A Scalable Method for the Reductive Dissolution, Purification, and Separation of Single-Walled Carbon Nanotubes

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Abstract: As synthesized, bulk single-walled carbon nanotube (SWNT) samples are typically highly agglomerated and heterogeneous. However, their most promising applications require the isolation of individualized, purified nanotubes, often with specific optoelectronic characteristics. A wide range of dispersion and separation techniques have been developed, but the use of sonication or ultracentrifugation imposes severe limits on scalability, and may introduce damage. Here, we demonstrate a new, intrinsically-scalable method for SWNT dispersion and separation, using reductive treatment in sodium metal-ammonia solutions, optionally followed by selective dissolution in a polar aprotic organic solvent. In situ small angle neutron scattering demonstrates the presence of dissolved, unbundled SWNTs in solution, at concentrations reaching at least 2 mg/ml; the ability to isolate individual nanotubes is confirmed by atomic force microscopy. Spectroscopy data suggest that the soluble fraction contains predominately large metallic nanotubes; a potential new mechanism for nanotube separation is proposed. In addition, the G/D ratios observed during the dissolution sequence, as a function of metal:carbon ratio, demonstrate a new purification method for removing carbonaceous impurities from pristine SWNTs, which avoids traditional, damaging, competitive oxidation reactions.

The extraordinary electronic and optical properties of single walled carbon nanotubes (SWNTs) are determined by their structure; SWNTs are either metallic or semiconducting dependent on their diameter and chirality¹. Many valuable applications in, for example, transparent conductors,² solar cells,³ biosensors⁴ and nanoelectronics,⁵ require individualised carbon nanotubes of specific electronic character and high purity. For this reason, there has been a concerted effort to provide methods for separation of individual SWNT species. Despite recent progress in selective synthesis, current techniques still produce heterogeneous samples containing SWNTs of varying geometry and electronic character, as well as other carbonaceous contaminants.⁶⁻⁷ Post-production separation of metallic and semiconducting SWNTs follows various strategies based on physical (dielectrophoresis,⁸ density gradient ultracentrifugation,⁹ gel electrophoresis¹⁰ and chromatography¹¹) and chemical (diazonium salts,¹² ozonolysis,¹³ diporphyrin,¹⁴ bromine,¹⁵ amine¹⁶ and pyrene functionalisation¹⁷) means. All of

these methods suffer from problems with scalability, effectiveness, and often dissolution sensitivity. Additionally, chemical functionalisation methods are limited by the inherent potential to damage the outstanding electronic, optical and mechanical properties of interest.¹⁸

Whilst some larger volume separation techniques have been demonstrated^{8,10,19}, the initial dispersion step remains limiting. SWNTs form bundles due to strong van der Waals interactions; so, before any fractionation attempt, the SWNTs must be individualised by high power sonication in organic solvent or, most commonly, aqueous surfactant solution.^{8,10,16} Sonication is known to damage the SWNT structure by both shortening and functionalisation.²⁰ The resulting solutions are intrinsically low concentration and almost always require ultracentrifugation for the complete removal of small bundles,²¹ greatly limiting the scalability of these methods. The only means to dissolve SWNTs without sonication rely on charging effects, either protonation in superacids²² or reduction (indirectly) with alkali metals.²³ Penicaud *et al.* showed that SWNTs can be reduced with sodium naphthalene in THF, leading to spontaneous dissolution in various aprotic polar solvents, and the identification of individual SWNTs in subsequent, dry AFM.²³ However, no selective electronic separation of nanotubes or impurities has been reported using either approach; in Penicaud's method, the sodium naphthalide remains in the system as a contaminant.

Using liquid ammonia as the charging solvent avoids this contamination problem, since the ammonia can be simply removed by evaporation. When an alkali metal is dissolved in liquid ammonia, solvated electrons are formed (seen as an intense blue color within the solution²⁴); these solvated electrons can then be transferred to the SWNTs.²⁵⁻²⁸ The resulting reduced SWNTs have been reported previously to disperse more easily with "continued stirring" in liquid ammonia,²⁹ with AFM evidence of individual SWNTs after chemical grafting reactions.²⁶ The mechanism for this reactive debundling is still under some debate, with intertube repulsions and alkali metal intercalation most commonly suggested.^{26,29} Using ammonia as a solvent, Wunderlich *et al.* suggested that metallic and small diameter SWNTs are

preferentially chemically functionalised due to strain and charging effects,^{28,30} although no fractionation was undertaken.

In this work, we use *in situ* small angle neutron scattering (SANS) to demonstrate that SWNTs can indeed be unbundled to give individual tubes in solution by reductive charging in ammonia. This process requires control of the metal:carbon (M:C) ratio, to much lower values (approximately 1:10) than typical in the literature ($>1:1$), but does not inherently involve covalent functionalisation of the tubes. Transfer of these reduced nanotubes, as a dry powder, to an organic solvent allows for straightforward fractionation of the spontaneously dissolved material, which was found to contain predominantly larger metallic SWNTs and impurities.

b

Figure 1. Liquid ammonia reduction of ARC SWNTs. **a.** Process scheme for the reduction and subsequent dissolution in sodium-ammonia, illustrated by atomistic models and photographs of relevant phases (M:C 1:20). **b.** SANS pattern from the spontaneously dissolved SWNTs (M:C 1:24) in sodium-ammonia (Na:ND₃). The red and green lines show power-law best fits, giving exponents of -1.0 and -2.65, as marked. The insert shows a schematic of the corresponding nanotube mesh, and associated length scales.

Results and Discussion

SWNT reduction in liquid ammonia

The initial treatment of SWNTs in liquid sodium metal ammonia solution is a key step in the separation. Here, unlike in the THF/naphthalide route, the electrons are directly solvated and no secondary charge transfer agent is required. When liquid ammonia is condensed onto a mixture of ARC SWNT and alkali metal, the initially colorless liquid ammonia turns blue, due to the dissolution of the metal and the concomitant 1s - 2p transition of solvated electrons.²⁴ In the absence of stirring, the SWNT powder

swells, and gradually over the course of approximately two hours, the solution changes from blue to clear to black, as the nanotubes accept the solvated electrons, spontaneously debundle, and dissolve into the liquid ammonia (reaction scheme Fig. 1a).

The structure of the SWNTs in liquid sodium-ammonia was characterized by small angle neutron scattering (SANS). SANS is a powerful technique for probing the structure of nanoparticles in solution³¹⁻³⁴, and specifically can be used to determine whether the SWNTs are present as isolated species or as bundles or clusters.³³ In the dilute regime, the SANS intensity is $I(Q) \propto Q^{-D}$, where $Q = 4\pi \sin(\theta) / \lambda$ is the magnitude of the scattering wavevector and D is the fractal dimension of the scattering objects. Thus a plot of $\log_{10} I(Q)$ against $\log_{10} Q$ will yield a straight line of gradient $-D$. Fully dispersed rod-like objects such as individual SWNTs will have a dimensionality $D \approx 1$, and their expected SANS signal is therefore $I(Q) \propto Q^{-1}$. Heterogeneous dispersions of SWNTs, consisting of agglomerates and bundles (effectively rod networks), will exhibit a larger fractal dimension^{31,32,34}, typically in the range $D \approx 2$ to 4.

SANS data from a solution of ARC SWNTs dissolved in sodium-deuteroammonia (ND₃) solution show two clear scattering regimes (Fig. 1b). At higher Q -values, up to at least $Q \approx 0.2 \text{ \AA}^{-1}$, there is a Q^{-1} dependence that can be quantitatively fitted by a model containing cylindrical rods of diameter 15-20 Å (see Methods). This dimension is entirely consistent with the value of $\sim 14 \text{ \AA}$ reported for the diameter of ARC-SWNTs, since each tube will be surrounded by a dense ordered shell of solvent (Fig 1a). This solvation behavior is expected to be similar to that deduced from neutron diffraction studies of C₆₀ anions dissolved in ammonia, where isotopic labeling and the intrinsic monodispersity allowed the solvent structure to be explicitly determined.³⁵ However, we note that since SANS intensity typically scales as the particle volume squared, scattering from any fullerene species (as measured previously) will fall below the sensitivity of our data, in the current case (see ESI Figure S1). Conversely, if the dissolved material consisted of a significant fraction of graphitic nanoparticles that are known to occur

in arc soots³⁶ they would contribute noticeably to the SANS signal, giving a lower dimensionality; thus, a large fraction of dissolved nanoparticles is not consistent with the Q^{-1} dependence (rods) observed.

The second regime in figure 1b has $Q^{-2.65}$ behavior, with the cross-over indicating the smallest size of the larger scatterers. For concentrated solutions of SWNTs, effectively, there is a mesh formed³⁴ (see Fig.1b inset) which gives rise to this higher power law scattering at lower Q . Previously, a $Q^{-2.5}$ power law has been proposed in this regime³⁴, based on infinitely thin rods; the slightly higher value of the fractal dimension (2.65) obtained here, directly from a free fit, is consistent with rods of finite thickness, and is identical to that obtained from detailed analysis of SEM and TEM images of carbon nanotube films³⁷. The crossover in Fig 1b is found at $Q \approx 0.055 \text{ \AA}^{-1}$, equivalent to a mesh size $\xi \approx 115 \text{ \AA}$. From this model, it is possible to deduce that the approximate concentration of individual SWNTs in solution is in the range 2-12 mg/ml (up to 0.75 % SWNT by volume, see Methods). This value is consistent with Hough *et al.*³³ who reported that, in sodium dodecylbenzenesulfonate (NaDDBS) stabilized aqueous suspension, $\xi = 150 \text{ \AA}$ corresponds to a 0.75 wt% SWNT solution.

DMF addition to the nanotubide salt – SWNT separation

Following the removal of the liquid ammonia, a dry powder of sodium ‘nanotubide’ is formed (where ‘nanotubide’ is the proposed term for a pure nanotube anion). On the addition of dry DMF (Fig. 2), the ARC nanotubide swells and a fraction (approximately 40 wt%) of the SWNTs spontaneously dissolves. It is worth emphasising that no stirring, and more importantly, no ultrasound was used at any stage of the process. As in the work of Penicaud²³, this spontaneous dissolution is presumably driven by the solvation of the cations, leading to repulsion between the solvated nanotubide anions and the formation of an electrostatically-stabilised colloid (or polyelectrolyte molecule). The use of oxygen and water free solvent is essential to avoid quenching the charge through the formation of sodium (hydr)oxide. The solution in organic solvent allows straightforward fractionation via cannula, although reliable

discrimination of the phases becomes difficult at higher concentration. The maximum measurable concentration was found to be 2 mg/ml although, again, this value should be considered a lower bound.

Sodium nanotubide based on CoMoCAT SWNTs, produced by the same ammonia process, also spontaneously dissolves in DMF (dissolved yield 14 wt%) with a similar minimum solubility (2 mg/ml). SEM (Fig. 3a) and AFM (Fig. 3b & 3c) confirm that the dissolved fraction consists predominantly of nanotubes; the AFM, in particular, shows that the dissolved species are individualized SWNTs with diameters 1.0-1.5 nm.

Figure 2. Photographs showing the spontaneous dissolution of 10mg ARC nanotubide salt (M:C 1:20) into 10 cm³ dry DMF.

Figure 3. Microscopy images of CoMoCAT SWNTs spontaneously dissolved into dry DMF (M:C 1:10). **a.** SEM X and Y-scale 2 μm . **b.** AFM X and Y-scale 3 μm , Z-scale 0-45 nm. **c.** Height section from AFM images showing the height of three SWNTs in different areas of the sample.

The relative metallicity of the SWNT fractions was estimated from the positions of the radial breathing modes (RBMs) observed using Raman spectroscopy; the SWNT diameters³⁸ ($\omega_{\text{rbm}} = 218.7/d + 15.3$) were correlated with the diameter dependent transition energies between the Van Hove singularities using the Kataura plot³⁹, to assign bands to particular electronic type. Whilst the assessment is highly diameter dependent, selecting a probe wavelength in resonance with both metallic and semiconducting species provides an effective means to monitor relative changes with the sample. Both metallic and semiconducting species within the as-supplied CoMoCAT SWNTs conveniently couple to

the red (633 nm) and green (532 nm) lasers available. Excitation using the red laser, shows 7 distinct peaks; using previously established methods,^{38,39} the tube diameters can be calculated and the peaks indexed (Fig. 4a). The intensity of the tube peaks is only comparative, not absolute, as the efficiency of resonant energy transfer varies;⁴⁰ it is known, for example, that (6, 5) tubes dominate CoMoCAT samples, although they couple weakly to the red laser. In the spontaneous dissolution process, the intensity of the peaks correlating to the semiconducting SWNTs decrease strongly, disappearing almost entirely. The Raman data also suggest a diameter effect, with preferential dissolution of larger metallic nanotubes. Although apparent shifts in RBM distributions have been reported on debundling,⁴¹ due to changes in interband transition energies, no shifts were observed in the metallic region of the spectra, indicating any such effects are minimal; most likely, SWNTs rebundle during charge quenching in suspension before sampling for Raman measurements. Related overall trends, dependent on both electronic character and size, have been observed during the deliberate covalent functionalisation of charged SWNTs;^{28,30} in the current case, no explicit reagents were added and the possibility of significant, accidental, selective functionalisation is addressed by experiments discussed in the mechanism section below. The green Raman laser line (Fig. 4b) shows a similar relative decrease in intensity of the semiconducting peaks in the spontaneously-dissolved fraction, both before and after vacuum annealing. The Raman data for the original salt and subsequent dissolved fractions are homogeneous across the sample; the spectra for the remaining undissolved fractions are more heterogeneous, presumably reflecting the difficulty in isolating the sediment from the last drops of dispersion and the potential for otherwise soluble nanotubes to remain trapped due to entanglements or amorphous debris.

UV/Vis spectroscopy has the potential to provide complementary evidence of the degree of separation. Indeed, a reduction in the size of the semiconducting peak is observed following the dissolution in the DMF, Fig. 4c. The first Van Hove transition of the metallic species can be seen in the 350-500 nm range, although the low intensity of the peaks in the as-received sample parallels the small

fraction of metallic SWNTs present. The second Van Hove transitions of the semiconducting SWNTs fall in the 480-700 nm range. In this range, the largest peak at 560 nm relates to the (6,5) SWNTs which is known⁴⁰ to be the dominant tube in this CoMoCAT material. Despite extensive attempts using different approaches, the signal from the dissolved fraction remains weak; the result may relate to an increased concentration of impurities, as indicated by the Raman studies of the ‘D-peak’ discussed below, limited functionalisation during quenching, and/or the difficulty in redispersing dense bundles formed on drying individualized SWNTs from solution. Note that spectra cannot be obtained from the spontaneously-dissolved SWNTs due to the bleaching of optical transitions associated with charging or doping of the Van Hove features.⁴²

Figure 4. Optical characterisation of SWNT fraction spontaneously dissolved into DMF at M:C ratio 1:10 **a.** Red (633nm) RBM Raman Spectra of as-received CoMoCAT SWNTs (solid line), the spontaneously-dissolved CoMoCAT fraction (dashed line) and the spontaneously-dissolved CoMoCAT fraction following vacuum annealing (dotted line). The shaded area represents the cross over between the metallic and semiconducting SWNTs. **b.** Green (532nm) RBM Raman Spectra of as-received CoMoCAT SWNTs (solid line), the spontaneously-dissolved CoMoCAT fraction (dashed line) and the spontaneously-dissolved CoMoCAT fraction followed by vacuum annealing (dotted line). **c.** UV/Vis spectra of as-received CoMoCAT SWNTs (solid line) and spontaneously-dissolved fraction (dotted line).

Raman spectroscopy provides a semi-quantitative indication of purity or degree of functionalisation by comparing the relative intensity of the defect band (D-band) at $\sim 1350\text{ cm}^{-1}$ and the graphite band (G-band) at $\sim 1580\text{ cm}^{-1}$; although the details are complex, increasing D-band intensity is usually correlated either with the presence of other contaminating carbons, or damage to the nanotube framework. The initial G/D ratio of around 11 does not change following the ammonia treatment even after air-exposure, consistent with a lack of significant covalent functionalisation (see ESI Figure S2 for a summary plot of the G/D ratios obtained). Similarly, Energy-dispersive X-ray spectroscopy (EDX) showed no evidence of nitrogen incorporation in any of the liquid ammonia or DMF treated samples. However, the G/D ratio

drops to 7 in the spontaneously-dissolved fraction, whilst increasing to 18 in the undissolved material. This divergent trend suggests that carbonaceous impurities (or short/defective SWNTs) are preferentially dissolved with or before the metallic fraction, although the dissolved material is predominantly SWNTs at our standard M:C ratios, as discussed above. This hypothesis was confirmed using a much lower M:C ratio (1:100); in this case, only a small mass fraction dissolved spontaneously, giving a very low G/D ratio around unity, leaving the bulk material as a residue, again with an improved G/D of 18. Thus, the overall method can be used for purification as well as dispersion and separation. SEMs of the remaining, undissolved material (see ESI Figure S2), show well defined, apparently cleaner, SWNTs.

Yield and scalability

The yield (Fig. 5) of spontaneously-dissolved CoMoCAT SWNTs in DMF depends strongly on metal to carbon ratio (M:C). The optimum atomic ratio of 1:10 M:C consistently corresponds to a maximum dissolved yield of 15 ± 1 wt% of the original sample. Further washings with fresh solvent do not dissolve additional material; the remaining, undissolved SWNTs are either inappropriately charged, or physically trapped by entanglements or insoluble amorphous carbon. At lower metal loadings, the yield declines due to insufficient charge for electrostatic repulsion between the nanotubes and/or charge-driven solvation. At higher loadings, the yield declines due to excessive sodium ion concentration, leading to ‘salting’ out, as the Debye length declines. A similar optimum (1:12) was identified previously for the dissolution of C_{60}^{5-} in metal-ammonia solutions.⁴³ It is worth noting that the great majority of literature studies of nanotubes in ammonia use a metal to carbon ratio of 1:1 or higher which is not expected to yield individualized nanotubes, either for debundled functionalisation or separation. The reaction was found to be scalable at the optimum 1:10 M:C ratio. On increasing the initial quantity of SWNTs from 20mg to 50mg and 100mg, the reaction proceeded in the same manner, yielding 14wt% and 17wt% of spontaneously-dissolved nanotubes, respectively. Similar Raman spectra were obtained

for the spontaneously dissolved component in all cases (see ESI Figure S3)

Figure 5. Mass fraction dissolved spontaneously, following metal-ammonia and DMF treatment; 20mg of CoMoCAT SWNTs were used in each experiment.

Mechanism

The dissolution process is driven by the solvation of the sodium cations, electrostatic repulsion between the nanotubes, and the relatively favorable enthalpic interactions of nanotubes with both amides⁴⁴ and ammonia, particularly when charged.⁴³ The dominant importance of the charged state is clearly demonstrated (Fig. 6) by the long term stability of the DMF dispersions in the absence of air (>6 months), but rapid agglomeration on quenching the charge with air (minutes). Whilst the use of DMSO as a solvent for nanotubide has been reported to lead to covalent modification and permanent modification of solubility⁴⁵, the sensitivity to air in the current case strongly suggests that the solution stability is due to electrostatic repulsion between charged nanotubes rather than any permanent chemical reaction with the ammonia or DMF. Exposure to oxygen (in air) is reported to quench the charge on reduced SWNTs without oxidative addition.^{46,47} Whilst adventitious moisture may have an influence, a low level of functionalization is implied, in this study, by the G/D Raman data above and, specifically,

by vacuum-annealing the spontaneously-dissolved sample at 500°C, 10^{-6} mbar, for 24 hours. This process is often used to remove unwanted functional groups and (partially) restore original optoelectronic properties⁴⁸; however, the Raman spectra of the dissolved SWNT fraction was unchanged by this treatment. The spontaneous nature of the dissolution confirms that it is thermodynamically favored, and therefore likely to give rise to individualized SWNTs. Previous estimates for the Bjerrum length²³, the minimum separation of surface charge before counter-ion condensation becomes thermodynamically favorable, suggest that a 1:50 M:C ratio should be optimal for nanotubes; very recent refinements have identified a slightly lower optimum ratio (<1:20) and provide a detailed discussion of the thermodynamics of the electrostatic stabilization.⁴⁹ However, these estimates rely entirely on the dielectric constant of the solvent; in reality, the dielectric constant of the nanotube will make a major contribution, particularly for metallic tubes. Qualitatively, the higher dielectric constant of the metallic SWNTs should reduce the effective Bjerrum length, increasing the maximum surface charge that can be accommodated before ‘salting’ out occurs, and raising the electrostatic stability of the dispersed state relative to the semiconducting SWNTs; larger diameter nanotubes are also expected to have higher dielectric constants.⁵⁰ Detailed calculations require a revision of the traditional theory, beyond the scope of this study, but would be of great interest. Whilst this hypothesis for the size and electronic selectivity of the dissolution process appears most plausible, the variation of dielectric constant on charging SWNTs is not yet known, and other selective mechanisms may contribute. Intrinsically, different nanotubes have characteristic electron affinities, and thus charge will be partitioned heterogeneously when in contact. Whilst the Fermi level appears to be a function of diameter,⁵¹ metallic nanotubes will always have greater initial electron affinity than semiconducting species, an effect exploited in selective diazonium functionalisation reactions.¹² In the present case, however, there is sufficient charge available to saturate at least the first Van Hove singularities, and the distinction between the multiple electron affinities of the various nanotube types is less obvious. Although charge effects clearly dominate, the selective solubilizing interactions of DMF

may also play a role.⁵²

Figure 6. The spontaneously-dissolved fraction from a 1:10 metal to carbon ratio reduction stored in inert atmosphere for 6 months, 1 minute and 12 hours after exposure to air.

Conclusions

Liquid ammonia has been shown to be an excellent medium for creating powders of alkali metal nanotubide salts, avoiding the need of any additional charge transfer species through the intrinsic electron-solvating power of the ammonia. These salts can be conveniently dissolved in dry amide solvents to form solutions of individual charged SWNTs that can be readily handled. The dissolution of a variety of as-synthesised SWNTs occurs without any sonication, stirring, or chemical functionalisation, as long as the metal:carbon ratio is carefully controlled (to lower levels than typical in the literature); the high aspect ratios and characteristics of pristine SWNT can thus be retained. The process was found to work similarly for HiPco SWNTs (see ESI Figure S4). At very low metal contents, carbonaceous impurities are selectively removed, providing a non-destructive means of SWNT purification. At high metal contents, the spontaneously-dissolved fraction appears to consist mainly of larger metallic SWNTs, suggesting the possibility of straightforward separation by size or electronic character without use of ultrasound or centrifugation. The exact nature of the separated samples could be clarified by specialists in individual SWNT methods; whilst laborious, measurements of individual properties, particularly transport characteristics, are highly relevant to application. This new separation

process has been shown to scale consistently to the 100mg level, and there is no intrinsic reason that it could not provide larger scale separation and/or purification of SWNTs in the future (see ESI Figure S3). The availability of large quantities of individualized, purified, undamaged SWNTs will enable further fundamental studies of SWNT phenomena and a wide range of applications.

Methods

Preparation of separated SWNTs. SWNTs were purchased from Carbon Solutions (ARC grown) and Southwestern Nanotechnologies (CoMoCat, CVD grown) and used as-received. In a typical experiment, 20 mg SWNTs was accurately weighed in a high purity argon glovebox and placed in a specially designed glass reaction tube which was then outgassed at 400°C for 48 hours, at $<10^{-6}$ mbar. 2 mg of sodium was accurately weighed and transferred to the reaction tube which was cooled at -50°C in either a propanol/dry ice bath or in propanol cooled using a chiller unit. 0.12 moles of ammonia (Aldrich at 99.9% purity) was then condensed onto the reaction mixture and after approximately 2 hours and the ammonia was removed from the mixture by cryo-pumping. With the complete exclusion of air, dry DMF (Aldrich, anhydrous 99.8%) was then added to the SWNT powder using standard Schlenk line techniques after which a fraction could be seen to dissolve spontaneously. Samples were vacuum annealed at 500 °C under a vacuum of 10^{-6} mbar for 24 hours. In some experiments, the mass of SWNTs and/or metal was varied to control metal to carbon ratio, or absolute reaction scale.

Raman spectra measurements. Raman spectra of dry samples were collected with a LabRam Infinity Raman Instrument using red (633 nm) and green (532 nm) lasers.

Optical absorbance measurements. UV/Vis/Near-IR spectra were collected using a Perkin Elmer Lambda 950 following the dispersion of the SWNT sample in D₂O and sodium cholate using 30 W sonication for 5 minutes, and high speed centrifugation.

Atomic force microscopy images. AFM images were obtained using a Nanoscope IV Digital

Instruments AFM in tapping mode, using samples prepared by spin-coating the spontaneously-dissolved fraction onto a mica substrate.

Small Angle Neutron Scattering measurements. SANS data were collected on the LOQ beam line at the ISIS pulsed neutron source, using neutrons of wavelengths 2.2 to 10 Å recorded by time-of-flight at a 64 cm square position sensitive detector, 4.19 m from the sample. Samples were contained in a 2 mm path fused silica cell inside an evacuated closed cycle He refrigerator (CCR). The solutions were made *in situ* on the beam line by condensation of 2cm⁻³ deuterammonia (ND₃) onto excess (20mg) ARC SWNTs and sodium metal at a carbon to metal ratio of 1:24. The 12 mm diameter neutron beam sampled only the upper half of the cell, containing spontaneously dissolved material. Background subtractions, which were dominated by scattering from the cell and CCR vacuum tails, were measured for the same cell filled with ND₃. Data were corrected for wavelength dependent sample transmissions and detector efficiencies, and then scaled to absolute scattering cross-section, $I(Q)$, by comparison to data from a partially deuterated polystyrene standard.

SANS data were fitted using the program Fish⁵³, which employs a standard iterative linear least squares method involving computation of first derivatives of each calculated data point with respect to each parameter in the model. The best fit shown in figure 1b was obtained by using a model in which the data were fitted in two regimes.

The low- Q region ($Q < 0.055 \text{ Å}^{-1}$) was fitted to a power law;

$$I(Q) = CQ^{-D} + B$$

where C is a constant and B the incoherent background, and the fractal dimension was found to be $D = 2.65$.

The high- Q region ($Q \geq 0.055 \text{ Å}^{-1}$) was fitted to the expression⁵³;

$$I(Q) \propto N_p V_p^2 \Delta\rho^2 P(Q) + B$$

where N_P is the number concentration of scattering bodies, V_P is the volume of one scattering body, $\Delta\rho$ is the difference in neutron scattering length density between the scattering bodies and the solvent (the so-called ‘contrast’), B is again the incoherent background, and $P(Q)$ is the form (shape) factor, corresponding to randomly oriented cylindrical rods of effective diameter ~ 20 Å. This model yields a nanotube concentration of approximately 0.17 - 0.75 vol % or 2-12 mg/mL, for effective rod diameters of 20 Å and 15 Å, respectively.

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