers, Eric, , Thomas Black, , Brad Ferrier, , Ying Lin, , David Parrish, & , and Geoffrey  
23  
24 DiMego. (2001). National Oceanic and Atmospheric Administration Changes to the  
25  
26 NCEP Meso Eta Analysis and Forecast System: Increase in resolution, new cloud  
27  
28 microphysics, modified precipitation assimilation, modified 3DVAR analysis. *NOAA*.  
29  
30  
31 Saha, S., Moorthi, S., Pan, H. L., Wu, X. R., Wang, J. D., Nadiga, S., Goldberg, M. (2010).  
32  
33 The Ncep Climate Forecast System Reanalysis. *Bulletin of the American*  
34  
35 *Meteorological Society*, 91(8), 1015-1057. doi: Doi 10.1175/2010bams3001.1  
36  
37  
38 Shchepetkin, A. F., & McWilliams, J. C. (2005). The regional oceanic modeling system  
39  
40 (ROMS): a split-explicit, free-surface, topography-following-coordinate oceanic  
41  
42 model. *Ocean Modelling*, 9(4), 347-404. doi: DOI 10.1016/j.ocemod.2004.08.002  
43  
44  
45  
46 Sheng, Y. P., Alymov, V., & Paramygin, V. A. (2010). Simulation of storm surge, wave,  
47  
48 currents, and inundation in the Outer Banks and Chesapeake Bay during Hurricane  
49  
50 Isabel in 2003: The importance of waves. *Journal of Geophysical Research-Oceans*,  
51  
52 115. doi: Artn C04008 Doi 10.1029/2009jc005402  
53  
54  
55  
56  
57  
58  
59  
60

- 1  
2  
3 Sheng, Y. P., Davis, J. R., Figueiredo, R., Liu, B., Liu, H., Luettich, R., Zheng, L. (2012). A  
4  
5 regional testbed for storm surge and coastal inundation models - an overview.  
6  
7 *Estuarine and Coastal Modeling, American Society of Civil Engineers*, 476-495.  
8  
9  
10 Skamarock, W. C, Klemp, J. B, Dudhia, J, Gill, D. O, Baker, D. M, Duda, M. G, JG., Power.  
11  
12 (2008). A Description of the Advanced Research WRF Version 3. . *NCAR/TN-*  
13  
14 *475+STR*.  
15  
16  
17 Smith, W. H. F., and Sandwell D. T. . (1997). Global seafloor topography from satellite  
18  
19 altimetry and ship depth soundings. *Science*, 277, 1957-1962.  
20  
21  
22 Stewart, Robert H. (2009). Introduction to Physical Oceanography. *University Press of*  
23  
24 *Florida*.  
25  
26  
27 Umlauf, L., & Burchard, H. (2003). A generic length-scale equation for geophysical  
28  
29 turbulence models. *Journal of Marine Research*, 61(2), 235-265. doi:  
30  
31 10.1357/002224003322005087  
32  
33  
34 Van Dorn, W. G. (1953). Wind stress on an artificial pond. *Journal of marine research*, 12,  
35  
36 249-276.  
37  
38  
39 Wang, Shuai, Toumi, Ralf, Czaja, Arnaud, & Van Kan, Adrian. (2015). An Analytic Model  
40  
41 of Tropical Cyclone Wind Profiles. *Quarterly Journal of the Royal Meteorological*  
42  
43 *Society*, n/a-n/a. doi: 10.1002/qj.2586  
44  
45  
46 Watson, P. J. (2011). Is There Evidence Yet of Acceleration in Mean Sea Level Rise around  
47  
48 Mainland Australia? *Journal of Coastal Research*, 27(2), 368-377. doi: Doi  
49  
50 10.2112/Jcoastres-D-10-00141.1  
51  
52  
53 Weisberg, R. H., & Zheng, L. Y. (2006). Hurricane storm surge simulations for Tampa Bay.  
54  
55 *Estuaries and Coasts*, 29(6A), 899-913.  
56  
57  
58  
59  
60

- 1  
2  
3 Westerink, J. J., Luettich, R. A., Baptista, A. M., Scheffner, N. W., & Farrar, P. (1992). Tide  
4  
5 and Storm-Surge Predictions Using Finite-Element Model. *Journal of Hydraulic*  
6  
7 *Engineering-Asce*, 118(10), 1373-1390. doi: Doi 10.1061/(Asce)0733-  
8  
9 9429(1992)118:10(1373)  
10  
11  
12 Westerink, J. J., Luettich, R. A., Feyen, J. C., Atkinson, J. H., Dawson, C., Roberts, H. J.,  
13  
14 Pourtaheri, H. (2008). A basin- to channel-scale unstructured grid hurricane storm  
15  
16 surge model applied to southern Louisiana. *Monthly Weather Review*, 136(3), 833-  
17  
18 864. doi: Doi 10.1175/2007mwr1946.1  
19  
20  
21  
22 Wu, C. C. (2013). TYPHOON MORAKOT Key Findings from the Journal TAO for  
23  
24 Improving Prediction of Extreme Rains at Landfall. *Bulletin of the American*  
25  
26 *Meteorological Society*, 94(2), 155-160. doi: Doi 10.1175/Bams-D-11-00155.1  
27  
28  
29 Zhong, L. J., Li, M., & Zhang, D. L. (2010). How do uncertainties in hurricane model  
30  
31 forecasts affect storm surge predictions in a semi-enclosed bay? *Estuarine Coastal*  
32  
33 *and Shelf Science*, 90(2), 61-72. doi: DOI 10.1016/j.ecss.2010.07.001  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60